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Hankel and Toeplitz Determinants for a Subclass of *q*-Starlike Functions Associated with a General Conic Domain

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Received: 29 January 2019; Accepted: 12 February 2019; Published: 15 February 2019



Abstract: By using a certain general conic domain as well as the quantum (or q-) calculus, here we define and investigate a new subclass of normalized analytic and starlike functions in the open unit disk \mathbb{U} . In particular, we find the Hankel determinant and the Toeplitz matrices for this newly-defined class of analytic q-starlike functions. We also highlight some known consequences of our main results.

Keywords: analytic functions; starlike and *q*-starlike functions; *q*-derivative operator; *q*-hypergeometric functions; conic and generalized conic domains; Hankel determinant; Toeplitz matrices

MSC: Primary 05A30, 30C45; Secondary 11B65, 47B38

1. Introduction and Definitions

Let the class of functions, which are analytic in the open unit disk

$$\mathbb{U} = \{ z : z \in \mathbb{C} \text{ and } |z| < 1 \}$$
 ,

be denoted by $\mathcal{L}(\mathbb{U})$. Also let \mathcal{A} denote the class of all functions f, which are analytic in the open unit disk \mathbb{U} and normalized by

$$f(0) = 0$$
 and $f'(0) = 1$.

Then, clearly, each $f \in A$ has a Taylor–Maclaurin series representation as follows:

$$f(z) = z + \sum_{n=2}^{\infty} a_n z^n \qquad (z \in \mathbb{U}).$$
⁽¹⁾

Suppose that S is the subclass of the analytic function class A, which consists of all functions which are also univalent in \mathbb{U} .

A function $f \in A$ is said to be starlike in \mathbb{U} if it satisfies the following inequality:

$$\Re\left(\frac{zf'(z)}{f(z)}\right) > 0 \qquad (z \in \mathbb{U}).$$



We denote by S^* the class of all such starlike functions in \mathbb{U} .

For two functions f and g, analytic in \mathbb{U} , we say that the function f is subordinate to the function g and write this subordination as follows:

$$f \prec g$$
 or $f(z) \prec g(z)$.

if there exists a Schwarz function w which is analytic in \mathbb{U} , with

$$w(0) = 0$$
 and $|w(z)| < 1$,

such that

$$f(z) = g(w(z)).$$

In the case when the function g is univalent in \mathbb{U} , then we have the following equivalence (see, for example, [1]; see also [2]):

$$f(z) \prec g(z) \quad (z \in \mathbb{U}) \iff f(0) = g(0) \text{ and } f(\mathbb{U}) \subset g(\mathbb{U})$$

Next, for a function $f \in A$ given by (1) and another function $g \in A$ given by

$$g(z) = z + \sum_{n=2}^{\infty} b_n z^n \qquad (z \in \mathbb{U}),$$

the convolution (or the Hadamard product) of *f* and *g* is defined here by

$$(f * g)(z) := z + \sum_{n=2}^{\infty} a_n b_n z^n =: (g * f)(z).$$
⁽²⁾

Let \mathcal{P} denote the well-known Carathéodory class of functions p, analytic in the open unit disk \mathbb{U} , which are normalized by

$$p(z) = 1 + \sum_{n=1}^{\infty} c_n z^n,$$
 (3)

such that

$$\Re(p(z)) > 0 \qquad (z \in \mathbb{U})$$

Following the works of Kanas et al. (see [3,4]; see also [5]), we introduce the conic domain Ω_k ($k \ge 0$) as follows:

$$\Omega_k = \left\{ u + iv : u > k\sqrt{(u-1)^2 + v^2} \right\}.$$
(4)

In fact, subjected to the conic domain Ω_k ($k \ge 0$), Kanas and Wiśniowska (see [3,4]; see also [6]) studied the corresponding class k-ST of k-starlike functions in \mathbb{U} (see Definition 1 below). For fixed k, Ω_k represents the conic region bounded successively by the imaginary axis (k = 0), by a parabola (k = 1), by the right branch of a hyperbola (0 < k < 1), and by an ellipse (k > 1).

For these conic regions, the following functions play the role of extremal functions.

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$$\begin{cases} \frac{1+z}{1-z} = 1 + 2z + 2z^2 + \cdots & (k=0) \\ 1 + \frac{2}{\pi^2} \left[\log\left(\frac{1+\sqrt{z}}{1-\sqrt{z}}\right) \right]^2 & (k=1) \end{cases}$$

$$p_{k}(z) = \begin{cases} 1 + \frac{2}{1-k^{2}} \sinh^{2} \left[\left(\frac{2}{\pi} \arccos k \right) \arctan \left(h \sqrt{z} \right) \right] & (0 \leq k < 1) \\ 1 + \frac{1}{k^{2} - 1} \left[1 + \sin \left(\frac{\pi}{2K(\kappa)} \int_{0}^{\frac{u(z)}{\sqrt{\kappa}}} \frac{dt}{\sqrt{(1-t^{2})(1-\kappa^{2}t^{2})}} \right) \right] & (k > 1), \end{cases}$$
(5)

where

$$u(z) = rac{z - \sqrt{\kappa}}{1 - \sqrt{\kappa}z}$$
 $(z \in \mathbb{U})$,

and $\kappa \in (0, 1)$ is so chosen that

$$k = \cosh\left(\frac{\pi K'(\kappa)}{4K(\kappa)}\right).$$

Here $K(\kappa)$ is Legendre's complete elliptic integral of first kind and

$$K'(\kappa) = K\left(\sqrt{1-\kappa^2}\right),$$

that is, $K'(\kappa)$ is the complementary integral of $K(\kappa)$ (see, for example, ([7], p. 326, Equation 9.4 (209))). Indeed, from (5), we have

$$p_k(z) = 1 + p_1 z + p_2 z^2 + p_3 z^3 + \cdots$$
 (6)

The class k-ST is defined as follows.

Definition 1. A function $f \in A$ is said to be in the class k-ST if and only if

$$\frac{zf'\left(z\right)}{f\left(z\right)} \prec p_k\left(z\right) \quad \left(\forall \ z \in \mathbb{U}; \ k \geqq 0\right).$$

We now recall some basic definitions and concept details of the *q*-calculus which will be used in this paper (see, for example, ([7], p. 346 et seq.)). Throughout the paper, unless otherwise mentioned, we suppose that 0 < q < 1 and

$$\mathbb{N}=\{1,2,3\cdots\}=\mathbb{N}_0\setminus\{0\}\qquad (\mathbb{N}_0:=\{0,1,2,\cdots\})\,.$$

Definition 2. *Let* $q \in (0, 1)$ *and define the* q*-number* $[\lambda]_q$ *by*

$$[\lambda]_q = \begin{cases} \frac{1-q^{\lambda}}{1-q} & (\lambda \in \mathbb{C}) \\\\ \sum_{k=0}^{n-1} q^k = 1+q+q^2+\dots+q^{n-1} & (\lambda = n \in \mathbb{N}). \end{cases}$$

Definition 3. Let $q \in (0, 1)$ and define the *q*-factorial $[n]_q!$ by

$$[n]_q! = \begin{cases} 1 & (n=0) \\ \\ \prod_{k=1}^n [k]_q & (n \in \mathbb{N}) \,. \end{cases}$$

Definition 4 (see [8,9]). The q-derivative (or q-difference) operator D_q of a function f defined, in a given subset of \mathbb{C} , by

$$(D_q f) (z) = \begin{cases} \frac{f(z) - f(qz)}{(1-q)z} & (z \neq 0) \\ f'(0) & (z = 0), \end{cases}$$

$$(7)$$

provided that f'(0) exists.

From Definition 4, we can observe that

$$\lim_{q \to 1^{-}} (D_q f)(z) = \lim_{q \to 1^{-}} \frac{f(z) - f(qz)}{(1 - q)z} = f'(z)$$

for a differentiable function f in a given subset of \mathbb{C} . It is also known from (1) and (7) that

$$(D_q f)(z) = 1 + \sum_{n=2}^{\infty} [n]_q a_n z^{n-1}.$$
 (8)

Definition 5. The *q*-Pochhammer symbol $[\xi]_{n,q}$ ($\xi \in \mathbb{C}$; $n \in \mathbb{N}_0$) is defined as follows:

$$[\xi]_{n,q} = \frac{(q^{\xi};q)_n}{(1-q)^n} = \begin{cases} 1 & (n=0) \\ [\xi]_q [\xi+1]_q [\xi+2]_q \cdots [\xi+n-1]_q & (n \in \mathbb{N}) . \end{cases}$$

Moreover, the q-gamma function is defined by the following recurrence relation:

$$\Gamma_q(z+1) = [z]_q \Gamma_q(z)$$
 and $\Gamma_q(1) = 1$.

Definition 6 (see [10]). For $f \in A$, let the *q*-Ruscheweyh derivative operator \mathcal{R}_q^{λ} be defined, in terms of the Hadamard product (or convolution) given by (2), as follows:

$$\mathcal{R}_{q}^{\lambda}f(z) = f(z) * \mathcal{F}_{q,\lambda+1}(z) \qquad (z \in \mathbb{U}; \lambda > -1),$$

where

$$\mathcal{F}_{q,\lambda+1}\left(z\right) = z + \sum_{n=2}^{\infty} \frac{\Gamma_q\left(\lambda+n\right)}{[n-1]_q! \Gamma_q\left(\lambda+1\right)} z^n = z + \sum_{n=2}^{\infty} \frac{[\lambda+1]_{q,n-1}}{[n-1]_q!} z^n.$$

We next define a certain *q*-integral operator by using the same technique as that used by Noor [11].

Definition 7. For $f \in A$, let the *q*-integral operator $\mathcal{F}_{q,\lambda}$ be defined by

$$\mathcal{F}_{q,\lambda+1}^{-1}(z) * \mathcal{F}_{q,\lambda+1}(z) = z \left(D_q f \right)(z).$$

Then

$$\mathcal{I}_{q}^{\lambda}f(z) = f(z) * \mathcal{F}_{q,\lambda+1}^{-1}(z)$$

= $z + \sum_{n=2}^{\infty} \psi_{n-1}a_{n}z^{n}$ $(z \in \mathbb{U}; \lambda > -1),$ (9)

where

$$\mathcal{F}_{q,\lambda+1}^{-1}(z) = z + \sum_{n=2}^{\infty} \psi_{n-1} z^n$$

and

$$\psi_{n-1} = \frac{[n]_q!\Gamma_q(\lambda+1)}{\Gamma_q(\lambda+n)} = \frac{[n]_q!}{[\lambda+1]_{q,n-1}}.$$

Clearly, we have

$$\mathcal{I}_{q}^{0}f\left(z\right)=z\left(D_{q}f\right)\left(z\right) \quad \text{and} \quad \mathcal{I}_{q}^{1}f\left(z\right)=f\left(z\right).$$

We note also that, in the limit case when $q \to 1-$, the *q*-integral operator $\mathcal{F}_{q,\lambda}$ given by Definition 7 would reduce to the integral operator which was studied by Noor [11].

The following identity can be easily verified:

$$zD_q\left(\mathcal{I}_q^{\lambda+1}f(z)\right) = \left(1 + \frac{[\lambda]_q}{q^{\lambda}}\right)\mathcal{I}_q^{\lambda}f(z) - \frac{[\lambda]_q}{q^{\lambda}}\mathcal{I}_q^{\lambda+1}f(z).$$
(10)

When $q \rightarrow 1-$, this last identity in (10) implies that

$$z\left(\mathcal{I}^{\lambda+1}f(z)\right)' = (1+\lambda)\mathcal{I}^{\lambda}f(z) - \lambda\mathcal{I}^{\lambda+1}f(z),$$

which is the well-known recurrence relation for the above-mentioned integral operator which was studied by Noor [11].

In geometric function theory, several subclasses belonging to the class of normalized analytic functions class \mathcal{A} have already been investigated in different aspects. The above-defined *q*-calculus gives valuable tools that have been extensively used in order to investigate several subclasses of \mathcal{A} . Ismail et al. [12] were the first who used the *q*-derivative operator D_q to study the *q*-calculus analogous of the class \mathcal{S}^* of starlike functions in \mathbb{U} (see Definition 8 below). However, a firm footing of the *q*-calculus in the context of geometric function theory was presented mainly and basic (or *q*-) hypergeometric functions were first used in geometric function theory in a book chapter by Srivastava (see, for details, ([13], p. 347 et seq.); see also [14]).

Definition 8 (see [12]). A function $f \in A$ is said to belong to the class S_q^* if

$$f(0) = f'(0) - 1 = 0 \tag{11}$$

and

$$\left|\frac{z}{f(z)}\left(D_{q}f\right)z - \frac{1}{1-q}\right| \leq \frac{1}{1-q}.$$
(12)

It is readily observed that, as $q \rightarrow 1-$, the closed disk:

$$\left|w - \frac{1}{1-q}\right| \le \frac{1}{1-q}$$

becomes the right-half plane and the class S_q^* of *q*-starlike functions reduces to the familiar class S^* of normalized starlike functions in \mathbb{U} with respect to the origin (z = 0). Equivalently, by using the principle of subordination between analytic functions, we can rewrite the conditions in (11) and (12) as follows (see [15]):

$$\frac{z}{f(z)} \left(D_q f \right)(z) \prec \hat{p}(z) \qquad \left(\hat{p}(z) = \frac{1+z}{1-qz} \right).$$
(13)

The notation S_q^* was used by Sahoo and Sharma [16].

Now, making use of the principle of subordination between analytic functions and the above-mentioned *q*-calculus, we present the following definition.

Definition 9. A function p is said to be in the class k- \mathcal{P}_q if and only if

$$p(z) \prec \frac{2p_k(z)}{(1+q)+(1-q)p_k(z)},$$

where $p_k(z)$ is defined by (5).

Geometrically, the function $p(z) \in k-\mathcal{P}_q$ takes on all values from the domain $\Omega_{k,q}$ $(k \ge 0)$ which is defined as follows:

$$\Omega_{k,q} = \left\{ w : \Re\left(\frac{(1+q)w}{(q-1)w+2}\right) > k \left| \frac{(1+q)w}{(q-1)w+2} - 1 \right| \right\}.$$

The domain $\Omega_{k,q}$ represents a generalized conic region. It can be seen that

$$\lim_{q\to 1-}\Omega_{k,q}=\Omega_k,$$

where Ω_k is the conic domain considered by Kanas and Wiśniowska [3]. Below, we give some basic facts about the class $k-\mathcal{P}_q$.

Remark 1. First of all, we see that

$$k-\mathcal{P}_q \subseteq \mathcal{P}\left[\frac{2k}{2k+1+q}\right],$$

where $\mathcal{P}\left[\frac{2k}{2k+1+q}\right]$ is the well-known class of functions with real part greater than $\frac{2k}{2k+1+q}$. Secondly, we have

$$\lim_{q\to 1-}k\mathcal{P}_q=\mathcal{P}\left(p_k\right),$$

where $\mathcal{P}(p_k)$ is the well-known function class introduced by Kanas and Wiśniowska [3]. Thirdly, we have

$$\lim_{q \to 1-} 0 - \mathcal{P}_q = \mathcal{P}_q$$

where \mathcal{P} is the well-known class of analytic functions with positive real part.

Definition 10. A function f is said to be in the class $ST(k, \lambda, q)$ if and only if

$$\frac{z\left(D_{q}\mathcal{I}_{q}^{\lambda}f\right)(z)}{f\left(z\right)}\in k\mathcal{P}_{q} \qquad (k\geqq 0;\;\lambda\geqq 0),$$

or, equivalently,

$$\Re\left(\frac{(1+q)\,\frac{z\left(D_{q}\mathcal{I}_{q}^{\lambda}f\right)(z)}{f(z)}}{(q-1)\,\frac{z\left(D_{q}\mathcal{I}_{q}^{\lambda}f\right)(z)}{f(z)}+2}\right) > k\left|\frac{(1+q)\,\frac{z\left(D_{q}\mathcal{I}_{q}^{\lambda}f\right)(z)}{f(z)}}{(q-1)\,\frac{z\left(D_{q}\mathcal{I}_{q}^{\lambda}f\right)(z)}{f(z)}+2}-1\right|$$

Remark 2. First of all, it is easily seen that

$$\mathcal{ST}\left(0,1,q\right)=\mathcal{S}_{q}^{*},$$

where S_q^* is the function class introduced and studied by Ismail et al. [12]. Secondly, we have

$$\lim_{q \to 1-} \mathcal{ST}(k, 1, q) = k - \mathcal{ST},$$

where k-ST is a function class introduced and studied by Kanas and Wiśniowska [4]. Finally, we have

$$\lim_{q \to 1-} \mathcal{ST}(0,1,q) = \mathcal{S}^*,$$

where S^* is the well-known class of starlike functions in \mathbb{U} with respect to the origin (z = 0).

Remark 3. Further studies of the new q-starlike function class $ST(k, \lambda, q)$, as well as of its more consequences, can next be determined and investigated in future papers.

Let $n \in \mathbb{N}_0$ and $j \in \mathbb{N}$. The following *j*th Hankel determinant was considered by Noonan and Thomas [17]:

$\mathcal{H}_{j}\left(n ight)=$	a _n	a_{n+1}			a_{n+j-1}	
	a_{n+1}	•			•	
	•					
	•					ľ
	•	•			•	
	a_{n+j-1}	•	•	•	$a_{n+2(j-1)}$	

where $a_1 = 1$. In fact, this determinant has been studied by several authors, and sharp upper bounds on $\mathcal{H}_2(2)$ were obtained by several authors (see [18–20]) for various classes of functions. It is well-known that the Fekete–Szegö functional $|a_3 - a_2^2|$ can be represented in terms of the Hankel determinant as $\mathcal{H}_2(1)$. This functional has been further generalized as $|a_3 - \mu a_2^2|$ for some real or complex μ . Fekete and Szegö gave sharp estimates of $|a_3 - \mu a_2^2|$ for μ real and $f \in S$, the class of normalized univalent functions in \mathbb{U} . It is also known that the functional $|a_2a_4 - a_3^2|$ is equivalent to $\mathcal{H}_2(2)$ (see [18]). Babalola [21] studied the Hankel determinant $\mathcal{H}_3(1)$ for some subclasses of normalized analytic functions in \mathbb{U} . The symmetric Toeplitz determinant $\mathcal{T}_i(n)$ is defined by

$$\mathcal{T}_{j}(n) = \begin{vmatrix} a_{n} & a_{n+1} & \dots & a_{n+j-1} \\ a_{n+1} & \ddots & & \ddots \\ \ddots & \ddots & & \ddots \\ \vdots & \ddots & \ddots & \ddots \\ a_{n+j-1} & \ddots & \ddots & \ddots & a_{n} \end{vmatrix} ,$$

so that

$$\mathcal{T}_{2}(2) = \begin{vmatrix} a_{2} & a_{3} \\ & & \\ a_{3} & a_{2} \end{vmatrix}, \qquad \mathcal{T}_{2}(3) = \begin{vmatrix} a_{3} & a_{4} \\ & & \\ a_{4} & a_{3} \end{vmatrix}, \qquad \mathcal{T}_{3}(2) = \begin{vmatrix} a_{2} & a_{3} & a_{4} \\ & & \\ a_{3} & a_{2} & a_{3} \\ & & \\ a_{4} & a_{3} & a_{2} \end{vmatrix}$$

and so on.

For $f \in S$, the problem of finding the best possible bounds for $||a_{n+1}| - |a_n||$ has a long history (see, for details, [22]). It is a known fact from [22] that

$$||a_{n+1}| - |a_n|| < c$$

for a constant *c*. However, the problem of finding exact values of the constant *c* for *S* and its various subclasses has proved to be difficult. In a very recent investigation, Thomas and Abdul-Halim [23] succeeded in obtaining some sharp estimates for $\mathcal{T}_j(n)$ for the first few values of *n* and *j* involving symmetric Toeplitz determinants whose entries are the coefficients a_n of starlike and close-to-convex functions.

In the present investigation, our focus is on the Hankel determinant and the Toeplitz matrices for the function class $ST(k, \lambda, q)$ given by Definition 10.

2. A Set of Lemmas

In order to prove our main results in this paper, we need each of the following lemmas.

Lemma 1 (see [20]). *If the function* p(z) *given by* (3) *is in the Carathéodory class* \mathcal{P} *of analytic functions with positive real part in* \mathbb{U} *, then*

$$2c_2 = c_1^2 + x\left(4 - c_1^2\right)$$

and

$$4c_{3} = c_{1}^{3} + 2\left(4 - c_{1}^{2}\right)c_{1}x - c_{1}\left(4 - c_{1}^{2}\right)x^{2} + 2\left(4 - c_{1}^{2}\right)\left(1 - \left|x^{2}\right|\right)z$$

for some $x, z \in \mathbb{C}$ *with* $|x| \leq 1$ *and* $|z| \leq 1$ *.*

Lemma 2 (see [24]). Let the function p(z) given by (3) be in the Carathéodory class \mathcal{P} of analytic functions with positive real part in \mathbb{U} . Also let $\mu \in \mathbb{C}$. Then

$$|c_n - \mu c_k c_{n-k}| \le 2 \max(1, |2\mu - 1|) \qquad (1 \le k \le n - 1).$$

Lemma 3 (see [22]). Let the function p(z) given by (3) be in the Carathéodory class \mathcal{P} of analytic functions with positive real part in \mathbb{U} . Then

$$|c_n| \leq 2 \qquad (n \in \mathbb{N}).$$

This last inequality is sharp.

3. Main Results

Throughout this section, unless otherwise mentioned, we suppose that

$$q \in (0,1)$$
, $\lambda > -1$ and $k \in [0,1]$.

Theorem 1. *If the function* f(z) *given by* (1) *belongs to the class* $ST(k, \lambda, q)$ *, where* $k \in [0, 1]$ *, then*

$$|a_{2}| \leq \frac{(1+q) p_{1}}{2q\psi_{1}},$$
$$a_{3} \leq \frac{1}{2q\psi_{2}} \left(p_{1} + \left| p_{2} - p_{1} + \frac{(q^{2}+1) p_{1}^{2}}{2q} \right| \right)$$

and

$$a_{4} \leq \frac{(1+q)}{4(q+q^{2}+q^{3})\psi_{3}} \left(2p_{1}+4\left|p_{2}-p_{1}+\frac{(2+q^{2})p_{1}^{2}}{4q}\right| + \left|2p_{3}+2p_{1}-4p_{2}-\frac{(2(1+q^{2})-q)p_{1}^{2}}{q}+\frac{(4q^{2}-3q+2)}{q}p_{1}p_{2} + \frac{(q^{2}+2q-1)}{2q^{2}}p_{1}^{3}\right|\right),$$
(14)

where p_j (j = 1, 2, 3) are positive and are the coefficients of the functions $p_k(z)$ defined by (6). Each of the above results is sharp for the function g(z) given by

$$g(z) = \frac{2p_k(z)}{(1+q) + (1-q) p_k(z)}.$$

Proof. Let $f(z) \in ST(k, \lambda, q)$. Then, we have

$$\frac{z\left(D_{q}f\right)\left(z\right)}{f\left(z\right)} = \mathfrak{q}\left(z\right) \prec S_{k}\left(z\right),\tag{15}$$

where

$$S_{k}(z) = \frac{2p_{k}(z)}{(1+q) + (1-q)p_{k}(z)},$$

and the functions $p_k(z)$ are defined by (6).

We now define the function p(z) with p(0) = 1 and with a positive real part in \mathbb{U} as follows:

$$p(z) = \frac{1 + S_k^{-1}(\mathfrak{q}(z))}{1 - S_k^{-1}(\mathfrak{q}(z))} = 1 + c_1 z + c_2 z^2 + \cdots .$$
(16)

After some simple computation involving (16), we get

$$\mathfrak{q}(z) = S_k\left(\frac{p(z)+1}{p(z)-1}\right).$$

We thus find that

$$S_{k}\left(\frac{p(z)+1}{p(z)-1}\right)$$

$$=1+\left(\frac{q+1}{2}\right)\left[\frac{p_{1}c_{1}}{2}z+\left\{\frac{p_{1}c_{2}}{2}+\left(\frac{p_{2}}{4}-\frac{p_{1}}{4}+\left(\frac{(q-1)p_{1}^{2}}{8}\right)\right)c_{1}^{2}\right\}z^{2}$$

$$+\left\{\frac{p_{1}c_{3}}{2}+\left(\frac{p_{2}}{2}-\frac{p_{1}}{2}+\left(\frac{(q-1)p_{1}^{2}}{4}\right)\right)c_{1}c_{2}$$

$$+\left(\frac{p_{1}}{8}-\frac{p_{2}}{4}-\frac{(q-1)p_{1}^{2}}{8}+\frac{p_{3}}{8}-\frac{(q-1)p_{1}p_{2}}{8}+\frac{(q-1)^{2}p_{1}^{3}}{32}\right)c_{1}^{3}\right\}z^{3}\right]+\cdots.$$
(17)

Now, upon expanding the left-hand side of (15), we have

$$\frac{z\left(D_{q}\mathcal{I}_{q}^{\lambda}f\right)(z)}{f(z)} = 1 + q\psi_{1}a_{2}z + \left\{\left(q+q^{2}\right)\psi_{2}a_{3} - q\psi_{1}^{2}a_{2}^{2}\right\}z^{2} + \left\{\left(q+q^{2}+q^{3}\right)\psi_{3}a_{4} - \left(2q+q^{2}\right)\psi_{1}\psi_{2}a_{2}a_{3} + q\psi_{1}^{3}a_{2}^{3}\right\}z^{3} + \cdots$$
(18)

Finally, by comparing the corresponding coefficients in (17) and (18) along with Lemma 3, we obtain the result asserted by Theorem 1. \Box

Theorem 2. *If the function f* (*z*) *given by* (1) *belongs to the class* $ST(k, \lambda, q)$ *, then*

$$\begin{aligned} \mathcal{T}_{3}\left(2\right) &\leq \left[\left(\frac{1+q}{2q\psi_{1}}\right) p_{1}^{2} + \left(\frac{1+q}{4\left(q+q^{2}+q^{3}\right)\psi_{3}}\right) \left[\Omega_{1}+\Omega_{2}\right] \right] \\ &\cdot \left[4 \left(\frac{(1+q)^{2}}{16q^{2}\psi_{1}^{2}}\right) p_{1}^{2} + 16\left|\Omega_{3}\right| + \frac{p_{1}^{2}}{4q^{2}\psi_{2}^{2}} + 2\Omega_{5}p_{1}^{2} \left|2 - \frac{\Omega_{4}}{\Omega_{5}p_{1}^{2}}\right| \right], \end{aligned}$$

where

$$\begin{split} \Omega_{1} &= 2p_{1} + 4 \left| p_{2} - p_{1} + \frac{(2+q^{2})}{4q} p_{1}^{2} \right|, \\ \Omega_{2} &= \left| 2p_{3} + 2p_{1} - 4p_{2} - \left(2\left(1+q^{2}\right) - q\right) p_{1}^{2} \right. \\ &+ \left(\frac{4q^{2} - 3q + 2}{q}\right) p_{1}p_{2} + \left(\frac{q^{2} + q + 1}{2q^{2}} p_{1}^{3}\right) \right|, \\ \Omega_{3} &= \frac{1}{2q^{2}\psi_{2}^{2}} \left(\frac{p_{2}}{4} - \frac{p_{1}}{4} + \frac{(q^{2} + 1) p_{1}^{2}}{8q} \right)^{2} - \Omega_{5} \cdot \left[\frac{p_{3}}{4} + \frac{p_{1}}{4} - \frac{p_{2}}{2} \right. \\ &- \frac{\left[2\left(1+q^{2}\right) - q\right] p_{1}^{2}}{8q} + \frac{4q^{2} - 3q + 2}{8q} p_{1}p_{2} + \left(\frac{q^{2} + 2q - 1}{16q^{2}}\right) p_{1}^{3} \right], \\ \Omega_{4} &= \frac{p_{1}}{2q^{2}\psi_{2}^{2}} \left(\frac{p_{2}}{4} - \frac{p_{1}}{4} + \frac{(q^{2} + 1) p_{1}^{2}}{8q} \right) - \Omega_{5}p_{1} \left(p_{2} - p_{1} + \frac{(2+q^{2}) p_{1}^{2}}{4q} \right), \\ \Omega_{5} &= \frac{(1+q)^{2}}{16q^{2}\left(1+q+q^{2}\right) \psi_{1}\psi_{3}} \end{split}$$

and p_{j} (j = 1, 2) are positive and are the coefficients of the functions $p_{k}(z)$ defined by (6).

Proof. Upon comparing the corresponding coefficients in (17) and (18), we find that

$$a_2 = \frac{(1+q)\,p_1c_1}{4q\psi_1},\tag{19}$$

$$a_3 = \frac{1}{2q\psi_2} \left[\frac{p_1c_2}{2} + \left(\frac{p_2}{4} - \frac{p_1}{4} + \frac{(q^2+1)p_1^2}{8q} \right) c_1^2 \right],$$
(20)

$$a_{4} = \frac{(1+q)}{4(q+q^{2}+q^{3})\psi_{3}} \left[p_{1}c_{3} + \left(p_{2} - p_{1} + \frac{(2+q^{2})p_{1}^{2}}{4q} \right) c_{1}c_{2} + \left(\frac{p_{3}}{4} + \frac{p_{1}}{4} - \frac{p_{2}}{2} - \frac{(2(1+q^{2})-q)p_{1}^{2}}{8q} + \frac{(4q^{2}-3q+2)}{8q}p_{1}p_{2} + \frac{(q^{2}+2q-1)}{16q^{2}}p_{1}^{3} \right)c_{1}^{3} \right].$$
(21)

By a simple computation, $\mathcal{T}_{3}\left(2\right)$ can be written as follows:

$$\mathcal{T}_{3}(2) = (a_{2} - a_{4}) \left(a_{2}^{2} - 2a_{3}^{2} + a_{2}a_{4}\right).$$

Now, if $f \in \mathcal{ST}(k, \lambda, q)$, then it is clearly seen that

$$\begin{aligned} |a_2 - a_4| &\leq |a_2| + |a_4| \\ &\leq \left(\frac{1+q}{2q\psi_1}\right) p_1^2 + \left(\frac{1+q}{4(q+q^2+q^3)\psi_3}\right) \left(\Omega_1 + \Omega_2\right). \end{aligned}$$

We need to maximize $|a_2^2 - 2a_3^2 + a_2a_4|$ for a function $f \in ST(k, \lambda, q)$. So, by writing a_2, a_3 , and a_4 in terms of c_1, c_2 , and c_3 , with the help of (19)–(21), we get

$$\begin{vmatrix} a_2^2 - 2a_3^2 + a_2 a_4 \end{vmatrix} = \left| \left(\frac{(1+q)^2}{16q^2 \psi_1^2} \right) p_1^2 c_1^2 - \Omega_3 c_1^4 - \Omega_4 c_1^2 c_2 - \frac{p_1^2}{8q^2 \psi_2^2} c_2^2 + \Omega_5 p_1^2 c_1 c_3 \end{vmatrix} .$$
(22)

Finally, by applying the trigonometric inequalities, Lemmas 2 and 3 along with (22), we obtain the result asserted by Theorem 2. \Box

As an application of Theorem 2, we first set $\psi_{n-1} = 1$ and k = 0 and then let $q \to 1-$. We thus arrive at the following known result.

Corollary 1 (see [25]). *If the function* f(z) *given by* (1) *belongs to the class* S^* *, then*

$$\mathcal{T}_{3}(2) \leq 84$$

Theorem 3. If the function f(z) given by (1) belongs to the class $ST(k, \lambda, q)$, then

$$\left|a_{2}a_{4}-a_{3}^{2}\right| \leq \frac{1}{4q^{2}\psi_{2}^{2}} p_{1}^{2}, \tag{23}$$

where $k \in [0, 1]$ and p_j (j = 1, 2, 3) are positive and are the coefficients of the functions $p_k(z)$ defined by (6).

Proof. Making use of (19)–(21), we find that

$$\begin{aligned} a_{2}a_{4} - a_{3}^{2} &= \frac{A\left(q\right)}{16q^{2}\psi_{1}\psi_{3}} \ p_{1}^{2}c_{1}c_{3} + \left(\frac{A\left(q\right)\psi_{2}^{2} - \psi_{1}\psi_{3}}{16q^{2}\psi_{1}\psi_{2}^{2}\psi_{3}} \ p_{1}p_{2} - \frac{A\left(q\right)\psi_{2}^{2} - \psi_{1}\psi_{3}}{16q^{2}\psi_{1}\psi_{2}^{2}\psi_{3}} \ p_{1}^{2} \\ &+ \frac{A\left(q\right)\left(2 + q^{2}\right)\psi_{2}^{2} - 2\left(1 + q^{2}\right)\psi_{1}\psi_{3}}{64q^{2}\psi_{1}\psi_{3}} \ p_{1}^{3}\right)c_{1}^{2}c_{2} + \frac{1}{16q^{2}\psi_{2}^{2}} \ p_{1}^{2}c_{2}^{2} \\ &+ \left[\frac{A\left(q\right)}{64q^{2}\psi_{1}\psi_{3}} \ p_{1}p_{3} + \left(\frac{A\left(q\right)\psi_{2}^{2} - \psi_{1}\psi_{3}}{64q^{2}\psi_{1}\psi_{2}^{2}\psi_{3}}\right)p_{1}^{2} + \left(\frac{\psi_{1}\psi_{3} - A\left(q\right)\psi_{2}^{2}}{32q^{2}\psi_{1}\psi_{2}^{2}\psi_{3}}\right)p_{1}p_{2} \\ &+ \left(\frac{2\left(1 + q^{2}\right)\psi_{1}\psi_{3} - \left(2\left(1 + q^{2}\right) - q\right)A\left(q\right)\psi_{2}^{2}}{128q^{3}\psi_{1}\psi_{2}^{2}\psi_{3}}\right)p_{1}^{3} \\ &+ \left(\frac{A\left(q\right)\left(4q^{2} - 3q + 2\right)\psi_{2}^{2} - 2\left(1 + q^{2}\right)\psi_{1}\psi_{3}}{128q^{3}\psi_{1}\psi_{2}^{2}\psi_{3}}\right)p_{1}^{4} - \frac{1}{64q^{2}\psi_{2}^{2}} \ p_{2}^{2}\right]c_{1}^{4}, \end{aligned} \tag{24}$$

where

$$A(q) = \frac{(1+q)^2}{1+q+q^2}.$$

We substitute the values of c_2 and c_3 from the above Lemma and, for simplicity, take $Y = 4 - c_1^2$ and $Z = (1 - |x|^2)z$. Without loss of generality, we assume that $c = c_1$ ($0 \le c \le 2$), so that

$$a_{2}a_{4} - a_{3}^{2} = \left[\frac{q\left(1-q\right)A\left(q\right)\psi_{2}^{2}}{128q^{2}\psi_{1}\psi_{3}}p_{1}^{3} + \frac{A\left(q\right)}{64q^{2}\psi_{1}\psi_{3}}p_{1}p_{3} + \left(\frac{A\left(q\right)\left(4q^{2}-3q+2\right)\psi_{2}^{2}-2\left(1+q^{2}\right)\psi_{1}\psi_{3}}{128q^{3}\psi_{1}\psi_{2}^{2}\psi_{3}}\right)p_{1}^{2}p_{2} + \left(\frac{A\left(q\right)\left(q^{2}+2q-1\right)\psi_{2}^{2}-\left(1+q^{2}\right)^{2}\psi_{1}\psi_{3}}{256q^{4}\psi_{1}\psi_{2}^{2}\psi_{3}}\right)p_{1}^{4} - \frac{1}{64q^{2}\psi_{2}^{2}}p_{2}^{2}\right]c^{4} + \left[\frac{A\left(q\right)\psi_{2}^{2}-\psi_{1}\psi_{3}}{32q^{2}\psi_{1}\psi_{2}^{2}\psi_{3}}p_{1}p_{2} + \frac{A\left(q\right)\left(2+q^{2}\right)\psi_{2}^{2}-2\left(1+q^{2}\right)\psi_{1}\psi_{3}}{128q^{2}\psi_{1}\psi_{3}}p_{1}^{3}\right]c^{2}xY \\ \cdot \left[-\frac{A\left(q\right)}{64q^{2}\psi_{1}\psi_{3}}p_{1}^{2}c^{2}Yx^{2} - \frac{1}{64q^{2}\psi_{2}^{2}}p_{1}^{2}x^{2}Y^{2} + \frac{A\left(q\right)}{32q^{2}\psi_{1}\psi_{3}}p_{1}^{2}cYZ\right].$$
(25)

Upon setting $Z = (1 - |x|^2)z$ and taking the moduli in (25) and using trigonometric inequality, we find that

$$\begin{aligned} \left| a_{2}a_{4} - a_{3}^{2} \right| &\leq \left| \lambda_{1} \right| c^{4} + \left| \lambda_{2} \right| \left| x \right| Yc^{2} + \frac{A\left(q\right)}{64q^{2}\psi_{1}\psi_{3}} p_{1}^{2}Y \left| x \right|^{2}c^{2} \\ &+ \frac{1}{64q^{2}\psi_{2}^{2}} p_{1}^{2} \left| x \right|^{2}Y^{2} + \frac{A\left(q\right)}{32q^{2}\psi_{1}\psi_{3}} p_{1}^{2}c^{2}Y \left(1 - \left| x \right|^{2} \right) \\ &= \Lambda\left(c, \left| x \right| \right), \end{aligned}$$

$$(26)$$

where

$$\begin{split} \lambda_1 &= \frac{q \left(1-q\right) A \left(q\right) \psi_2^2}{128 q^2 \psi_1 \psi_3} \ p_1^3 + \frac{A \left(q\right)}{64 q^2 \psi_1 \psi_3} \ p_1 p_3 \\ &+ \left(\frac{A \left(q\right) \left(4 q^2 - 3 q + 2\right) \psi_2^2 - 2 \left(1+q^2\right) \psi_1 \psi_3}{128 q^3 \psi_1 \psi_2^2 \psi_3}\right) p_1^2 p_2 \\ &+ \left(\frac{A \left(q\right) \left(q^2 + 2 q - 1\right) \psi_2^2 - \left(1+q^2\right)^2 \psi_1 \psi_3}{256 q^4 \psi_1 \psi_2^2 \psi_3}\right) p_1^4 - \frac{1}{64 q^2 \psi_2^2} \ p_2^2 \\ \lambda_2 &= \frac{A \left(q\right) \psi_2^2 - \psi_1 \psi_3}{32 q^2 \psi_1 \psi_2^2 \psi_3}; p_1 p_2 + \frac{A \left(q\right) \left(2+q^2\right) \psi_2^2 - 2 \left(1+q^2\right) \psi_1 \psi_3}{128 q^2 \psi_1 \psi_3} \ p_1^3. \end{split}$$

Now, trivially, we have

$$\Lambda'\left(|x|\right) > 0$$

on [0, 1], and so

$$\Lambda\left(\left|x\right|\right) \leqq \Lambda\left(1\right).$$

Hence, by puting $Y = 4 - c_1^2$ and after some simplification, we have

$$\begin{aligned} \left| a_{2}a_{4} - a_{3}^{2} \right| &= \left(\left| \lambda_{1} \right| - \left| \lambda_{2} \right| + \frac{\psi_{1}\psi_{3} - A\left(q\right)\psi_{2}^{2}}{64q^{2}\psi_{1}\psi_{3}} p_{1}^{2} \right) c^{4} \\ &+ \left(4\left| \lambda_{2} \right| + \left(\frac{A\left(q\right)\psi_{2}^{2} - \psi_{1}\psi_{3}}{16q^{2}\psi_{1}\psi_{3}} p_{1}^{2} \right) \right) c^{2} + \frac{1}{4q^{2}\psi_{2}^{2}} p_{1}^{2} \\ &= G\left(c\right). \end{aligned}$$

$$(27)$$

For optimum value of G(c), we consider G'(c) = 0, which implies that c = 0. So G(c) has a maximum value at c = 0. We therefore conclude that the maximum value of G(c) is given by

$$\frac{1}{4q^2\psi_2^2}p_1^2,$$

which occurs at c = 0 or

$$c^{2} = -\frac{128 |\lambda_{2}| q^{2} \psi_{1} \psi_{3} + 4A(q) \psi_{2}^{2} - 2\psi_{1} \psi_{3} p_{1}^{2}}{(64q^{2} (|\lambda_{1}| - |\lambda_{2}|) \psi_{1} \psi_{3} + \psi_{1} \psi_{3} - A(q) \psi_{2}^{2} p_{1}^{2})}$$

This completes the proof of Theorem 3. \Box

If we put $\psi_{n-1} = 1$ and let $q \to 1-$ in Theorem 3, we have the following known result.

Corollary 2 (see [26]). *If the function* f(z) *given by* (1) *belongs to the class* k-ST*, where* $k \in [0, 1]$ *, then*

$$\left|a_2a_4-a_3^2\right| \leq \frac{p_1^2}{4}.$$

If we put

$$p_1 = 2$$
 and $\psi_{n-1} = 1$,

by letting $q \rightarrow 1-$ in Theorem 3, we have the following known result.

Corollary 3 (see [18]). *If* $f \in S^*$, *then*

$$\left|a_2a_4-a_3^2\right| \leq 1.$$

By letting k = 1, $\psi_{n-1} = 1$, $q \rightarrow 1$ – and

$$p_1 = \frac{8}{\pi^2}, \quad p_2 = \frac{16}{3\pi^2} \quad \text{and} \quad p_3 = \frac{184}{45\pi^2}$$

in Theorem 3, we have the following known result.

Corollary 4 (see [27]). If the function f(z) given by (1) belong to the class SP, then

$$\left|a_2a_4-a_3^2\right| \leq \frac{16}{\pi^4}.$$

4. Concluding Remarks and Observations

Motivated significantly by a number of recent works, we have made use of a certain general conic domain and the quantum (or q-) calculus in order to define and investigate a new subclass of normalized analytic functions in the open unit disk \mathbb{U} , which we have referred to as q-starlike functions. For this q-starlike function class, we have successfully derived several properties and characteristics. In particular, we have found the Hankel determinant and the Toeplitz matrices for this newly-defined class of q-starlike functions. We also highlight some known consequences of our main results which are stated and proved as theorems and corollaries.

Author Contributions: conceptualization, Q.Z.A. and N.K. (Nazar Khan); methodology, N.K. (Nasir Khan); software, B.K.; validation, H.M.S.; formal analysis, H.M.S.; writing—original draft preparation, H.M.S.; writing—review and editing, H.M.S.; supervision, H.M.S.

Funding: This research received no external funding.

Conflicts of Interest: The authors declare no conflict of interest.

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