

## SUPPLEMENTARY INFORMATION

### Deriving an Electric Wave Equation from Weber's Electrodynamics

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Further details relating to the integrations performed in the main manuscript are given here:

1. Integration on spherical shell of radius  $r$  for the expression below

$$\int (\hat{r} \cdot \vec{v}) \hat{r} ds$$

where  $\hat{r}$  is a unit vector pointing from origin to shell,  $\vec{v}$  is a constant vector.

$$\hat{r} = \sin \theta \cos \varphi \hat{x} + \sin \theta \sin \varphi \hat{y} + \cos \theta \hat{z}$$

$$\vec{v} = v_x \hat{x} + v_y \hat{y} + v_z \hat{z}$$

$$ds = r^2 \sin \theta d\varphi d\theta$$

where  $\theta$  and  $\varphi$  are angles in spherical coordinates.

$$\begin{aligned} \int (\hat{r} \cdot \vec{v}) \hat{r} ds &= \int_0^{2\pi} d\varphi \int_0^\pi (\sin \theta \cos \varphi v_x + \sin \theta \sin \varphi v_y + \cos \theta v_z) (\sin \theta \cos \varphi \hat{x} + \sin \theta \sin \varphi \hat{y} \\ &\quad + \cos \theta \hat{z}) r^2 \sin \theta d\theta = \frac{4}{3} \pi r^2 (v_x \hat{x} + v_y \hat{y} + v_z \hat{z}) = \frac{4}{3} \pi r^2 \vec{v} \end{aligned} \quad (A1)$$

2. Integration on spherical shell of radius  $r$  for the expression below

$$\int (\hat{r} \cdot \vec{v})^2 (\hat{r} \cdot \vec{u}) \hat{r} ds$$

where  $\vec{u}$  is a constant vector.

$$\vec{u} = u_x \hat{x} + u_y \hat{y} + u_z \hat{z}$$

Thus

$$\begin{aligned} \int (\hat{r} \cdot \vec{v})^2 (\hat{r} \cdot \vec{u}) \hat{r} ds &= \int_0^{2\pi} d\varphi \int_0^\pi (\sin \theta \cos \varphi v_x + \sin \theta \sin \varphi v_y + \cos \theta v_z)^2 (\sin \theta \cos \varphi u_x \\ &\quad + \sin \theta \sin \varphi u_y + \cos \theta u_z) (\sin \theta \cos \varphi \hat{x} + \sin \theta \sin \varphi \hat{y} + \cos \theta \hat{z}) r^2 \sin \theta d\theta \\ &= \frac{4}{15} \pi r^2 ((\vec{v} \cdot \vec{v}) \vec{u} + 2(\vec{v} \cdot \vec{u}) \vec{v}) \end{aligned} \quad (A2)$$

3. Integration in equation (6)

$$\vec{F}_c = -\frac{1}{4\pi\epsilon_0 r^3} \rho_-(\vec{0}) dV \int (2\rho_-(\vec{0}) + 2\vec{r} \cdot \nabla \rho_-(\vec{0}) - 2\rho) \vec{r} ds dr$$

Because of spherical symmetry,  $\int \vec{r} ds = \int -\vec{r} ds = 0$ . Thus

$$\vec{F}_c = -\frac{1}{4\pi\epsilon_0 r^3} \rho_-(\vec{0}) dV \int (2\vec{r} \cdot \nabla \rho_-(\vec{0})) \vec{r} ds dr$$

Using equation (A1), we obtain

$$\vec{F}_c = -\frac{1}{4\pi\epsilon_0 r^3} \rho_-(\vec{0}) dV 2r^2 dr \frac{4}{3} \pi r^2 \nabla \rho_-(\vec{0}) = -\frac{2r}{3\epsilon_0} \rho_-(\vec{0}) dV \nabla \rho_-(\vec{0}) dr$$

4. Integration in equation (10)

$$\begin{aligned}\vec{F}_v = & -\frac{1}{4\pi\epsilon_0 r^3} \rho_-(\vec{0}) dV \int \frac{1}{c^2} \left( (\vec{r} \cdot \nabla \vec{v}_-(\vec{0})) \cdot (\vec{r} \cdot \nabla \vec{v}_-(\vec{0})) - \frac{3}{2r^2} (\vec{r} \cdot (\vec{r} \cdot \nabla \vec{v}_-(\vec{0})))^2 \right) (\rho_-(\vec{0}) + \vec{r} \\ & \cdot \nabla \rho_-(\vec{0})) \vec{r} ds dr \\ & + \frac{1}{4\pi\epsilon_0 r^3} \rho_-(\vec{0}) dV \int \frac{1}{c^2} \left( (2\vec{v}_-(\vec{0}) + \vec{r} \cdot \nabla \vec{v}_-(\vec{0})) \cdot (2\vec{v}_-(\vec{0}) + \vec{r} \cdot \nabla \vec{v}_-(\vec{0})) \right. \\ & \left. - \frac{3}{2r^2} (\vec{r} \cdot (2\vec{v}_-(\vec{0}) + \vec{r} \cdot \nabla \vec{v}_-(\vec{0})))^2 \right) (2\rho - \rho_-(\vec{0}) - \vec{r} \cdot \nabla \rho_-(\vec{0})) \vec{r} ds dr\end{aligned}$$

We drop the higher order terms, which has a product of  $\vec{r} \cdot \nabla \rho_-(\vec{0})$  and  $\vec{r} \cdot \nabla \vec{v}_-(\vec{0})$ .

$$\begin{aligned}\vec{F}_v = & -\frac{1}{4\pi\epsilon_0 r^3} \rho_-(\vec{0}) dV \int \frac{1}{c^2} \left( (\vec{r} \cdot \nabla \vec{v}_-(\vec{0})) \cdot (\vec{r} \cdot \nabla \vec{v}_-(\vec{0})) \right. \\ & \left. - \frac{3}{2r^2} (\vec{r} \cdot (\vec{r} \cdot \nabla \vec{v}_-(\vec{0})))^2 \right) (\rho_-(\vec{0})) \vec{r} ds dr \\ & + \frac{1}{4\pi\epsilon_0 r^3} \rho_-(\vec{0}) dV \int \frac{1}{c^2} \left( (2\vec{v}_-(\vec{0}) + \vec{r} \cdot \nabla \vec{v}_-(\vec{0})) \cdot (2\vec{v}_-(\vec{0}) + \vec{r} \cdot \nabla \vec{v}_-(\vec{0})) \right. \\ & \left. - \frac{3}{2r^2} (\vec{r} \cdot (2\vec{v}_-(\vec{0}) + \vec{r} \cdot \nabla \vec{v}_-(\vec{0})))^2 \right) (2\rho - \rho_-(\vec{0})) \vec{r} ds dr \\ & + \frac{1}{4\pi\epsilon_0 r^3} \rho_-(\vec{0}) dV \int \frac{1}{c^2} \left( (2\vec{v}_-(\vec{0})) \cdot (2\vec{v}_-(\vec{0})) - \frac{3}{2r^2} (\vec{r} \cdot (2\vec{v}_-(\vec{0})))^2 \right) (-\vec{r} \\ & \cdot \nabla \rho_-(\vec{0})) \vec{r} ds dr\end{aligned}$$

Because of spherical symmetry,  $\int \vec{r} ds = \int -\vec{r} ds = 0$ . Thus the first and second terms of the above equation equal zero.

$$\vec{F}_v = \frac{1}{4\pi\epsilon_0 r^3} \rho_-(\vec{0}) dV \int \frac{1}{c^2} \left( (2\vec{v}_-(\vec{0})) \cdot (2\vec{v}_-(\vec{0})) - \frac{3}{2r^2} (\vec{r} \cdot (2\vec{v}_-(\vec{0})))^2 \right) (-\vec{r} \cdot \nabla \rho_-(\vec{0})) \vec{r} ds dr$$

Using equation (A1 & A2), we obtain

$$\begin{aligned}\vec{F}_v = & \frac{1}{4\pi\epsilon_0 r^3} \rho_-(\vec{0}) dV \frac{4}{c^2} (\vec{v}_-(\vec{0}) \cdot \vec{v}_-(\vec{0})) r^2 dr \frac{-4}{3} \pi r^2 \nabla \rho_-(\vec{0}) \\ & + \frac{1}{4\pi\epsilon_0 r^3} \rho_-(\vec{0}) dV \frac{4}{c^2} \frac{3}{2r^2} r^2 dr \frac{4}{15} \pi r^2 ((\vec{v}_-(\vec{0}) \cdot \vec{v}_-(\vec{0})) \nabla \rho_-(\vec{0}) \\ & + 2 (\vec{v}_-(\vec{0}) \cdot \nabla \rho_-(\vec{0})) \vec{v}_-(\vec{0})) \\ & = -\frac{14r}{15\epsilon_0} \rho_-(\vec{0}) dV \frac{1}{c^2} (\vec{v}_-(\vec{0}) \cdot \vec{v}_-(\vec{0})) \nabla \rho_-(\vec{0}) dr \\ & + \frac{4r}{5\epsilon_0} \rho_-(\vec{0}) dV \frac{1}{c^2} (\vec{v}_-(\vec{0}) \cdot \nabla \rho_-(\vec{0})) \vec{v}_-(\vec{0}) dr\end{aligned}$$

5. Integration in equation (14)

$$\begin{aligned}\vec{F}_a = & -\frac{1}{4\pi\epsilon_0 r^3} \rho_-(\vec{0}) dV \int \frac{2\rho}{c^2} \vec{r} \cdot (\vec{r} \cdot \nabla \vec{a}_-(\vec{0})) \vec{r} ds dr \\ & - \frac{1}{4\pi\epsilon_0 r^3} \rho_-(\vec{0}) dV \int \frac{2}{c^2} (\rho_+(\vec{0}) + \vec{r} \cdot \nabla \rho_+(\vec{0})) \vec{r} \cdot \vec{a}_-(\vec{0}) \vec{r} ds dr\end{aligned}$$

We drop the higher order terms, which has a product of  $\vec{r} \cdot \nabla \rho_+(\vec{0})$  and  $\vec{r} \cdot 2\vec{a}_-(\vec{0})$ .

$$\begin{aligned}\vec{F}_a = & -\frac{1}{4\pi\epsilon_0 r^3} \rho_-(\vec{0}) dV \int \frac{2\rho}{c^2} \vec{r} \cdot (\vec{r} \cdot \nabla \vec{a}_-(\vec{0})) \vec{r} ds dr \\ & - \frac{1}{4\pi\epsilon_0 r^3} \rho_-(\vec{0}) dV \int \frac{2}{c^2} (\rho_+(\vec{0})) \vec{r} \cdot \vec{a}_-(\vec{0}) \vec{r} ds dr\end{aligned}$$

Because of spherical symmetry,  $\int \vec{r} \cdot (\vec{r} \cdot \vec{u}) \vec{r} ds = \int (-\vec{r}) \cdot ((-\vec{r}) \cdot \vec{u}) (-\vec{r}) ds = 0$ . Thus

$$\vec{F}_a = -\frac{1}{4\pi\epsilon_0 r^3} \rho_-(\vec{0}) dV \int \frac{2}{c^2} (\rho_+(\vec{0})) \vec{r} \cdot \vec{a}_-(\vec{0}) \vec{r} ds dr$$

We also use the approximation  $\rho_+(0) \approx \rho$ , and use equation (A1). Then

$$\begin{aligned}
\vec{F}_a &= -\frac{1}{4\pi\epsilon_0 r^3} \rho_-(\vec{0}) dV \int \frac{2}{c^2} \rho \vec{r} \cdot \vec{a}_-(\vec{0}) \vec{r} ds dr = -\frac{1}{4\pi\epsilon_0 r^3} \rho_-(\vec{0}) dV \frac{2}{c^2} \rho r^2 dr \frac{4}{3} \pi r^2 \vec{a}_-(\vec{0}) \\
&= -\frac{2r}{3\epsilon_0} \rho_-(\vec{0}) dV \frac{\rho}{c^2} \vec{a}_-(\vec{0}) dr
\end{aligned}$$