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Oscillatory Behavior of a Type of Generalized Proportional Fractional Differential Equations with Forcing and Damping Terms

Jehad Alzabut 1,*,†, James Viji 2,†, Velu Muthulakshmi 2,†, and Weerawat Sudsutad 3,†

- Department of Mathematics and General Sciences, Prince Sultan University, Riyadh 11586, Saudi Arabia
- Department of Mathematics, Periyar University, Salem 636 011, Tamilnadu, India; vijimaths25@gmail.com (J.V.); vmuthumath@periyaruniversity.ac.in (V.M.)
- Department of General Education, Navamindradhiraj University, Bangkok 10300, Thailand; weerawat@nmu.ac.th
- * Correspondence: jalzabut@psu.edu.sa
- † These authors contributed equally to this work.

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Abstract: In this paper, we study the oscillatory behavior of solutions for a type of generalized proportional fractional differential equations with forcing and damping terms. Several oscillation criteria are established for the proposed equations in terms of Riemann-Liouville and Caputo settings. The results of this paper generalize some existing theorems in the literature. Indeed, it is shown that for particular choices of parameters, the obtained conditions in this paper reduce our theorems to some known results. Numerical examples are constructed to demonstrate the effectiveness of the our main theorems. Furthermore, we present and illustrate an example which does not satisfy the assumptions of our theorem and whose solution demonstrates nonoscillatory behavior.

Keywords: generalized proportional fractional operator; oscillation criteria; nonoscillatory behavior; damping and forcing terms

1. Introduction

Fractional calculus is a mathematical branch investigating the properties of derivatives and integrals of non-integer orders. The significance of this subject falls in the fact that the fractional derivative has the feature of nonlocal nature. This property makes these derivatives suitable to simulate more physical phenomena such as earthquake vibrations, polymers, and so forth; see, for example, References [1–10] and the references cited therein.

In recent years, there have appeared different types of fractional derivatives. However, it has been realized that most of these derivatives lose some of their basic properties that classical derivatives have such as the product rule and the chain rule. Fortunately, Khalil et al. [11] defined a new well-behaved fractional derivative, called the "conformable fractional derivative", which depends entirely on the classical limit definition of the derivative. Thereafter, researchers developed the conformable derivative and obtained different results exposing its features [12–14]. Recently, Jarad et al. [15] introduced the generalized proportional fractional (GPF) derivative of Caputo and Riemann-Liouville type involving exponential functions in their kernels. The GPF derivative not only preserves classical properties but also verifies semi group property and of nonlocal behavior. For recent results involving GPF derivative, one can refer to References [16–18].

In 2012, Grace et al. [19] initiated the study of oscillation theory for fractional differential equations. Thereafter, many researchers have investigated the oscillatory properties of fractional differential

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equations; see for instance References [20–25]. In 2019, Aphithana et al. [24] studied forced oscillatory properties of solutions to the conformable initial value problem of the form

$$\begin{cases} aD^{1+\alpha,\rho}x(t) + p(t)_a D^{\alpha,\rho}x(t) + q(t)f(x(t)) = g(t), & t > a, \\ \lim_{t \to a^+} a \mathcal{J}^{j-\alpha,\rho}x(t) = b_j & (j = 1, 2, \dots, m), \end{cases}$$

where $m = \lceil \alpha \rceil$, $0 < \rho \le 1$, p, $g \in \mathbb{C}(\mathbb{R}^+, \mathbb{R})$, $q \in \mathbb{C}(\mathbb{R}^+, \mathbb{R}^+)$, $f \in \mathbb{C}(\mathbb{R}, \mathbb{R})$ are continuous functions, ${}_aD^{\alpha,\rho}$ is the left conformable derivative of order $\alpha \in \mathbb{C}$ of x, $Re(\alpha) \ge 0$ in the Riemann-Liouville setting and ${}_a\mathcal{J}^{j-\alpha,\rho}$ is the left conformable integral operator of order $j-\alpha \in \mathbb{C}$, $b_j \in \mathbb{R}$, $j=1,2,\ldots,m$.

They also studied the forced oscillation of conformable initial value problems in the Caputo setting of the form

$$\begin{cases} {}^{C}_{a}D^{1+\alpha,\rho}x(t) + p(t)^{C}_{a}D^{\alpha,\rho}x(t) + q(t)f(x(t)) = g(t), & t > a, \\ {}^{k}_{a}D^{\rho}x(a) = b_{k} & (k = 0, 1, \dots, m-1), \end{cases}$$

where $m = \lceil \alpha \rceil$, $0 < \rho \le 1$, and ${}^{\mathbb{C}}_{a}D^{\alpha,\rho}$ is the left conformable derivative of order $\alpha \in \mathbb{C}$ of x, $Re(\alpha) \ge 0$ in the Caputo setting.

In 2020, Sudsutad et al. [26] established some oscillation criteria for the following generalized proportional fractional differential equation

$$\begin{cases} aD^{\alpha,\rho}x(t) + \xi_1(t,x(t)) = \mu(t) + \xi_2(t,x(t)), & t > a \ge 0, \\ \lim_{t \to a^+} aI^{j-\alpha,\rho}x(t) = b_j & (j = 1,2,\ldots,n), \end{cases}$$

with $n = \lceil \alpha \rceil$, ${}_aD^{\alpha,\rho}$ is the generalized proportional fractional derivative operator of order $\alpha \in \mathbb{C}$, $Re(\alpha) \geq 0$, $0 < \rho \leq 1$ in the Riemann-Liouville setting and ${}_aI^{\alpha,\rho}$ is the generalized proportional fractional integral operator.

In this paper, motivated by the above papers, we establish some sufficient conditions for forced oscillation criteria of all solutions of the generalized proportional fractional (GPF) initial value problem with damping term in the Riemann-Liouville type of the form:

$$\begin{cases}
 {a}D^{1+\alpha,\rho}y(l) + p(l){a}D^{\alpha,\rho}y(l) + q(l)f(y(l)) = g(l), & l > a \ge 0, \\
 \lim_{l \to a^{+}} aI^{j-\alpha,\rho}y(l) = b_{j} & (j = 1, 2, ..., m),
\end{cases}$$
(1)

where $m = \lceil \alpha \rceil$, $0 < \rho \le 1$, ${}_aD^{\alpha,\rho}$ is the left GPF derivative of order $\alpha \in \mathbb{C}$ of y, $Re(\alpha) \ge 0$ in the Riemann-Liouville setting and ${}_aI^{j-\alpha,\rho}$ is the left GPF integral of order $j-\alpha \in \mathbb{C}$, $Re(j-\alpha) > 0$, $b_j \in \mathbb{R}$, $j=1,2,\ldots,m$ and $p,g \in \mathbb{C}(\mathbb{R}^+,\mathbb{R})$, $q \in \mathbb{C}(\mathbb{R}^+,\mathbb{R}^+)$, $f \in \mathbb{C}(\mathbb{R},\mathbb{R})$.

Moreover, we study the forced oscillation criteria of all solutions of the GPF initial value problem with damping term in the Caputo type of the form

where $n = \lceil \alpha \rceil$, $0 < \rho \le 1$, ${}^C_a D^{\alpha,\rho}$ is the left GPF derivative of order $\alpha \in \mathbb{C}$ of y, $Re(\alpha) \ge 0$ in the Caputo setting and $D^{k,\rho} = \underbrace{D^\rho D^\rho \cdots D^\rho}$, and D^ρ is the proportional derivative defined in Reference [13].

We claim that the results of this paper improve and generalize previously existing oscillation results in Reference [24].

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Definition 1. The solution y of problem (1) (respectively (2)) is called oscillatory if it has arbitrarily large zeros on $(0, \infty)$; otherwise, it is called nonoscillatory. An equation is called oscillatory if all its solutions are oscillatory.

2. Preliminaries

In this section, we provide some basic definitions and results which will be used throughout this paper. For the justifications and proofs, the reader can consult References [13,15].

Definition 2. [15] (Modified Conformable Derivatives).

For $\rho \in [0,1]$, let the functions $k_0, k_1 : [0,1] \times \mathbb{R} \to [0,\infty)$ be continuous such that for all $l \in \mathbb{R}$ we have

$$\lim_{\rho \to 0^+} k_1(\rho, l) = 1, \quad \lim_{\rho \to 0^+} k_0(\rho, l) = 0, \quad \lim_{\rho \to 1^-} k_1(\rho, l) = 0, \quad \lim_{\rho \to 1^-} k_0(\rho, l) = 1, \tag{3}$$

and $k_1(\rho, l) \neq 0$, $\rho \in [0, 1)$, $k_0(\rho, l) \neq 0$, $\rho \in (0, 1]$.

Then, Anderson et al. [13] defined the modified conformable differential operator of order ρ by

$$D^{\rho}f(l) = k_1(\rho, l)f(l) + k_0(\rho, l)f'(l), \tag{4}$$

provided that the right-hand side exists at $l \in \mathbb{R}$ and $f'(l) = \frac{d}{dl}f$. The derivative given in (4) is called a proportional derivative. For more details about the control theory of the proportional derivatives and its component functions k_0 and k_1 , we refer the reader to [27].

Of special interest, we shall restrict ourselves to the case when $k_1(\rho, l) = (1 - \rho)$ and $k_0(\rho, l) = \rho$. Therefore, (4) becomes

$$D^{\rho}f(l) = (1 - \rho)f(l) + \rho f'(l). \tag{5}$$

Notice that $\lim_{\rho \to 0^+} D^\rho f(l) = f(l)$ and $\lim_{\rho \to 1^-} D^\rho f(l) = f'(l)$. It is clear that the derivative (5) is somehow more general than the conformable derivative which does not tend to the original function as ρ tends to 0.

To find the associated integral to the proportional derivative in (5), we solve the following equation

$$D^{\rho}g(l) = (1-\rho)g(l) + \rho g'(l) = f(l), \quad l > a.$$

The above equation is a first order linear differential equation and its solution is given by

$$g(l) = \frac{1}{\rho} \int_a^l e^{\frac{\rho - 1}{\rho}(l - s)} f(s) ds.$$

Define the proportional integral associated to D^{ρ} by

$$_{a}I^{1,\rho}f(l) = \frac{1}{\rho} \int_{a}^{l} e^{\frac{\rho-1}{\rho}(l-s)} f(s) ds,$$
 (6)

where we accept that $_aI^{0,\rho}f(l)=f(l)$.

Lemma 1. [15] Let f be defined on $[a, \infty)$ and differentiable on (a, ∞) and $\rho \in (0, 1]$. Then, we have

$$_{a}I^{1,\rho}D^{\rho}f(l) = f(l) - e^{\frac{\rho-1}{\rho}(l-a)}f(a).$$
 (7)

Definition 3. [15] For $\rho \in (0,1]$ and $\alpha \in \mathbb{C}$, $Re(\alpha) > 0$, we define the left GPF integral of f by

$$({}_{a}I^{\alpha,\rho}f)(l) = \frac{1}{\rho^{\alpha}\Gamma(\alpha)} \int_{a}^{l} e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1}f(s)ds = \rho^{-\alpha} e^{\frac{\rho-1}{\rho}l} {}_{a}I^{\alpha} \left(e^{\frac{1-\rho}{\rho}l}f(l)\right),$$
 (8)

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where $_aI^{\alpha}$ is the left Riemann-Liouville fractional integral of order α . The right GPF integral ending at b, however, can be defined by

$$(I_b^{\alpha,\rho}f)(l) = \frac{1}{\rho^{\alpha}\Gamma(\alpha)} \int_l^b e^{\frac{\rho-1}{\rho}(s-l)} (s-l)^{\alpha-1} f(s) ds = \rho^{-\alpha} e^{\frac{\rho-1}{\rho}l} I_b^{\alpha} \left(e^{\frac{1-\rho}{\rho}l} f(l) \right), \tag{9}$$

where I_h^{α} is the right Riemann-Liouville fractional integral of order α .

Definition 4. [15] For $\rho \in (0,1]$ and $\alpha \in \mathbb{C}$, $Re(\alpha) \geq 0$, we define the left GPF derivative of f by

$$({}_{a}D^{\alpha,\rho}f)(l) = D^{n,\rho}{}_{a} I^{n-\alpha,\rho}f(l)$$

$$= \frac{D^{n,\rho}_{l}}{\rho^{n-\alpha}\Gamma(n-\alpha)} \int_{a}^{l} e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{n-\alpha-1} f(s) ds.$$
(10)

The right GPF derivative ending at b is defined by

$$(D_b^{\alpha,\rho}f)(l) = {}_{\ominus}D^{n,\rho} I_b^{n-\alpha,\rho}f(l)$$

$$= \frac{{}_{\ominus}D_l^{n,\rho}}{\rho^{n-\alpha}\Gamma(n-\alpha)} \int_l^b e^{\frac{\rho-1}{\rho}(s-l)} (s-l)^{n-\alpha-1} f(s) ds, \tag{11}$$

where $n = [Re(\alpha)] + 1$.

If we let $\rho = 1$ in Definition 4 , then one can obtain the left and right Riemann-Liouville fractional derivatives as in [6]. Moreover, it is clear that

$$\lim_{\alpha \to 0} D^{\alpha,\rho} f(l) = f(l) \text{ and } \lim_{\alpha \to 1} D^{\alpha,\rho} f(l) = D^{\rho} f(l).$$

Lemma 2. [15] Let $Re(\alpha) > 0$, $n = -[-Re(\alpha)]$, $f \in L_1(a,b)$ and $({}_aI^{\alpha,\rho}f)(l) \in AC^n[a,b]$. Then,

$$({}_{a}I^{\alpha,\rho}{}_{a}D^{\alpha,\rho}f)(l) = f(l) - e^{\frac{\rho-1}{\rho}(l-a)} \sum_{j=1}^{n} ({}_{a}I^{j-\alpha,\rho}f)(a^{+}) \frac{(l-a)^{\alpha-j}}{\rho^{\alpha-j}\Gamma(\alpha+1-j)}.$$
 (12)

Definition 5. [13] (Partial Conformable Derivatives). Let $\rho \in [0,1]$, and let the functions $k_0, k_1 : [0,1] \times \mathbb{R} \to [0,\infty)$ be continuous and satisfy (3). Given a function $f: \mathbb{R}^2 \to \mathbb{R}$ such that $\frac{\partial}{\partial l} f(l,s)$ exists for each fixed $s \in \mathbb{R}$, define the partial differential operator D_l^ρ via

$$D_l^{\rho} f(l,s) = k_1(\rho, l) f(l,s) + k_0(\rho, l) \frac{\partial}{\partial l} f(l,s). \tag{13}$$

Definition 6. [13] (Conformable Exponential Function). Let $\rho \in (0,1]$, the points $s,l, \in \mathbb{R}$ with $s \leq l$, and let the function $p:[s,l] \to \mathbb{R}$ be continuous. Let $k_0,k_1:[0,1] \times \mathbb{R} \to [0,\infty)$ be continuous and satisfy (3), with p/k_0 and k_1/k_0 Riemann integrable on [s,l]. Then the exponential function with respect to D^ρ in (4) is defined to be

$$e_p(l,s) := e^{\int_s^l \frac{p(\tau) - k_1(\rho, \tau)}{k_0(\rho, \tau)} d\tau}, \quad e_0(l,s) := e^{-\int_s^l \frac{k_1(\rho, \tau)}{k_0(\rho, \tau)} d\tau}. \tag{14}$$

Using (4) and (14), we have the following basic results.

Lemma 3. [13] (Basic Derivatives). Let the conformable differential operator D^{ρ} be given as in (4), where $\rho \in [0,1]$. Let the function $p:[s,l] \to \mathbb{R}$ be continuous. Let $k_0, k_1:[0,1] \times \mathbb{R} \to [0,\infty)$ be continuous and satisfy (3), with p/k_0 and k_1/k_0 Riemann integrable on [s,l]. Assume the functions f and g are differentiable as needed. Then

(i)
$$D^{\rho}[af + bg] = aD^{\rho}[f] + bD^{\rho}[g]$$
 for all $a, b \in \mathbb{R}$;

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- (ii) $D^{\rho}c = ck_1(\rho, \cdot)$ for all constants $c \in \mathbb{R}$; (iii) $D^{\rho}[fg] = fD^{\rho}[g] + gD^{\rho}[f] fgk_1(\rho, \cdot)$; (iv) $D^{\rho}[f/g] = \frac{gD^{\rho}[f] fD^{\rho}[g]}{g^2} + \frac{f}{g}k_1(\rho, \cdot)$;
- (v) for $\rho \in (0,1]$ and fixed $s \in \mathbb{R}$, the exponential function satisfies

$$D_l^{\rho}[e_p(l,s)] = p(l)e_p(l,s)$$

for $e_p(l,s)$ given in (14);

(vi) for $\rho \in (0,1]$ and for the exponential function e_0 given in (14), we have

$$D^{\rho} \left[\int_a^l \frac{f(s)e_0(l,s)}{k_0(\rho,s)} ds \right] = f(l).$$

Definition 7. [15] For $\rho \in (0,1]$ and $\alpha \in \mathbb{C}$ with $Re(\alpha) \geq 0$, we define the left GPF derivative of Caputo type starting at a by

The right GPF derivative of Caputo ending at b is defined by

$$({}^{C}D_{b}^{\alpha,\rho}f)(l) = I_{b}^{n-\alpha,\rho} \left({}_{\ominus}D^{n,\rho}f\right)(l)$$

$$= \frac{1}{\rho^{n-\alpha}\Gamma(n-\alpha)} \int_{l}^{b} e^{\frac{\rho-1}{\rho}(s-l)} (s-l)^{n-\alpha-1} ({}_{\ominus}D^{n,\rho}f)(s) ds, \tag{16}$$

where $n = [Re(\alpha)] + 1$.

Lemma 4. [15] For $\rho \in (0,1]$ and $n = [Re(\alpha)] + 1$, we have

$${}_{a}I^{\alpha,\rho}({}_{a}^{C}D^{\alpha,\rho}f)(l) = f(l) - \sum_{k=0}^{n-1} \frac{(D^{k,\rho}f)(a)}{\rho^{k}k!}(l-a)^{k}e^{\frac{\rho-1}{\rho}}(l-a). \tag{17}$$

Proposition 1. [15] Let $\alpha, \beta \in \mathbb{C}$ be such that $Re(\alpha) \geq 0$ and $Re(\beta) > 0$. Then, for any $0 < \rho \leq 1$ and

$$\begin{split} &(i) \quad \left({}_aI^{\alpha,\rho}e^{\frac{\rho-1}{\rho}l}(l-a)^{\beta-1}\right)(y) = \frac{\Gamma(\beta)}{\Gamma(\beta+\alpha)\rho^\alpha}e^{\frac{\rho-1}{\rho}y}(y-a)^{\beta+\alpha-1}, \quad Re(\alpha) > 0. \\ &(ii) \quad \left({}_aD^{\alpha,\rho}e^{\frac{\rho-1}{\rho}l}(l-a)^{\beta-1}\right)(y) = \frac{\rho^\alpha\Gamma(\beta)}{\Gamma(\beta-\alpha)}e^{\frac{\rho-1}{\rho}y}(y-a)^{\beta-\alpha-1}, \quad Re(\alpha) \geq 0. \\ &(iii) \quad \left({}_a^CD^{\alpha,\rho}e^{\frac{\rho-1}{\rho}l}(l-a)^{\beta-1}\right)(y) = \frac{\rho^\alpha\Gamma(\beta)}{\Gamma(\beta-\alpha)}e^{\frac{\rho-1}{\rho}y}(y-a)^{\beta-\alpha-1}, \quad Re(\alpha) > n. \end{split}$$

$$(ii) \quad \left({}_aD^{\alpha,\rho}e^{\frac{\rho-1}{\rho}l}(l-a)^{\beta-1}\right)(y) = \frac{\rho^{\alpha}\Gamma(\beta)}{\Gamma(\beta-\alpha)}e^{\frac{\rho-1}{\rho}y}(y-a)^{\beta-\alpha-1}, \quad Re(\alpha) \geq 0.$$

$$(iii) \left({}^{C}_{a}D^{\alpha,\rho}e^{\frac{\rho-1}{\rho}l}(l-a)^{\beta-1} \right)(y) = \frac{\rho^{\alpha}\Gamma(\beta)}{\Gamma(\beta-\alpha)}e^{\frac{\rho-1}{\rho}y}(y-a)^{\beta-\alpha-1}, \quad Re(\alpha) > n.$$

3. Oscillation Results via Riemann-Liouville Operator

In this section, we establish the oscillation criteria for the GPF initial value problem (1). We prove our results under the following assumption:

$$(H) \ \ p \in \mathbb{C}(\mathbb{R}^+,\mathbb{R}), \ q \in \mathbb{C}(\mathbb{R}^+,\mathbb{R}^+), \ g \in \mathbb{C}(\mathbb{R}^+,\mathbb{R}), \ f \in C(\mathbb{R},\mathbb{R}) \ \text{with} \ \frac{f(u)}{u} > 0 \ \text{for all} \ u \neq 0.$$

For our convenience, we set the following notations:

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$$\Phi(l) := \Gamma(\alpha)e^{\frac{\rho-1}{\rho}(l-a)} \sum_{j=1}^{m} \frac{\rho^{j}b_{j}(l-a)^{\alpha-j}}{\Gamma(\alpha+1-j)}, \tag{18}$$

$$\Lambda(l,L) := \int_{a}^{L} e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left(\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M + l_1 I^{1,\rho} (\rho g(s)V(s))}{V(s)} \right) ds, \tag{19}$$

$$V(l) := \exp \int_{l_1}^{l} \frac{\rho p(\tau) - (1 - \rho)}{\rho} d\tau, \tag{20}$$

$$M := {}_{l_1}D^{\alpha,\rho}y(l_1)V(l_1), \quad M \text{ is an arbitrary constant.}$$
 (21)

Theorem 1. Assume that (H) holds. If

$$\lim_{l \to \infty} \inf l^{1-\alpha} \int_{L}^{l} e^{\frac{\rho - 1}{\rho}(l - s)} (l - s)^{\alpha - 1} \left[\frac{e^{\frac{\rho - 1}{\rho}(s - l_1)} M + l_1 I^{1,\rho} (\rho g(s) V(s))}{V(s)} \right] ds = -\infty$$
 (22)

and

$$\lim_{l \to \infty} \sup l^{1-\alpha} \int_{L}^{l} e^{\frac{\rho - 1}{\rho}(l - s)} (l - s)^{\alpha - 1} \left[\frac{e^{\frac{\rho - 1}{\rho}(s - l_1)} M + l_1 I^{1, \rho} (\rho g(s) V(s))}{V(s)} \right] ds = \infty, \tag{23}$$

for every sufficiently large L, where V(l) and M are defined as in (20) and (21) respectively, then every solution of problem (1) is oscillatory.

Proof. Suppose that y(l) is a nonoscillatory solution of problem (1). Without loss of generality, let L > a be large enough and $l_1 \ge L$ such that y(l) > 0 for all $l \ge l_1$. Using Lemma 3 (iii), Equations (5) and (13), we have

$$\begin{split} D^{\rho}[{}_{a}D^{\alpha,\rho}y(l)V(l)] &= {}_{a}D^{\alpha,\rho}y(l)D^{\rho}V(l) + V(l)D^{\rho}\left({}_{a}D^{\alpha,\rho}y(l)\right) - (1-\rho){}_{a}D^{\alpha,\rho}y(l)V(l) \\ &= {}_{a}D^{\alpha,\rho}y(l)D^{\rho}V(l) + V(l)\left[(1-\rho){}_{a}D^{\alpha,\rho}y(l) + \rho\frac{d}{dl}\left({}_{a}D^{\alpha,\rho}y(l)\right)\right] \\ &- (1-\rho){}_{a}D^{\alpha,\rho}y(l)V(l) \\ &= \rho\left[{}_{a}D^{1+\alpha,\rho}y(l) + p(l){}_{a}D^{\alpha,\rho}y(l)\right]V(l) \\ &= \rho\left[-q(l)f(y(l)) + g(l)\right]V(l) \\ &< \rho g(l)V(l). \end{split}$$

Taking the proportional integral operator $l_1I^{1,\rho}$ on both sides to the above inequality, we obtain

$$_{l_{1}}I^{1,\rho}\left(D^{\rho}\left[{}_{a}D^{\alpha,\rho}y(l)V(l)\right]\right) < {}_{l_{1}}I^{1,\rho}\left(\rho g(l)V(l)\right). \tag{24}$$

Using Lemma 1 on the L.H.S of (24), we have

$$_{a}D^{\alpha,\rho}y(l)<rac{e^{rac{
ho-1}{
ho}(l-l_{1})}M+_{l_{1}}I^{1,
ho}\left(
ho g(l)V(l)
ight)}{V(l)}.$$

Taking the left GPF integral operator ${}_aI^{\alpha,\rho}$ on both sides to the above inequality, we get

$${}_{a}I^{\alpha,\rho}\left({}_{a}D^{\alpha,\rho}y(l)\right) < {}_{a}I^{\alpha,\rho}\left[\frac{e^{\frac{\rho-1}{\rho}(l-l_{1})}M + {}_{l_{1}}I^{1,\rho}\left(\rho g(l)V(l)\right)}{V(l)}\right]. \tag{25}$$

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Using Lemma 2 on the L.H.S of (25), we have

$$y(l) - e^{\frac{\rho - 1}{\rho}(l - a)} \sum_{j=1}^{m} \frac{b_{j}(l - a)^{\alpha - j}}{\rho^{\alpha - j}\Gamma(\alpha + 1 - j)} <_{a}I^{\alpha, \rho} \left[\frac{e^{\frac{\rho - 1}{\rho}(l - l_{1})}M + _{l_{1}}I^{1, \rho}\left(\rho g(l)V(l)\right)}{V(l)} \right]. \tag{26}$$

Applying the left GPF integral formula on the R.H.S of (26), we have

$$y(l) < e^{\frac{\rho-1}{\rho}(l-a)} \sum_{j=1}^{m} \frac{b_{j}(l-a)^{\alpha-j}}{\rho^{\alpha-j}\Gamma(\alpha+1-j)}$$

$$+ \frac{1}{\rho^{\alpha}\Gamma(\alpha)} \int_{a}^{l} e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_{1})}M + l_{1}I^{1}\rho}(\rho g(s)V(s))}{V(s)} \right] ds,$$

for every sufficiently large *L*. If we multiply the above inequality by $\rho^{\alpha}\Gamma(\alpha)$, we get

$$\rho^{\alpha}\Gamma(\alpha)y(l) < \Gamma(\alpha)e^{\frac{\rho-1}{\rho}(l-a)} \sum_{j=1}^{m} \frac{\rho^{j}b_{j}(l-a)^{\alpha-j}}{\Gamma(\alpha+1-j)}
+ \int_{a}^{L} e^{\frac{\rho-1}{\rho}(l-s)}(l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_{1})} M + l_{1}I^{1,\rho}(\rho g(s)V(s))}{V(s)} \right] ds
+ \int_{L}^{l} e^{\frac{\rho-1}{\rho}(l-s)}(l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_{1})} M + l_{1}I^{1,\rho}(\rho g(s)V(s))}{V(s)} \right] ds$$

$$= \Phi(l) + \Lambda(l,L) + \int_{L}^{l} e^{\frac{\rho-1}{\rho}(l-s)}(l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_{1})} M + l_{1}I^{1,\rho}(\rho g(s)V(s))}{V(s)} \right] ds,$$
(27)

where $\Phi(l)$ and $\Lambda(l, L)$ are defined in (18) and (19), respectively. Multiplying (27) by $l^{1-\alpha}$, we get

$$0 < l^{1-\alpha} \rho^{\alpha} \Gamma(\alpha) y(l) < l^{1-\alpha} \Phi(l) + l^{1-\alpha} \Lambda(l, L)$$

$$+ l^{1-\alpha} \int_{L}^{l} e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_{1})} M + l_{1} I^{1,\rho} (\rho g(s) V(s))}{V(s)} \right] ds.$$

$$(28)$$

Let us consider the following two cases for $L_1 \geq L$. Case(i): Let $0 < \alpha \leq 1$. Then m = 1. Since $\left| e^{\frac{\rho - 1}{\rho}(l - a)} \right| \leq 1$ and the function $h_1(l) = \left(\frac{l - a}{l} \right)^{\alpha - 1}$ is decreasing for $\rho > 0$, $0 < \alpha < 1$, we get for $l \ge L_1$.

$$\left|l^{1-\alpha}\Phi(l)\right| = \left|e^{\frac{\rho-1}{\rho}(l-a)}\rho b_1\left(\frac{l-a}{l}\right)^{\alpha-1}\right| \le \rho|b_1|\left(\frac{L_1-a}{L_1}\right)^{\alpha-1} := C_1(L_1),\tag{29}$$

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and

$$\left| l^{1-\alpha} \Lambda(l,L) \right| = \left| l^{1-\alpha} \int_{a}^{L} e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left(\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M + l_1 I^{1,\rho} (\rho g(s)V(s))}{V(s)} \right) ds \right| \\
\leq \int_{a}^{L} \left| e^{\frac{\rho-1}{\rho}(l-s)} \right| \left(\frac{l-a}{l} \right)^{\alpha-1} \left| \frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M + l_1 I^{1,\rho} (\rho g(s)V(s))}{V(s)} \right| ds \\
\leq \int_{a}^{L} \left(\frac{L_1-a}{L_1} \right)^{\alpha-1} \left| \frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M + l_1 I^{1,\rho} (\rho g(s)V(s))}{V(s)} \right| ds \\
\leq C_2(L,L_1). \tag{30}$$

From (28), (29) and (30), we get, for $l \ge L_1$,

$$l^{1-\alpha} \int_{L}^{l} e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M + l_1 I^{1,\rho} \left(\rho g(s) V(s)\right)}{V(s)} \right] ds \geq -[C_1(L_1) + C_2(L,L_1)].$$

Since the R.H.S of the above inequality is a negative constant, it follows that

$$\lim_{l \to \infty} \inf l^{1-\alpha} \int_{L}^{l} e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M + l_1 I^{1,\rho} \left(\rho g(s) V(s) \right)}{V(s)} \right] ds > -\infty,$$

which leads to a contradiction with (22).

Case(ii): Let $\alpha > 1$. Then $m \ge 2$ and $\left(\frac{l-a}{l}\right)^{\alpha-1} < 1$ for $\alpha > 1$ and $\rho > 0$. Since $\left|e^{\frac{\rho-1}{\rho}(l-a)}\right| \le 1$ and the function $h_2(l) = (l-a)^{1-j}$ is decreasing for j > 1 and $\rho > 0$, for $l \ge L_1$, we have

$$\begin{aligned}
\left|l^{1-\alpha}\Phi(l)\right| &= \left|l^{1-\alpha}\Gamma(\alpha)e^{\frac{\rho-1}{\rho}(l-a)}\sum_{j=1}^{m}\frac{\rho^{j}b_{j}(l-a)^{\alpha-j}}{\Gamma(\alpha+1-j)}\right| \\
&\leq \Gamma(\alpha)\left(\frac{l-a}{l}\right)^{\alpha-1}\sum_{j=1}^{m}\frac{\rho^{j}|b_{j}|(l-a)^{1-j}}{\Gamma(\alpha+1-j)} \\
&\leq \Gamma(\alpha)\sum_{j=1}^{m}\frac{\rho^{j}|b_{j}|(L_{1}-a)^{1-j}}{\Gamma(\alpha+1-j)} \\
&\leq C_{3}(L_{1}),
\end{aligned} (31)$$

and

$$\left| l^{1-\alpha} \Lambda(l,L) \right| = \left| \int_{a}^{L} e^{\frac{\rho-1}{\rho}(l-s)} \left(\frac{l-s}{l} \right)^{\alpha-1} \left(\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M + l_1 I^{1,\rho} \left(\rho g(s) V(s) \right)}{V(s)} \right) ds \right|$$

$$\leq \int_{a}^{L} \left| \frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M + l_1 I^{1,\rho} \left(\rho g(s) V(s) \right)}{V(s)} \right| ds$$

$$:= C_4(L).$$
(32)

From (28), (31) and (32), we conclude that for $l \ge L_1$,

$$l^{1-\alpha} \int_{L}^{l} e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M + l_1 I^{1,\rho} \left(\rho g(s) V(s)\right)}{V(s)} \right] ds \ge -[C_3(L_1) + C_4(L)].$$

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Since, the R.H.S of the above inequality is a negative constant, it follows that

$$\lim_{l \to \infty} \inf l^{1-\alpha} \int_{L}^{l} e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M + l_1 I^{1,\rho} \left(\rho g(s) V(s) \right)}{V(s)} \right] ds > -\infty,$$

which is a contradiction to (22).

Therefore, y(l) is oscillatory. If y(l) is eventually negative, by a similar argument, we get a contradiction with condition (23). Hence the theorem. \Box

4. Oscillation Results via Caputo Operator

In this section, we establish the oscillation criteria for the GPF initial value problem (2) under the assumption (H):

We set

$$\Psi(l) := \Gamma(\alpha) \sum_{k=0}^{m-1} \frac{b_k}{\rho^{k-\alpha} k!} (l-a)^k e^{\frac{\rho-1}{\rho}(l-a)}, \tag{33}$$

$$\Omega(l,L) := \int_{a}^{L} e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left(\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M^* + l_1 I^{1,\rho} \left(\rho g(s) V(s)\right)}{V(s)} \right) ds, \tag{34}$$

$$M^* := {}^{C}_{a} D^{\alpha, \rho} y(a) V(a), \quad M^* \text{ is an arbitrary constant.}$$
 (35)

Theorem 2. Assume that (H) holds. If

$$\lim_{l \to \infty} \inf l^{1-n} \int_{L}^{l} e^{\frac{\rho - 1}{\rho}(l - s)} (l - s)^{\alpha - 1} \left[\frac{e^{\frac{\rho - 1}{\rho}(s - l_1)} M^* + l_1 I^{1,\rho} \left(\rho g(s) V(s)\right)}{V(s)} \right] ds = -\infty$$
 (36)

and

$$\lim_{l \to \infty} \sup l^{1-n} \int_{L}^{l} e^{\frac{\rho - 1}{\rho}(l - s)} (l - s)^{\alpha - 1} \left[\frac{e^{\frac{\rho - 1}{\rho}(s - l_1)} M^* + l_1 I^{1,\rho} \left(\rho g(s) V(s)\right)}{V(s)} \right] ds = \infty, \tag{37}$$

for every sufficiently large L, where V(l) and M^* are defined as in (20) and (35), respectively, then every solution of problem (2) is oscillatory.

Proof. Suppose that y(l) is a nonoscillatory solution of problem (2). Without loss of generality, let L > a be large enough and $l_1 \ge L$ such that y(l) > 0 for $l \ge l_1$. Using Lemma 3 (iii), Equations (5) and (13), we have

$$\begin{split} D^{\rho} [{}^{C}_{a}D^{\alpha,\rho}y(l)V(l)] &= {}^{C}_{a}D^{\alpha,\rho}y(l)D^{\rho}V(l) + V(l)D^{\rho} \left({}^{C}_{a}D^{\alpha,\rho}y(l)\right) - (1-\rho){}^{C}_{a}D^{\alpha,\rho}y(l)V(l) \\ &= {}^{C}_{a}D^{\alpha,\rho}y(l)D^{\rho}V(l) + V(l) \left[(1-\rho){}^{C}_{a}D^{\alpha,\rho}y(l) + \rho \frac{d}{dl} \left({}^{C}_{a}D^{\alpha,\rho}y(l)\right) \right] \\ &- (1-\rho){}^{C}_{a}D^{\alpha,\rho}y(l)V(l) \\ &= \rho \left[{}^{C}_{a}D^{1+\alpha,\rho}y(l) + p(l){}^{C}_{a}D^{\alpha,\rho}y(l) \right] V(l) \\ &= \rho \left[-q(l)f(y(l)) + g(l) \right] V(l) \\ &< \rho g(l)V(l). \end{split}$$

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Taking the proportional integral operator $l_1I^{1,\rho}$ on both sides to the above inequality, we obtain

$$_{l_{1}}I^{1,\rho}\left(D^{\rho}\left[_{a}^{C}D^{\alpha,\rho}y(l)V(l)\right]\right) <_{l_{1}}I^{1,\rho}\left(\rho g(l)V(l)\right). \tag{38}$$

Using Lemma (1) on the L.H.S of (38), we have

$${}_{a}^{C}D^{\alpha,\rho}y(l) < \frac{e^{\frac{\rho-1}{\rho}(l-l_{1})}M^{*} + {}_{l_{1}}I^{1,\rho}\left(\rho g(l)V(l)\right)}{V(l)}.$$

Applying the left GPF integral operator ${}_aI^{\alpha,\rho}$ on both sides to the above inequality, we get

$${}_{a}I^{\alpha,\rho}\left({}_{a}^{C}D^{\alpha,\rho}y(l)\right) < {}_{a}I^{\alpha,\rho}\left[\frac{e^{\frac{\rho-1}{\rho}(l-l_{1})}M^{*} + {}_{l_{1}}I^{1,\rho}\left(\rho g(l)V(l)\right)}{V(l)}\right]. \tag{39}$$

Using Lemma 4 on the L.H.S of (39), we have

$$y(l) - \sum_{k=0}^{n-1} \frac{(D^{k,\rho}y)(a)}{\rho^k k!} (l-a)^k e^{\frac{\rho-1}{\rho}(l-a)} < {}_{a}I^{\alpha,\rho} \left[\frac{e^{\frac{\rho-1}{\rho}(l-l_1)}M^* + {}_{l_1}I^{1,\rho} (\rho g(l)V(l))}{V(l)} \right]. \tag{40}$$

Applying the left GPF integral formula on the R.H.S of (40), we have

$$y(l) < \sum_{k=0}^{n-1} \frac{b_k}{\rho^k k!} (l-a)^k e^{\frac{\rho-1}{\rho}(l-a)}$$

$$+ \frac{1}{\rho^{\alpha} \Gamma(\alpha)} \int_a^l e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M^* + l_1 I^{1,\rho} \left(\rho g(s) V(s)\right)}{V(s)} \right] ds,$$

for every sufficiently large L. If we multiply the above inequality by $\rho^{\alpha}\Gamma(\alpha)$, we get

$$\rho^{\alpha}\Gamma(\alpha)y(l) < \rho^{\alpha}\Gamma(\alpha)\sum_{k=0}^{n-1} \frac{b_{k}}{\rho^{k}k!}(l-a)^{k} e^{\frac{\rho-1}{\rho}(l-a)}
+ \int_{a}^{L} e^{\frac{\rho-1}{\rho}(l-s)}(l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_{1})}M^{*} + l_{1}I^{1,\rho}\left(\rho g(s)V(s)\right)}{V(s)} \right] ds
+ \int_{L}^{l} e^{\frac{\rho-1}{\rho}(l-s)}(l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_{1})}M^{*} + l_{1}I^{1,\rho}\left(\rho g(s)V(s)\right)}{V(s)} \right] ds
= \Psi(l) + \Omega(l,L)
+ \int_{L}^{l} e^{\frac{\rho-1}{\rho}(l-s)}(l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_{1})}M^{*} + l_{1}I^{1,\rho}\left(\rho g(s)V(s)\right)}{V(s)} \right] ds,$$
(41)

where $\Psi(l)$ and $\Omega(l,L)$ are defined in (33) and (34), respectively. Multiplying (41) by l^{1-n} , we get

$$0 < l^{1-n}\Gamma(\alpha)y(l) < l^{1-n}\Psi(l) + l^{1-n}\Omega(l,L) + l^{1-n} \int_{L}^{l} e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_1)}M^* + l_1I^{1,\rho}(\rho g(s)V(s))}{V(s)} \right] ds.$$
 (42)

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Let us consider the following two cases for $L_1 \ge L$. Case(i): Let $0 < \alpha \le 1$. Then n = 1. Since $\left| e^{\frac{\rho - 1}{\rho}(l - a)} \right| \le 1$ and the function $h_3(l) = (l - s)^{\alpha - 1}$ is decreasing for $0 < \alpha < 1$, we get for $l \ge L_1$,

$$\left| l^{1-n} \Psi(l) \right| = \left| l^{1-n} \Gamma(\alpha) \sum_{k=0}^{n-1} \frac{b_k}{\rho^{k-\alpha} k!} (l-a)^k e^{\frac{\rho-1}{\rho}(l-a)} \right| \le \rho^{\alpha} \Gamma(\alpha) |b_0| := C_5(L), \tag{43}$$

and

$$\left| l^{1-n} \Omega(l,L) \right| = \left| l^{1-n} \int_{a}^{L} e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left(\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M^* + l_1 I^{1,\rho} (\rho g(s) V(s))}{V(s)} \right) ds \right| \\
\leq \int_{a}^{L} (l-s)^{\alpha-1} \left| \frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M^* + l_1 I^{1,\rho} (\rho g(s) V(s))}{V(s)} \right| ds \\
\leq \int_{a}^{L} (L_1-s)^{\alpha-1} \left| \frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M^* + l_1 I^{1,\rho} (\rho g(s) V(s))}{V(s)} \right| ds \\
\leq C_6(L,L_1). \tag{44}$$

Then, from (42) and $l \ge L_1$, we get

$$l^{1-n} \int_{L}^{l} e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M^* + l_1 I^{1,\rho} \left(\rho g(s) V(s)\right)}{V(s)} \right] ds \geq -[C_5(L) + C_6(L, L_1)].$$

Since, the R.H.S of the above inequality is a negative constant, it follows that

$$\lim_{l \to \infty} \inf l^{1-n} \int_{L}^{l} e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M^* + _{l_1} I^{1,\rho} \left(\rho g(s) V(s)\right)}{V(s)} \right] ds > -\infty,$$

which leads to a contradiction with the condition (36).

Case(ii): Let $\alpha > 1$. Then $n \ge 2$ and $\left(\frac{l-a}{l}\right)^{n-1} < 1$ for $n \ge 2$ and $\alpha > 1$. Since $\left|e^{\frac{\rho-1}{\rho}(l-a)}\right| \le 1$ and the function $h_4(l) = (l-a)^{k-n+1}$ is decreasing for k > n-1 and for $l \ge L_1$, we have

$$\begin{aligned}
\left| l^{1-n} \Psi(l) \right| &= \left| l^{1-n} \Gamma(\alpha) \sum_{k=0}^{n-1} \frac{b_k}{\rho^{k-\alpha} k!} (l-a)^k e^{\frac{\rho-1}{\rho}(l-a)} \right| \\
&= \left| \rho^{\alpha} \Gamma(\alpha) \left(\frac{l-a}{l} \right)^{n-1} \sum_{k=0}^{n-1} \frac{b_k}{\rho^k k!} (l-a)^{k-n+1} e^{\frac{\rho-1}{\rho}(l-a)} \right| \\
&\leq \rho^{\alpha} \Gamma(\alpha) \sum_{k=0}^{n-1} \frac{|b_k|}{\rho^k k!} (l-a)^{k-n+1} \\
&\leq \rho^{\alpha} \Gamma(\alpha) \sum_{k=0}^{n-1} \frac{|b_k| (L_1-a)^{k-n+1}}{\rho^k k!} \\
&= C_7(L_1),
\end{aligned} \tag{45}$$

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and

$$\left| l^{1-n} \Omega(l,L) \right| = \left| l^{1-n} \int_{a}^{L} e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left(\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M^* + l_1 I^{1,\rho} (\rho g(s) V(s))}{V(s)} \right) ds \right|
= \left| l^{\alpha-n} \int_{a}^{L} e^{\frac{\rho-1}{\rho}(l-s)} \left(\frac{l-s}{l} \right)^{\alpha-1} \left(\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M^* + l_1 I^{1,\rho} (\rho g(s) V(s))}{V(s)} \right) ds \right|
\leq \int_{a}^{L} \left| \frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M^* + l_1 I^{1,\rho} (\rho g(s) V(s))}{V(s)} \right| ds
:= C_8(L).$$
(46)

From Equations (42), (45) and (46), we conclude that for $l \ge L_1$,

$$l^{1-n} \int_{L}^{l} e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M^* + l_1 I^{1,\rho} \left(\rho g(s) V(s)\right)}{V(s)} \right] ds \ge -\left[C_7(L_1) + C_8(L) \right].$$

Since, the R.H.S of the above inequality is a negative constant, it follows that

$$\lim_{l \to \infty} \inf l^{1-n} \int_{L}^{l} e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M^* + l_1 I^{1,\rho} \left(\rho g(s) V(s) \right)}{V(s)} \right] ds > -\infty,$$

which contradicts the (36).

Therefore, y(l) is oscillatory. If y(l) is eventually negative, by a similar argument, we get a contradiction with condition (37). Hence the theorem. \Box

Remark 1. *If we put* $\rho = 1$ *in Theorem* (1) *and Theorem* (2), *then they reduced to Theorem* 3.1 *and Theorem* 4.1, *respectively, of* [24].

5. Examples

This section include some examples for the illustration of our main results.

Example 1. Consider the following GPF initial value problem

$$\begin{cases}
 {0}D^{\frac{3}{2},1}y(l) - {}{0}D^{\frac{1}{2},1}y(l) + (l+7)^{2}(y+3)e^{\cos 2y} = e^{2l}\sin l, & l > 0, \\
 \lim_{l \to 0^{+}} {}_{0}I^{\frac{1}{2},1}y(l) = b_{1}.
\end{cases}$$
(47)

Setting $\alpha = \frac{1}{2}$, $\rho = 1$, a = 0, p(l) = -1, $q(l) = (l+7)^2$, $f(y) = (y+3)e^{\cos 2y}$, $g(l) = e^{2l} \sin l$ and $V(l) = e^{l_1 - l}$. The assumption (H) is satisfied if y(l) > 0. Then,

$$\begin{split} l_1 I^{1,\rho} \left(\rho g(s) V(s) \right) &= \frac{1}{\rho} \int_{l_1}^s e^{\frac{\rho - 1}{\rho} (s - \tau)} \rho g(\tau) V(\tau) d\tau \\ &= \int_{l_1}^s e^{l_1 + \tau} \sin \tau d\tau \\ &= \frac{e^{l_1 + s}}{2} \left(\sin s - \cos s \right) - \frac{e^{2l_1}}{2} \left(\sin l_1 - \cos l_1 \right) \\ &= \frac{\sqrt{2} e^{l_1 + s}}{2} \sin \left(s - \frac{\pi}{4} \right) - \frac{e^{2l_1}}{2} \left(\sin l_1 - \cos l_1 \right). \end{split}$$

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Set a point $l_1 = \frac{\pi}{2}$. Hence, we compute that

$$\begin{split} l^{1-\alpha} & \int_0^l e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M + {}_{l_1} I^{1,\rho} \left(\rho g(s) V(s)\right)}{V(s)} \right] ds \\ = & l^{\frac{1}{2}} \int_0^l (l-s)^{-\frac{1}{2}} e^{s-\frac{\pi}{2}} \left[\left(M - \frac{e^{\pi}}{2} \right) + \frac{\sqrt{2}}{2} e^{\frac{\pi}{2} + s} \sin\left(s - \frac{\pi}{4}\right) \right] ds. \end{split}$$

By setting $l - s = \tau^2$, we can get the above integral as

$$\begin{split} & l^{\frac{1}{2}} \int_{0}^{l} (l-s)^{-\frac{1}{2}} e^{s-\frac{\pi}{2}} \left[\left(M - \frac{e^{\pi}}{2} \right) + \frac{\sqrt{2}}{2} e^{\frac{\pi}{2} + s} \sin \left(s - \frac{\pi}{4} \right) \right] ds \\ & = l^{\frac{1}{2}} \int_{\sqrt{l}}^{0} \frac{1}{\tau} e^{l-\tau^{2} - \frac{\pi}{2}} \left[\left(\frac{2M - e^{\pi}}{2} \right) + \frac{\sqrt{2}}{2} e^{\frac{\pi}{2} + l - \tau^{2}} \sin \left(l - \tau^{2} - \frac{\pi}{4} \right) \right] (-2\tau) d\tau \\ & = (2M - e^{\pi}) l^{\frac{1}{2}} e^{l-\frac{\pi}{2}} \int_{0}^{\sqrt{l}} e^{-\tau^{2}} d\tau + \sqrt{2} l^{\frac{1}{2}} e^{2l} \int_{0}^{\sqrt{l}} e^{-2\tau^{2}} \sin \left(l - \tau^{2} - \frac{\pi}{4} \right) d\tau \\ & = (2M - e^{\pi}) l^{\frac{1}{2}} e^{l-\frac{\pi}{2}} \int_{0}^{\sqrt{l}} e^{-\tau^{2}} d\tau + \sqrt{2} l^{\frac{1}{2}} e^{2l} \sin \left(l - \frac{\pi}{4} \right) \int_{0}^{\sqrt{l}} e^{-2\tau^{2}} \cos \tau^{2} d\tau \\ & - \sqrt{2} l^{\frac{1}{2}} e^{2l} \cos \left(l - \frac{\pi}{4} \right) \int_{0}^{\sqrt{l}} e^{-2\tau^{2}} \sin \tau^{2} d\tau. \end{split}$$

Let $l \to +\infty$ as the result of $|e^{-2\tau^2}\cos\tau^2| \le e^{-2\tau^2}$, $|e^{-2\tau^2}\sin\tau^2| \le e^{-2\tau^2}$ and $\lim_{l\to +\infty} \int_0^{\sqrt{l}} e^{-2\tau^2} d\tau = \frac{\sqrt{2\pi}}{4}$. Thus, we know that $\lim_{l\to +\infty} \int_0^{\sqrt{l}} e^{-2\tau^2} \cos\tau^2 d\tau$ and $\lim_{l\to +\infty} \int_0^{\sqrt{l}} e^{-2\tau^2} \sin\tau^2 d\tau$ are convergent.

Hence, we set $\lim_{l\to +\infty}\int_0^{\sqrt{l}}e^{-2\tau^2}\cos\tau^2d\tau=A$ and $\lim_{l\to +\infty}\int_0^{\sqrt{l}}e^{-2\tau^2}\sin\tau^2d\tau=B$. Select the sequence $\{l_k\}=\Big\{\frac{3\pi}{2}+\frac{\pi}{4}+2k\pi-\arctan\Big(-\frac{B}{A}\Big)\Big\}$, $\lim_{l\to +\infty}l_k=\infty$, then

$$\lim_{k \to +\infty} \left\{ l_k^{\frac{1}{2}} e^{l_k} \left[(2M - e^{\pi}) e^{-\frac{\pi}{2}} \int_0^{\sqrt{l_k}} e^{-\tau^2} d\tau + \sqrt{2} e^{l_k} \left(\sin \left(l_k - \frac{\pi}{4} \right) \int_0^{\sqrt{l_k}} e^{-2\tau^2} \cos \tau^2 d\tau \right) \right] \right\}.$$

$$\left. -\cos \left(l_k - \frac{\pi}{4} \right) \int_0^{\sqrt{l_k}} e^{-2\tau^2} \sin \tau^2 d\tau \right] \right\}.$$
(48)

Firstly, we consider the following limit:

$$\lim_{k \to +\infty} \left\{ \sin \left(l_k - \frac{\pi}{4} \right) \int_0^{\sqrt{l_k}} e^{-2\tau^2} \cos \tau^2 d\tau - \cos \left(l_k - \frac{\pi}{4} \right) \int_0^{\sqrt{l_k}} e^{-2\tau^2} \sin \tau^2 d\tau \right\}$$

$$= A \cdot \lim_{k \to +\infty} \sin \left(\frac{3\pi}{2} + 2k\pi - \arctan \left(-\frac{B}{A} \right) \right) - B \cdot \lim_{k \to +\infty} \cos \left(\frac{3\pi}{2} + 2k\pi - \arctan \left(-\frac{B}{A} \right) \right)$$

$$= A \cdot \sin \left(\frac{3\pi}{2} - \arctan \left(-\frac{B}{A} \right) \right) - B \cdot \cos \left(\frac{3\pi}{2} - \arctan \left(-\frac{B}{A} \right) \right)$$

$$= -\sqrt{A^2 + B^2}.$$

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Secondly, we know that $\lim_{k\to+\infty}e^{l_k}=+\infty$ and $\lim_{k\to+\infty}2Me^{l_k}e^{-\frac{\pi}{2}}\int_0^{\sqrt{l_k}}e^{-\tau^2}d\tau=2Me^{-\frac{\pi}{2}}\frac{\sqrt{\pi}}{2}=\sqrt{\pi}Me^{-\frac{\pi}{2}}$. Hence, for (48), we have

$$\begin{split} &\lim_{k \to +\infty} \left\{ l_k^{\frac{1}{2}} e^{l_k} \Bigg[\left(2M - e^{\pi} \right) e^{-\frac{\pi}{2}} \int_0^{\sqrt{l_k}} e^{-\tau^2} d\tau + \sqrt{2} e^{l_k} \Bigg(\sin \left(l_k - \frac{\pi}{4} \right) \int_0^{\sqrt{l_k}} e^{-2\tau^2} \cos \tau^2 d\tau \\ &- \cos \left(l_k - \frac{\pi}{4} \right) \int_0^{\sqrt{l_k}} e^{-2\tau^2} \sin \tau^2 d\tau \Bigg) \Bigg] \right\} \\ &= \left[\sqrt{\pi} M e^{-\frac{\pi}{2}} + (+\infty) \left(-\sqrt{A^2 + B^2} \right) \right] \\ &= -\infty. \end{split}$$

Then, we obtain

$$\lim_{l \to \infty} \inf l^{1-\alpha} \int_0^l e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M + {}_{l_1} I^{1,\rho} \left(\rho g(s) V(s) \right)}{V(s)} \right] ds = -\infty < 0.$$

Similarly, selecting the sequence $\{l_r\} = \left\{\frac{3\pi}{2} + \frac{\pi}{4} + 2r\pi - \arctan\left(-\frac{B}{A}\right)\right\}$, we can obtain

$$\lim_{l \to \infty} \sup l^{1-\alpha} \int_0^l e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M + l_1 I^{1,\rho} \left(\rho g(s) V(s) \right)}{V(s)} \right] ds = +\infty > 0.$$

Therefore, by Theorem 1 all solutions of the problem (47) are oscillatory.

Example 2. Consider the following GPF Caputo initial value problem

$$\begin{cases}
{}_{0}^{C}D^{\frac{3}{2},1}y(l) - {}_{0}^{C}D^{\frac{1}{2},1}y(l) + e^{(l+1)^{2}}\ln(y^{2} + e) = e^{2l}\cos l, & l > 0, \\
y(0) = b_{0}.
\end{cases} (49)$$

Setting $\alpha = \frac{1}{2}$, $\rho = 1$, a = 0, p(l) = -1, $q(l) = e^{(l+1)^2}$, $f(y) = \ln(y^2 + e)$, $g(l) = e^{2l} \cos l$ and $V(l) = e^{l_1 - l}$. The assumption (H) is satisfied if y(l) > 0. Then, we get

$$\begin{split} l_1 I^{1,\rho} \left(\rho g(s) V(s) \right) &= \frac{1}{\rho} \int_{l_1}^s e^{\frac{\rho - 1}{\rho} (s - \tau)} \rho g(\tau) V(\tau) d\tau \\ &= \int_{l_1}^s e^{l_1 + \tau} \cos \tau d\tau \\ &= \frac{e^{l_1 + s}}{2} \left(\sin s + \cos s \right) - \frac{e^{2l_1}}{2} \left(\sin l_1 + \cos l_1 \right) \\ &= \frac{\sqrt{2} e^{l_1 + s}}{2} \sin \left(s + \frac{\pi}{4} \right) - \frac{e^{2l_1}}{2} \left(\sin l_1 + \cos l_1 \right) . \end{split}$$

Set $l_1 = \frac{\pi}{2}$ with n = 1. Hence, we can compute that

$$\begin{split} l^{1-n} \int_0^l e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M^* + l_1 I^{1,\rho} \left(\rho g(s) V(s)\right)}{V(s)} \right] ds \\ = & \int_0^l (l-s)^{-\frac{1}{2}} e^{s-\frac{\pi}{2}} \left[\left(M^* - \frac{e^{\pi}}{2} \right) + \frac{\sqrt{2}}{2} e^{\frac{\pi}{2} + s} \sin\left(s + \frac{\pi}{4}\right) \right] ds. \end{split}$$

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By setting $l - s = \tau^2$, we can get the above integral as

$$\begin{split} & \int_0^l (l-s)^{-\frac{1}{2}} e^{s-\frac{\pi}{2}} \left[\left(M^* - \frac{e^{\pi}}{2} \right) + \frac{\sqrt{2}}{2} e^{\frac{\pi}{2} + s} \sin\left(s + \frac{\pi}{4}\right) \right] ds \\ &= \int_{\sqrt{l}}^0 \frac{1}{\tau} e^{l-\tau^2 - \frac{\pi}{2}} \left[\left(\frac{2M^* - e^{\pi}}{2} \right) + \frac{\sqrt{2}}{2} e^{\frac{\pi}{2} + l - \tau^2} \sin\left(l - \tau^2 + \frac{\pi}{4}\right) \right] (-2\tau) d\tau \\ &= \left(2M^* - e^{\pi} \right) e^{l - \frac{\pi}{2}} \int_0^{\sqrt{l}} e^{-\tau^2} d\tau + \sqrt{2} e^{2l} \int_0^{\sqrt{l}} e^{-2\tau^2} \sin\left(l - \tau^2 + \frac{\pi}{4}\right) d\tau \\ &= \left(2M^* - e^{\pi} \right) e^{l - \frac{\pi}{2}} \int_0^{\sqrt{l}} e^{-\tau^2} d\tau + \sqrt{2} e^{2l} \sin\left(l + \frac{\pi}{4}\right) \int_0^{\sqrt{l}} e^{-2\tau^2} \cos\tau^2 d\tau \\ &- \sqrt{2} e^{2l} \cos\left(l + \frac{\pi}{4}\right) \int_0^{\sqrt{l}} e^{-2\tau^2} \sin\tau^2 d\tau . \end{split}$$

Let $l \to +\infty$ as the result of $|e^{-2\tau^2}\cos\tau^2| \le e^{-2\tau^2}$, $|e^{-2\tau^2}\sin\tau^2| \le e^{-2\tau^2}$ and $\lim_{l\to +\infty}\int_0^{\sqrt{l}}e^{-2\tau^2}d\tau = \frac{\sqrt{2\pi}}{4}$. Thus, we know that $\lim_{l\to +\infty}\int_0^{\sqrt{l}}e^{-2\tau^2}\cos\tau^2d\tau$ and $\lim_{l\to +\infty}\int_0^{\sqrt{l}}e^{-2\tau^2}\sin\tau^2d\tau$ are convergent.

Hence, we can set $\lim_{l\to +\infty}\int_0^{\sqrt{l}}e^{-2\tau^2}\cos\tau^2d\tau=A$ and $\lim_{l\to +\infty}\int_0^{\sqrt{l}}e^{-2\tau^2}\sin\tau^2d\tau=B$. Select the sequence $\{l_k\}=\left\{\frac{7\pi}{2}-\frac{\pi}{4}+2k\pi-\arctan\left(-\frac{B}{A}\right)\right\}$, $\lim_{l\to +\infty}l_k=\infty$, then we compute the following term:

$$\lim_{k \to +\infty} \left\{ e^{l_k} \left[(2M^* - e^{\pi}) e^{-\frac{\pi}{2}} \int_0^{\sqrt{l_k}} e^{-\tau^2} d\tau + \sqrt{2} e^{l_k} \left(\sin \left(l_k + \frac{\pi}{4} \right) \int_0^{\sqrt{l_k}} e^{-2\tau^2} \cos \tau^2 d\tau \right) \right] \right\}.$$
(50)

Firstly, we consider the following limit:

$$\begin{split} &\lim_{k \to +\infty} \left\{ \sin \left(l_k + \frac{\pi}{4} \right) \int_0^{\sqrt{l_k}} e^{-2\tau^2} \cos \tau^2 d\tau - \cos \left(l_k + \frac{\pi}{4} \right) \int_0^{\sqrt{l_k}} e^{-2\tau^2} \sin \tau^2 d\tau \right\} \\ &= A \cdot \lim_{k \to +\infty} \sin \left(\frac{7\pi}{2} + 2k\pi - \arctan \left(-\frac{B}{A} \right) \right) - B \cdot \lim_{k \to +\infty} \cos \left(\frac{7\pi}{2} + 2k\pi - \arctan \left(-\frac{B}{A} \right) \right) \\ &= A \cdot \sin \left(\frac{7\pi}{2} - \arctan \left(-\frac{B}{A} \right) \right) - B \cdot \cos \left(\frac{7\pi}{2} - \arctan \left(-\frac{B}{A} \right) \right) \\ &= -\sqrt{A^2 + B^2}. \end{split}$$

Secondly, we know that $\lim_{k\to+\infty}e^{l_k}=+\infty$ and $\lim_{k\to+\infty}2M^*e^{l_k}e^{-\frac{\pi}{2}}\int_0^{\sqrt{l_k}}e^{-\tau^2}d\tau=2M^*e^{-\frac{\pi}{2}}\frac{\sqrt{\pi}}{2}=\sqrt{\pi}M^*e^{-\frac{\pi}{2}}$. Hence, for (50), we have

$$\lim_{k \to +\infty} \left\{ e^{l_k} \left[(2M^* - e^{\pi}) e^{-\frac{\pi}{2}} \int_0^{\sqrt{l_k}} e^{-\tau^2} d\tau + \sqrt{2} e^{l_k} \left(\sin \left(l_k + \frac{\pi}{4} \right) \int_0^{\sqrt{l_k}} e^{-2\tau^2} \cos \tau^2 d\tau \right) \right] \right\}$$

$$-\cos \left(l_k + \frac{\pi}{4} \right) \int_0^{\sqrt{l_k}} e^{-2\tau^2} \sin \tau^2 d\tau \right]$$

$$= \left[\sqrt{\pi} M^* e^{-\frac{\pi}{2}} + (+\infty) \left(-\sqrt{A^2 + B^2} \right) \right]$$

$$= -\infty.$$

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Then, we obtain

$$\lim_{l \to \infty} \inf l^{1-n} \int_0^l e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M^* + l_1 I^{1,\rho} \left(\rho g(s) V(s) \right)}{V(s)} \right] ds = -\infty < 0.$$

Similarly, selecting the sequence $\{l_r\} = \left\{\frac{\pi}{2} - \frac{\pi}{4} + 2r\pi - \arctan\left(-\frac{B}{A}\right)\right\}$, we can obtain

$$\lim_{l \to \infty} \sup l^{1-n} \int_0^l e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M^* + l_1 I^{1,\rho} \left(\rho g(s) V(s) \right)}{V(s)} \right] ds = +\infty > 0.$$

Therefore, by Theorem 2 all solutions of the problem (49) are oscillatory.

Example 3. Consider the following GPF Riemann-Liouville initial value problem

$$\begin{cases}
 _{0}D^{\frac{3}{2},1}y(l) + \sqrt{l}\left(\frac{4}{\sqrt{\pi}} + \frac{e^{3\sqrt{-y}}}{(-y)^{\frac{1}{4}}}\right) = e^{3l}, & l > 0, \\
 \lim_{l \to 0^{+}} {}_{0}I^{\frac{1}{2},1}y(l) = 0.
\end{cases}$$
(51)

Setting $\alpha = \frac{1}{2}$, $\rho = 1$, a = p(l) = 0, $q(l) = \sqrt{l}$, $f(y) = \frac{4}{\sqrt{\pi}} + \frac{e^{3\sqrt{-y}}}{(-y)^{\frac{1}{4}}}$, $g(l) = e^{3l}$ and V(l) = 1. The assumption (H) is satisfied if y(l) > 0. Then,

$$_{l_{1}}I^{1,\rho}\left(\rho g(s)V(s)\right)=\frac{1}{\rho}\int_{l_{1}}^{s}e^{\frac{\rho-1}{\rho}(s-\tau)}\rho g(\tau)V(\tau)d\tau=\int_{l_{1}}^{s}e^{3\tau}d\tau=\frac{e^{3s}}{3}-\frac{e^{3l_{1}}}{3}=\frac{1}{3}\left(e^{3s}-e^{3l_{1}}\right).$$

By setting $l_1 = \frac{1}{3}$ and $l - s = \tau^2$, it follows that

$$\begin{split} l^{1-\alpha} & \int_0^l e^{\frac{\rho-1}{\rho}(l-s)} (l-s)^{\alpha-1} \left[\frac{e^{\frac{\rho-1}{\rho}(s-l_1)} M + {}_{l_1} I^{1,\rho} \left(\rho g(s) V(s) \right)}{V(s)} \right] ds \\ = & l^{\frac{1}{2}} \int_0^l (l-s)^{-\frac{1}{2}} \left[M + \frac{1}{3} \left(e^{3s} - e \right) \right] ds. \\ = & l^{\frac{1}{2}} \left[\left(M - \frac{e}{3} \right) \int_0^l (l-s)^{-\frac{1}{2}} ds + \frac{1}{3} \int_0^l (l-s)^{-\frac{1}{2}} e^{3s} ds \right] \\ = & 2l^{\frac{1}{2}} \left[\left(M - \frac{e}{3} \right) \sqrt{l} + \frac{1}{3} e^{3l} \int_0^{\sqrt{l}} e^{-3\tau} d\tau \right]. \end{split}$$

However, the condition (22) does not holds since

$$\begin{split} &\lim_{l\to\infty}\inf l^{1-\alpha}\int_0^l e^{\frac{\rho-1}{\rho}(l-s)}(l-s)^{\alpha-1}\left[\frac{e^{\frac{\rho-1}{\rho}(s-a)}\;M+_{l_1}I^{1,\rho}\left(\rho g(s)V(s)\right)}{V(s)}\right]ds\\ &=&\lim_{l\to\infty}\inf\left\{2l^{\frac{1}{2}}\left[\left(M-\frac{e}{3}\right)\sqrt{l}+\frac{1}{3}e^{3l}\int_0^{\sqrt{l}}e^{-3\tau}d\tau\right]\right\}\\ &=&\left[\left(M-\frac{e}{3}\right)(+\infty)+(+\infty)\frac{\sqrt{\pi}}{2}\right]=\infty. \end{split}$$

Using Proposition 1 (ii) with $\alpha = \frac{3}{2}$, $\beta = 3$ and $\rho = 1$, we get ${}_aD^{\alpha,\rho}y(l) = -\frac{4\sqrt{l}}{\sqrt{\pi}}$, it is easy to verify that $y(l) = -l^2$ is a nonoscillatory solution of (51). Figure 1 demonstrates the solution $y(l) = -l^2$.

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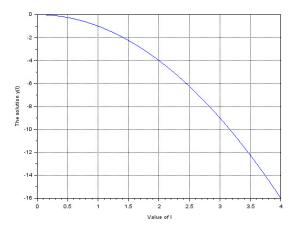


Figure 1. The nonoscillatory behavior of the solution $y(l) = -l^2$.

6. Conclusions

In this paper, the oscillatory behavior of solutions of generalized proportional fractional initial value problem is studied. Forced and damped oscillation results are obtained via GPF operators in the frame of Riemann-Liouville and Caputo settings. The main theorems of this paper improve and generalize some existing oscillation theorems reported in the literature. In particular, for the choice of $\rho=1$, our contributions obtained using GPF operators cover the results discussed in Reference [24] which are obtained via conformable operators. At the end, we presented some numerical examples with particular values of parameters to illustrate the validity of the proposed results. Interestingly, we provided an example demonstrating that the failure of any condition forces the existence of a nonoscillatory solution. This justifies the advantage of our findings.

We believe that the results of this paper are of great importance for the audience of interested researchers. Several types of oscillation conditions could be generalized by considering respective equations within GPF derivatives.

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