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Norm-Resolvent Convergence for Neumann Laplacians on Manifold Thinning to Graphs

Kirill D. Cherednichenko ¹ , Yulia Yu. Ershova ² and Alexander V. Kiselev ^{1,*} 

¹ Department of Mathematical Sciences, University of Bath, Claverton Down, Bath BA2 7AY, UK; k.cherednichenko@bath.ac.uk

² Department of Mathematics, Texas A&M University, College Station, TX 77843-3368, USA; yuershova@tamu.edu

* Correspondence: alexander.v.kiselev@gmail.com or ak2084@bath.ac.uk

Abstract: Norm-resolvent convergence with an order-sharp error estimate is established for Neumann Laplacians on thin domains in \mathbb{R}^d , $d \geq 2$, converging to metric graphs in the limit of vanishing thickness parameter in the “resonant” case. The vertex matching conditions of the limiting quantum graph are revealed as being closely related to those of the δ' type.

Keywords: PDE; quantum graphs; generalised resolvent; thin structures; norm-resolvent asymptotics

MSC: 35C20; 35J25; 47A10; 47A20; 81Q35; 74K99

1. Introduction

In the book [1] (see also references therein, in particular [2–4]), Neumann Laplacians A_ε on thin manifolds, converging to metric graphs as $\varepsilon \rightarrow 0$, were studied; see, e.g., Figure 1. Here, ε represents the “thickness” of the manifolds in those parts where they converge to the graph edges. The named works attacked the question of spectral (and, in the case of [1], norm-resolvent) convergence of such partial differential operators to a graph Laplacian with certain matching conditions at the graph vertices. This latter ordinary differential operator is introduced as follows. Denoting by E the set of edges e of the limiting graph G , each $e \in E$ can be identified with the interval $[0, l_e]$, where l_e is the length of e . We use notation $L^2(e)$ for the associated Hilbert space $L^2(0, l_e)$. Similarly, we denote $L^2(G) := \oplus_e L^2(e)$. The graph Laplacian A_G is generated by the differential expression $-u''$, $u \in H^2(e)$ on each edge $e \in E$ separately (see [5] for details), subject to the vertex conditions discussed below.

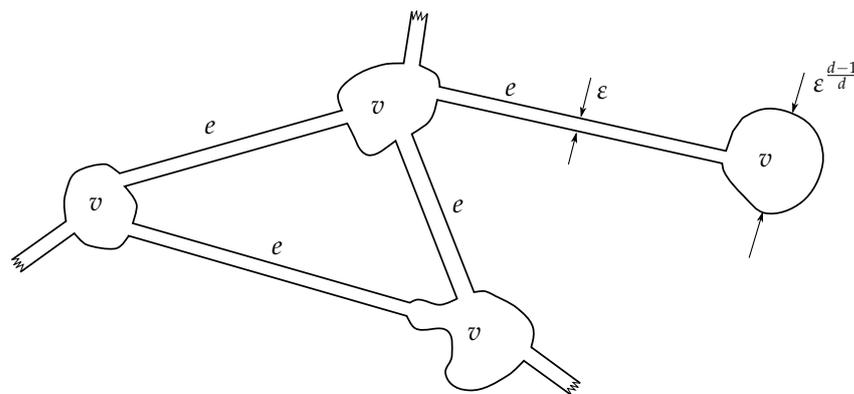


Figure 1. An example of a thin network in the case of a resonant scaling of the edge thickness and vertex diameter. The (non-dimensional) parameter ε is the ratio of the actual size to the wavelength. As $\varepsilon \rightarrow 0$, the structure converges to a graph; the related notation is introduced at the start of Section 2.



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It was proved in [2–4] that, within any compact $K \in \mathbb{C}$, the spectra of A_ε converge in the Hausdorff sense to the spectrum of a graph Laplacian A_G . In the book [1], the claimed convergence was enhanced to the norm-resolvent type, with an explicit control of the error as $O(\varepsilon^\gamma)$, where $\gamma > 0$ depends on whether the ambient space is two-dimensional. The matching conditions at the vertices of the limiting graph turn out to be any one of the following:

- (i) Kirchhoff (i.e., standard) if the vertex volumes are decaying, as $\varepsilon \rightarrow 0$, faster than the edge volumes;
- (ii) “Resonant”, which can be equivalently described in terms of δ -type matching conditions with coupling constants proportional to the spectral parameter z , see [2,3], if the vertex and edge volumes are of the same order;
- (iii) “Dirichlet-decoupled” (i.e., the graph Laplacian becomes completely decoupled) if the vertex volumes vanish slower than the edge ones.

We also refer the reader to the frequently overlooked papers [6,7], where an alternative approach to the asymptotic analysis of thin networks, aimed at capturing the resonant and scattering features, is developed; see also the references therein. In the present paper we consider the case of scalar PDEs, although our approach can be immediately generalised to the case of PDE systems, in particular, in the context of thin elastic structures with applications to, e.g., pentamodes (see [8,9]) by utilising the results of [10].

In the present paper, we are primarily interested in the most non-trivial resonant case (ii). We provide a straightforward, alternative to [1], proof of the fact that the Neumann Laplacians A_ε in this case converge in a norm-resolvent sense to a linear operator acting in the Hilbert space $L^2(G) \oplus \mathbb{C}^N$, where N is the number of vertices. The mentioned limiting operator is in fact the one first pointed out in [4] as a self-adjoint operator whose spectrum coincides with the Hausdorff limit of spectra for the family A_ε .

On the technical side, our approach can be seen as a modification of the one developed by us in [10,11]; see also [12–14]. We specifically point out that the framework originally developed for high-contrast homogenisation admits a natural generalisation to setups where the “contrast” is achieved by purely geometric means, including (but not limited to) thin networks, as in the present paper. This seems to widen significantly the range of dimension-reduction-type models that are amenable to this kind of analysis and thus to establish a transparent connection between previously unrelated physical contexts, providing for a possibility to develop new types of media in materials science.

We obtain a better error bound than [1] (Section 6.7) (and skip a visit to the “plumber’s shop” of [1,15]). Our estimate in the planar case is $O(\varepsilon |\log \varepsilon|)$ and in the case of \mathbb{R}^d , $d \geq 3$, it is $O(\varepsilon)$. Unlike that of [1], our method does not allow us to study the full set of asymptotic regimes in (iii). This is due to the fact that the argument of our paper [11] is based on the Dirichlet-to-Neumann (DN) machinery. There is a possibility to modify the approach by invoking Neumann-to-Dirichlet maps instead, which would have two advantages: one could consider all rates of vertex volume decay in (iii), and certain geometric smoothness requirements could be somewhat relaxed. Nevertheless, in this paper we stick with the DN version of the approach in order to align the exposition with that of [11]. The alternative strategy will be followed up elsewhere, both in the present context and in the setting of [11].

The above results of course imply the Hausdorff spectral convergence, at the same time yielding a sharp estimate on its rate. Moreover, in contrast to [1], our approach allows one to consider “high-frequency” regimes, i.e., setups where the spectral parameter (which in the wave propagation context may represent the square of the frequency) is no longer constrained to a compact set but is still constrained by some negative power of the small parameter ε . (In non-dimensional terms this corresponds to the wavelength being of the order of some positive power of ε). This leads to a sequence of “effective”, dimensionally reduced, models of the thin structure, which are sequentially applicable for a set of (asymptotic) frequency intervals. The complexity of the dimension reduction process for these models increases along the sequence. While an initial result regarding the

high-frequency situation is presented below (which suffices to reveal a metamaterial in a periodic thin network of [16]), we postpone the full analysis to a future publication.

Alongside the high-frequency analysis, yet another sequence of models will be revealed by a version of the same argument. This corresponds to the transition from the resonant regime to the Dirichlet-decoupled one and will allow us to reconcile the asymptotic analysis of [1] with that of [7] by introducing “transitional” models of increasing complexity. In these transitional regimes, from the point of view of [1], one gets arbitrarily close to the Dirichlet-decoupled situation, see (iii) above, whereas from the point of view of [7] (where the vertex volumes do not decay at all), one faces a highly non-trivial picture of resonant scattering.

Aiming at better clarity, in the present paper we restrict ourselves to the case where (without much loss of generality) the edge subdomains are assumed straight and uniformly thin, whereas the vertex subdomains are smooth with, possibly, the exception of the points where they meet the edges (see Section 2 for further details). These assumptions appear very natural in view of the possible examples shown in Figures 1 and 2.

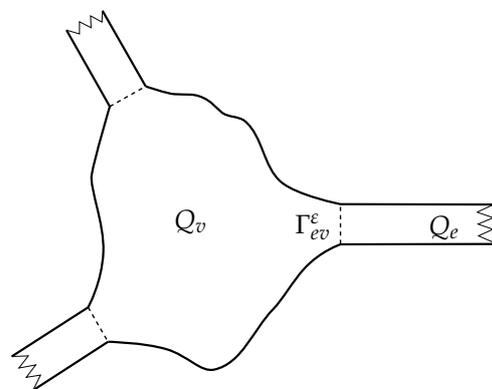


Figure 2. An example of a vertex subdomain Q_v and an edge subdomain Q_e separated by a contact plate Γ_{ev}^ϵ of linear size ϵ . The volumes of Q_v and Q_e are assumed to be of the same order ϵ^{d-1} , which is also the area of the contact plate. As in Figure 1, the (non-dimensional) parameter ϵ is the ratio of the actual size to the wavelength.

The ultimate section of the paper is dedicated to the spectral analysis of the effective graph Laplacian. We show that effective matching conditions at graph vertices, having been commonly treated as impedance type (namely, linear in the spectral parameter), are in fact of δ' type after an application of a unitary gauge. This latter, however, introduces constant magnetic potentials on the edges of the limit graph. It is notable that the presence of such a magnetic field leads to a phase transition to a medium exhibiting double-negative metamaterial properties; see [16]. This observation will permit us to construct thin periodic networks with negative group velocity without the need for an external magnetic field, by relying instead on the resonant geometric properties of the network.

2. Problem Setup and Preliminaries

In what follows, we consider a prototypical setup only, which already presents all the challenges appearing in the general case. We also refer the reader to [1], where the most general setup is meticulously introduced.

For the limiting graph, the following notation will be used: the metric graph G will be identified with the set of edges E , so each individual edge is denoted by $e \in E$ and is associated with an interval $[0, l_e]$. We denote by V the set of graph vertices and treat each $v \in V$ as the set of edge endpoints meeting at v . The graph G is assumed to be oriented throughout.

Proceeding to the setup for the Neumann Laplacian on a thin-graph-convergent structure, let a connected (open) domain Q be the union of the “vertex” part Q_V , the “edge” part Q_E , and the interface boundary Γ_v^ϵ between the two, where Q_E will be assumed to be a

finite collection of ε -thin cylinders (for $d = 2$, rectangular boxes), $Q_E = \cup_e Q_e$ of lengths l_e . For each e , the domain Q_e is assumed to be, up to a linear spatial transform (i.e., shift and rotation), defined by (for $d = 2$)

$$Q_e = \{x \in \mathbb{R}^2 : x_1 \in (0, l_e), x_2 \in (0, \varepsilon)\} = (0, l_e) \times (0, \varepsilon).$$

In dimensions $d \geq 3$, the set Q_e is defined as the direct product of the interval $(0, l_e)$ and a smooth cross-section Q_e^ε of volume ε^{d-1} .

It is further assumed that $Q_V = \cup_v Q_v$, where each of the disjoint domains Q_v is connected and has piecewise smooth boundary ∂Q_v , which can be decomposed as $\partial Q_v = \tilde{\Gamma}_v^\varepsilon \cup \Gamma_{ev}^\varepsilon$, where Γ_{ev}^ε is the interface between the vertex and edge domains and $\tilde{\Gamma}_v^\varepsilon$ is the remaining part of ∂Q_v . Henceforth, we often drop the dependence on ε in the notation for those geometric elements where it can be removed by a rescaling or is otherwise irrelevant in the analysis to follow.

The contact part Γ_{ev}^ε is further decomposed into a union of flat plates (for $d = 2$, straight segments): $\Gamma_{ev}^\varepsilon = \cup \Gamma_{ev}^\varepsilon$. Here, the union is taken over all edge domains Q_e connected to Q_v so that $\overline{Q_e} \cap \overline{Q_v} = \Gamma_{ev}^\varepsilon$. In what follows, we will refer to the segments Γ_{ev}^ε as *contact plates*. Since operators of Zaremba (or mixed) boundary value problem [17] will be used below, we further require that the contact plates Γ_{ev}^ε meet $\tilde{\Gamma}_v^\varepsilon$ at angles strictly less than π ; see [18,19] for further details. We will further assume that the curves $\tilde{\Gamma}_v^\varepsilon$ are smooth.

Furthermore, we assume that for all vertices v the domains Q_v have volumes of order ε^{d-1} and are obtained by suitable transformations of an ε -independent domain Q_v^0 containing the origin. More precisely, we assume that for all v and $\varepsilon > 0$ one has

$$Q_v = \varepsilon^{\frac{d-1}{d}} \mathcal{T}_v^\varepsilon(Q_v^0) + b_v^\varepsilon, \tag{1}$$

where $b_v^\varepsilon \in \mathbb{R}^d$, and the domain $\mathcal{T}_v^\varepsilon(Q_v^0)$ is obtained from Q_v^0 by a suitable homeomorphism $\mathcal{T}_v^\varepsilon$ that (a) preserves volume and (b) maintains the angles between the contact plates and the adjacent parts of the boundary to be uniformly less than π , cf. [17,18].

For example, if Q_v is star-shaped with respect to the origin (recall that a domain Q_v is said to be star-shaped with respect to $x_v \in Q_v$ if for all $x \in Q_v$ one has $\{x_v + t(x - x_v), t \in [0, 1]\} \subset Q_v$), the deformation $\mathcal{T}_v^\varepsilon$ can be constructed by “cut-and-glue surgery” applied to Q_v^0 followed by a suitable “radial” scaling, as follows. The domain Q_v^0 is first split into several conical (sectorial for $d = 2$) subdomains by making cuts alongside the cone “generatrices” (“radii” for $d = 2$) from the origin to the boundary points (endpoints for $d = 2$) of all subsets of ∂Q_v^0 pertaining to the contact plates of Q_v . Each cone is then transformed by a suitable change of variable so that (A) the images of the cone bases pertaining to the contact plates of Q_v have linear size $\varepsilon^{1/d}$ and (B) the image of the complementary part of Q_v^0 is reattached to the lateral boundaries of the cones in A. Note that this procedure in general changes the domain volume by a small quantity. In order to restore the volume, a suitable “radial” scaling (centred at the origin) is applied, ensuring that it does not affect any of the mentioned conical subsets.

We remark that the representation (1) guarantees that we are in the “resonant” (or “borderline”) case of [2,3], i.e., that the volumes of the contact plates are proportional to the volumes of vertex domains.

The requirement that the homeomorphism $\mathcal{T}_v^\varepsilon$ in (1) be volume preserving is not essential and can be removed, as long as the ratios of the volumes of the vertex domains to the volumes of the contact plates are bounded below. In this case, the corresponding final convergence estimates, similar to those we derive below (see Theorem 3), will in general depend on the said ratios, as is evident from the formula (14).

On the domain Q , we consider a family of self-adjoint operators A_ε defined by their sesquilinear forms and corresponding to the differential expression $-\Delta$ subject to Neumann boundary conditions. Other types of boundary conditions can be considered as well, including those of Robin and Dirichlet [6,20], in which case the lower edge of the spectrum is ε -dependent, which in the context of homogenisation corresponds to the so-called high-

frequency regime or, using the terminology of M. S. Birman, to “the neighbourhood of the edge of a gap” [21,22].

Consider $\sigma > 0$ and an arbitrary non-negative function $R = R(\varepsilon)$, $\varepsilon \in (0, 1)$ such that $R(\varepsilon)\varepsilon^{1/2}|\log \varepsilon|^\gamma \rightarrow 0$ as $\varepsilon \rightarrow 0$, where $\gamma = 1$ for $d = 2$ and $\gamma = 0$ otherwise. In what follows, we will deal with the family of resolvents $(A_\varepsilon - z)^{-1}$. We shall always assume that $z \in \mathbb{C}$ is separated from the spectrum of the original operator family. In particular cases, we will assume that either z is constrained to the ε -growing compact set

$$K_\sigma^\varepsilon := B_{R(\varepsilon)}(0) \setminus \{z \in \mathbb{C} \mid \text{dist}(z, \mathbb{R}) \geq \sigma\},$$

where $B_{R(\varepsilon)}(0)$ denotes the ball of radius $R(\varepsilon)$ centered at the origin, or

$$z \in K_\sigma := \{z \in \mathbb{C} \mid z \in K \text{ a compact set in } \mathbb{C}, \text{dist}(z, \mathbb{R}) \geq \sigma\}.$$

In particular, for all ε small enough, one has $K_\sigma \subset K_\sigma^\varepsilon$. After we have established the operator-norm asymptotics of $(A_\varepsilon - z)^{-1}$ for $z \in K_\sigma$, the result is extended by analyticity to a compact set K_σ^{ext} whose distance to the spectrum of the leading order of the asymptotics is bounded below by σ , and the same statement holds in the case of K_σ^ε .

Proceeding similarly to, e.g., [23] in the related area of critical-contrast homogenisation and facilitated by the abstract framework of [24], we consider A_ε as operators of transmission problems (see [25] and references therein) relative to the *internal boundary* $\Gamma^\varepsilon := \cup_{e,v} \Gamma_{ev}^\varepsilon$. The transmission problem is formulated as, given a function $f \in L^2(Q)$, finding the weak solution $u \in L^2(Q)$ of the boundary value problem

$$\begin{cases} -\Delta u(x) - zu(x) = f(x), & x \in Q_V \text{ or } x \in Q_E, \\ u_v(x) = u_e(x), \quad \frac{\partial u_v}{\partial n} - \frac{\partial u_e}{\partial n} = 0 & \text{on } \Gamma_{ev}^\varepsilon, \\ \frac{\partial u}{\partial n} = 0 & \text{on } \partial Q. \end{cases} \tag{2}$$

Here, $u_v := u|_{Q_v}$, $u_e := u|_{Q_e}$ for all admissible e and v (i.e. when $\Gamma_{ev}^\varepsilon \neq \emptyset$), and n represents the exterior normal on ∂Q and the “edge-inward” normal (i.e., directed from Q_v to Q_e) on any of contact plates Γ_{ev}^ε . By a classical argument, the solution of the above problem is shown to be equal to $(A_\varepsilon - z)^{-1}f$. Note that the boundary value problem for the Neumann Laplacian on Q is given by the first and third lines in (2). While the second line is, strictly speaking, redundant, it proves important in order to view the problem as a transmission one relative to Γ^ε .

It remains to be seen that the linear operator of the transmission problem (2) defined via the technique in [24], which we briefly recall below, is the same operator A_ε ; the proof of this fact follows easily by combining [24] and the main estimate of [25].

Following the approach of [24] (cf. [26,27] and references therein for alternative approaches), which is based on the ideas of the classical Birman–Kreĭn–Višik theory (see [28–30]), the linear operator of the transmission boundary value problem is introduced as follows. Let $\mathcal{H} := L^2(\Gamma^\varepsilon) = \oplus_{e,v} L^2(\Gamma_{ev}^\varepsilon)$, and consider the harmonic lift operators Π_V and Π_E defined on $\phi \in \mathcal{H}$ via

$$\begin{aligned} \Pi_V \phi &:= u_\phi, \text{ where } \begin{cases} \Delta u_\phi = 0, & u_\phi \in L^2(Q_V), \\ u_\phi|_{\Gamma^\varepsilon} = \phi, \end{cases} \\ \Pi_E \phi &:= u_\phi, \text{ where } \begin{cases} \Delta u_\phi = 0, & u_\phi \in L^2(Q_E), \\ u_\phi|_{\Gamma^\varepsilon} = \phi, \end{cases} \end{aligned}$$

subject to Neumann boundary conditions on ∂Q . These operators are first defined on $\phi \in C_0^2(\Gamma^\varepsilon)$, in which case the corresponding solutions u_ϕ can be seen as classical [17]. The results of [18] allow one to extend both harmonic lifts to bounded (in fact, compact) op-

erators on \mathcal{H} , in which case u_ϕ are to be treated as distributional solutions of the respective boundary value problems. The solution operator $\Pi : \mathcal{H} \mapsto L^2(Q) = L^2(Q_V) \oplus L^2(Q_E)$ is defined as follows:

$$\Pi\phi := \Pi_V\phi \oplus \Pi_E\phi.$$

Consider the self-adjoint operator family A_0 (choosing not to reflect the ε -dependence in the notation) to be the Dirichlet decoupling of the operator family A_ε , i.e., the operator of the boundary value problem on both Q_V and Q_E , where the Dirichlet boundary conditions are imposed on Γ^ε together with Neumann boundary conditions on ∂Q . The operator A_0 is generated by the same differential expression as A_ε . Clearly, one has $A_0 = A_0^V \oplus A_0^E$ relative to the orthogonal decomposition $L^2(Q) = L^2(Q_V) \oplus L^2(Q_E)$; all three operators A_0, A_0^V and A_0^E are self-adjoint and positive-definite. Moreover, by [18,31] there exists a bounded inverse A_0^{-1} . Note that $\text{dom } A_0 \cap \text{ran } \Pi = \emptyset$; see [24].

Furthermore, denoting by $\tilde{\Gamma}_0^V$ and $\tilde{\Gamma}_0^E$ the left inverse of Π_V and Π_E , respectively, one introduces the trace operator Γ_0^V (respectively, Γ_0^E) as the null extension of $\tilde{\Gamma}_0^V$ (respectively, $\tilde{\Gamma}_0^E$) to the domain $\text{dom } A_0^V \dot{+} \text{ran } \Pi_V$ (respectively, $\text{dom } A_0^E \dot{+} \text{ran } \Pi_E$). In the same way, we introduce the operator $\tilde{\Gamma}_0$ and its null extension Γ_0 to the domain $\text{dom } A_0 \dot{+} \text{ran } \Pi$.

The solution operators S_z^V, S_z^E of the boundary value problems

$$\begin{cases} -\Delta u_\phi - zu_\phi = 0, & u_\phi \in \text{dom } A_0^V \dot{+} \text{ran } \Pi_V, \\ \Gamma_0^V u_\phi = \phi, \end{cases}$$

$$\begin{cases} -\Delta u_\phi - zu_\phi = 0, & u_\phi \in \text{dom } A_0^E \dot{+} \text{ran } \Pi_E, \\ \Gamma_0^E u_\phi = \phi \end{cases}$$

are defined as linear mappings from ϕ to u_ϕ , respectively. These operators are bounded from $L_2(\Gamma^\varepsilon)$ to $L_2(Q_V)$ and $L_2(Q_E)$, respectively, and admit the following representations:

$$S_z^E = (1 - z(A_0^E)^{-1})^{-1}\Pi_E, \quad S_z^V = (1 - z(A_0^V)^{-1})^{-1}\Pi_V.$$

The solution operator S_z from $L^2(\Gamma^\varepsilon)$ to $L^2(Q_V) \oplus L^2(Q_E)$ is now defined as $S_z = S_z^V \oplus S_z^E$; it admits the representation $S_z = (1 - z(A_0)^{-1})^{-1}\Pi$ and is bounded.

Having introduced orthogonal projections P_V and P_E from $L^2(Q)$ onto $L^2(Q_V)$ and $L^2(Q_E)$, respectively, one has the obvious identities

$$S_z^V = P_V S_z, \quad S_z^E = P_E S_z, \quad \Pi_V = P_V \Pi, \quad \Pi_E = P_E \Pi.$$

Fix self-adjoint (and, in general, unbounded) operators Λ^E, Λ^V defined on domains $\text{dom } \Lambda^E, \text{dom } \Lambda^V \subset L^2(\Gamma^\varepsilon)$ (in what follows these operators will be chosen as DN maps of Zaremba problems on Q_E and Q_V , respectively, and well-defined on $H^1(\Gamma^\varepsilon)$, where $H^1(\Gamma^\varepsilon)$ is the standard Sobolev space pertaining to the internal boundary Γ^ε). Still following [24], we define the “second boundary operators” Γ_1^E and Γ_1^V to be linear operators on the domains

$$\text{dom } \Gamma_1^E := \text{dom } A_0^E \dot{+} \Pi_E \text{dom } \Lambda^E, \quad \text{dom } \Gamma_1^V := \text{dom } A_0^V \dot{+} \Pi_V \text{dom } \Lambda^V.$$

The action of $\Gamma_1^{E(V)}$ is set by:

$$\Gamma_1^E : (A_0^E)^{-1}f \dot{+} \Pi_E\phi \mapsto \Pi_E^*f + \Lambda^E\phi, \quad \Gamma_1^V : (A_0^V)^{-1}f \dot{+} \Pi_V\phi \mapsto \Pi_V^*f + \Lambda^V\phi$$

for all $f \in L^2(Q_E), \phi \in \text{dom } \Lambda^E$ and $f \in L^2(Q_V), \phi \in \text{dom } \Lambda^V$, respectively.

Alongside $\Gamma_1^{E(V)}$, introduce a self-adjoint operator Λ on $\text{dom } \Lambda \subset \mathcal{H}$ and then the following “boundary” operator Γ_1 :

$$\text{dom } \Gamma_1 := \text{dom } A_0^{-1} \dot{+} \Pi \text{dom } \Lambda,$$

$$\Gamma_1 : A_0^{-1}f \dot{+} \Pi\phi \mapsto \Pi^*f + \Lambda\phi \quad \forall f \in L^2(Q), \phi \in \text{dom } \Lambda.$$

We remark that the operators $\Gamma_1, \Gamma_1^{E(V)}$ thus defined are assumed to be neither closed nor indeed closable.

In our setup, we make the following concrete choice of the operators $\Lambda^{E(V)}$: in what follows, they are the DN maps pertaining to the components Q_E and Q_V , respectively. More precisely, for the problem

$$\begin{aligned} \Delta u_\phi &= 0, & u_\phi &\in L^2(Q_E), \\ u_\phi|_{\Gamma^\varepsilon} &= \phi, & \partial_n u_\phi|_{\partial Q} &= 0, \end{aligned}$$

the operator Λ^E maps the boundary values ϕ of u_ϕ to the negative traces of its normal derivative

$$\partial_n u_\phi|_{\Gamma^\varepsilon} := \nabla u \cdot n|_{\Gamma^\varepsilon},$$

where $n = -n_E$ is as above the “edge-inward” normal. This operator is well defined by its sesquilinear form as a self-adjoint operator on $L^2(\Gamma^\varepsilon)$ (see, e.g., [32,33]), and $\text{dom } \Lambda^E \supset H^1(\Gamma^\varepsilon)$ by [18]. The above definition is inspired by [24]. Note that the operator thus defined is negative the classical DN map of, e.g., [23].

On the vertex part Q_V , we consider the problem

$$\begin{aligned} \Delta u_\phi &= 0, & u_\phi &\in L^2(Q_V), \\ u_\phi|_{\Gamma^\varepsilon} &= \phi, & \partial_n u_\phi|_{\partial Q} &= 0, \end{aligned}$$

and define Λ^V as the operator mapping the boundary values ϕ of u_ϕ to the negative traces of its normal derivative $-\partial_n u_\phi|_{\Gamma^\varepsilon}$, where $n = n_V$ is again the “edge-inward” normal. The self-adjointness of Λ^V on $\text{dom } \Lambda^V \supset H^1(\Gamma^\varepsilon)$ follows by an unchanged argument.

Finally we introduce the operator Λ which on $\phi \in H^1(\Gamma^\varepsilon)$ is the sum $\Lambda\phi = \Lambda^V\phi + \Lambda^E\phi$. It is also a self-adjoint operator on $\text{dom } \Lambda \supset H^1(\Gamma^\varepsilon)$. This can be ascertained either by the argument of [25], in which case it is defined as the inverse of a compact self-adjoint operator on the orthogonal complement to constants $L^2(\Gamma^\varepsilon) \ominus \{c\mathbf{1}\}$, extended to $\{c\mathbf{1}\}$ by zero or, alternatively, from its definition by a closed sesquilinear form.

The choice of $\Lambda^{E(V)}$ made above allows us to consider Γ_1 on the domain $\text{dom } A_0 \dot{+} \Pi \text{dom } \Lambda$. One then writes [24] the second Green identity in the following form:

$$\langle Au, v \rangle_{L^2(Q)} - \langle u, Av \rangle_{L^2(Q)} = \langle \Gamma_1 u, \Gamma_0 v \rangle_{L^2(\Gamma^\varepsilon)} - \langle \Gamma_0 u, \Gamma_1 v \rangle_{L^2(\Gamma^\varepsilon)}$$

for all $u, v \in \text{dom } \Gamma_1 = \text{dom } A_0 \dot{+} \Pi \text{dom } \Lambda$, where the operator A is the null extension (see [34]) of the operator A_0 onto $\text{dom } \Gamma_1$. Thus the triple $(\mathcal{H}, \Gamma_0, \Gamma_1)$ is closely related to a boundary quasi-triple of [26] (see also [35]) for the transmission problem considered; cf. [27] for an alternative approach.

The calculation of Π^* in [24] shows that $\Pi^* = \Gamma_1 A_0^{-1}$ and therefore Γ_1 as introduced above acts as follows:

$$\Gamma_1 : u = P_E u + P_V u \mapsto \partial_n P_E u|_{\Gamma^\varepsilon} - \partial_n P_V u|_{\Gamma^\varepsilon},$$

where P_E and P_V are the orthogonal projections of $L^2(Q)$ onto $L^2(Q_E), L^2(Q_V)$, respectively. Therefore, transmission problem at hand (at least, formally so far) corresponds to the interface (or “matching”) condition $\Gamma_1 u = 0$.

Definition 1 ([24]). *The operator-valued function $M(z)$ defined on the domain $\text{dom } \Lambda$ for $z \in \rho(A_0)$ (and in particular, for $z \in K_\sigma^\varepsilon$) by the formula*

$$M(z)\phi = \Gamma_1 S_z \phi = \Gamma_1 (1 - zA_0^{-1})^{-1} \Pi\phi \tag{3}$$

is referred to as the M-function of the problem (2).

The following result of [24] summarises the properties of the M -function which we will need in what follows.

Proposition 1 ([24], Theorem 3.3). *1. One has the following representation:*

$$M(z) = \Lambda + z\Pi^*(1 - zA_0^{-1})^{-1}\Pi, \quad z \in \rho(A_0). \tag{4}$$

2. $M(z)$ is an analytic operator-function with values in the set of closed operators in $L_2(\Gamma^\varepsilon)$ densely defined on the z -independent domain $\text{dom } \Lambda$.

3. For $z, \zeta \in \rho(A_0)$ the operator $M(z) - M(\zeta)$ is bounded and

$$M(z) - M(\zeta) = (z - \zeta)S_z^*S_\zeta$$

*In particular, $\Im M(z) = (\Im z)S_z^*S_z$ and $(M(z))^* = M(\bar{z})$.*

4. For $u_z \in \ker(A - zI) \cap \{\text{dom } A_0 \dot{+} \Pi \text{dom } \Lambda\}$, the following formula holds:

$$M(z)\Gamma_0 u_z = \Gamma_1 u_z.$$

Alongside $M(z)$, we define $M_V(z)$ and $M_E(z)$, which pertain to the vertex Q_V and edge Q_E parts of the domain Q , respectively, by the formulae

$$\begin{aligned} M_V(z)\phi &= \Gamma_1^V S_z^V \phi = \Gamma_1^V (1 - z(A_0^V)^{-1})^{-1} \Pi_V \phi, \\ M_E(z)\phi &= \Gamma_1^E S_z^E \phi = \Gamma_1^E (1 - z(A_0^E)^{-1})^{-1} \Pi_E \phi. \end{aligned} \tag{5}$$

As before, in the notation we suppress the dependence on the parameter ε for brevity.

The value of the fact that for $\phi \in H^1(\Gamma^\varepsilon)$ one has $M(z)\phi = M_E(z)\phi + M_V(z)\phi$ is clear: in contrast to A_ε , which cannot be additively decomposed into "independent" terms pertaining to the vertex and edge parts of the medium Q owing to the transmission interface conditions on Γ^ε , the M -function is additive (see, e.g., [11,36], where this property was observed and exploited in the related settings of homogenisation and scattering, respectively). In what follows, we will observe that the resolvent $(A_\varepsilon - z)^{-1}$ can be expressed in terms of $M(z)$ via a version of the celebrated Kreĭn formula, thus reducing the asymptotic analysis of the resolvent to that of the corresponding M -function (see, e.g., [26,37] for alternative approaches to derivation of the Kreĭn formula in our setting).

Alongside the transmission problem (2), the boundary conditions of which can be now (so far, formally) represented as $u \in \text{dom } A_0 \dot{+} \Pi L^2(\Gamma^\varepsilon)$, $\Gamma_1 u = 0$, in what follows we will require a wider class of problems of this type. This class is formally given by the transmission conditions

$$u \in \text{dom } A \equiv \text{dom } A_0 \dot{+} \Pi L^2(\Gamma^\varepsilon), \quad \beta_0 \Gamma_0 u + \beta_1 \Gamma_1 u = 0,$$

where β_1 is a bounded operator on $L^2(\Gamma^\varepsilon)$ and β_0 is a linear operator defined on the domain $\text{dom } \beta_0 \supset \text{dom } \Lambda$.

In general, the operator $\beta_0 \Gamma_0 + \beta_1 \Gamma_1$ is not defined on the domain $\text{dom } A$. This problem is being taken care of by the following assumption, which will be satisfied throughout:

$$\beta_0 + \beta_1 \Lambda \text{ defined on } \text{dom } \Lambda \text{ is closable in } \mathcal{H}.$$

We remark that by Proposition 1 the operators $\beta_0 + \beta_1 M(z)$ are then closable for all $z \in \rho(A_0)$, and the domains of their closures coincide with $\text{dom } \overline{\beta_0 + \beta_1 \Lambda}$.

For any $f \in H$, $\phi \in \text{dom } \Lambda$, the equality

$$(\beta_0 \Gamma_0 + \beta_1 \Gamma_1)(A_0^{-1} f + \Pi \phi) = \beta_1 \Pi^* f + (\beta_0 + \beta_1 \Lambda)\phi$$

shows that the operator $\beta_0\Gamma_0 + \beta_1\Gamma_1$ is correctly defined on $A_0^{-1}H \dot{+} \Pi \text{dom } \Lambda \subset \text{dom } A$. Denoting $\mathcal{B} := \overline{\beta_0 + \beta_1\Lambda}$ with the domain $\text{dom } \mathcal{B} \supset \text{dom } \Lambda$, one checks that $H_{\mathcal{B}} := A_0^{-1}H \dot{+} \Pi \text{dom } \mathcal{B}$ is a Hilbert space with respect to the norm

$$\|u\|_{\mathcal{B}}^2 := \|f\|_H^2 + \|\phi\|_{\mathcal{H}}^2 + \|\mathcal{B}\phi\|_{\mathcal{H}}^2, \quad u = A_0^{-1}f + \Pi\phi.$$

It is then proved [24] (Lemma 4.1) that $\beta_0\Gamma_0 + \beta_1\Gamma_1$ extends to a bounded operator from $H_{\mathcal{B}}$ to \mathcal{H} . For the sake of convenience, same notation $\beta_0\Gamma_0 + \beta_1\Gamma_1$ is preserved for this extension.

We will make use of the following version of the celebrated Krein formula.

Proposition 2 ([24], Theorem 5.1). *Let $z \in \rho(A_0)$ be such that the operator $\overline{\beta_0 + \beta_1M(z)}$ defined on $\text{dom } \mathcal{B}$ is boundedly invertible. Then*

$$R_{\beta_0, \beta_1}(z) := (A_0 - z)^{-1} + S_z Q_{\beta_0, \beta_1}(z) S_z^*, \quad \text{where } Q_{\beta_0, \beta_1} := -(\overline{\beta_0 + \beta_1M(z)})^{-1} \beta_1, \quad (6)$$

is the resolvent of a closed densely defined operator A_{β_0, β_1} with the domain

$$\text{dom } A_{\beta_0, \beta_1} = \{u \in H_{\mathcal{B}} | (\beta_0\Gamma_0 + \beta_1\Gamma_1)u = 0\} = \ker(\beta_0\Gamma_0 + \beta_1\Gamma_1).$$

In particular, the (self-adjoint) operator of the transmission problem (2), which corresponds to the choice $\beta_0 = 0, \beta_1 = I$, admits the following characterisation in terms of its resolvent:

$$R_{0, I}(z) = (A_0 - z)^{-1} - S_z M^{-1}(z) S_z^*. \quad (7)$$

In this case, one clearly has $H_{\mathcal{B}} = A_0^{-1}H \dot{+} \Pi \text{dom } \Lambda$ and $\text{dom } A_{0, I} = \{u \in H_{\mathcal{B}} | \Gamma_1 u = 0\}$, which, together with the discussion at the beginning of this section, yields $A_{0, I} = A_{\varepsilon}$.

We remark that the operators β_0 and β_1 above can be assumed z -dependent, as this change does not impact the corresponding proofs of [24]. In this case, however, the corresponding operator-function $R_{\beta_0, \beta_1}(z)$ is shown to be the resolvent of a z -dependent operator family. Within the self-adjoint setup of the present paper, $R_{\beta_0, \beta_1}(z)$ is guaranteed to represent a generalised resolvent in the sense of [38–40].

3. Auxiliary Estimates

In this section, we collect a number of auxiliary statements required in our proof of the main result.

We start with the analysis of the operators Π_V and S_z^V introduced in Section 2. First, we note that each of these operators admits a decomposition into an orthogonal sum over N vertex domains $\{Q_v\}$ of Q . It therefore suffices to consider a single vertex domain Q_v (we recall for readers' convenience that the volume of this domain is assumed to be decaying with $\varepsilon \rightarrow 0$). Its boundary ∂Q_v contains a disjoint set of straight segments belonging to the internal boundary Γ^ε , which are, in line with what has been said above, denoted as $\Gamma_{\varepsilon v}^\varepsilon$; the union of the latter is Γ_v^ε .

The decoupled operator A_0 has $L^2(Q_v)$ as its invariant subspace. We will denote by $A_0^{(v)}$ its self-adjoint restriction, $A_0^{(v)} := A_0|_{L^2(Q_v)}$. By construction, the operator $A_0^{(v)}$ is the Laplacian with the so-called Zaremba, or mixed Neumann–Dirichlet, boundary condition [17,41]. More precisely, it is subject to the Dirichlet boundary condition on Γ_v^ε and to Neumann boundary condition on its complement $\tilde{\Gamma}_v^\varepsilon$. Clearly, this operator is boundedly invertible; moreover, the following statement holds.

Proposition 3 (see [31,42]). *There exists a constant $C > 0$ such that for all ε one has*

$$\|(A_0^{(v)})^{-1}\| \leq C|Q_v| \begin{cases} |\log \varepsilon|, & d = 2, \\ \varepsilon^{2-d}, & d \geq 3 \end{cases} \asymp \varepsilon |\log \varepsilon|^\gamma,$$

where, as before, $\gamma = 1$ for $d = 2$ and $\gamma = 0$ otherwise.

Remark 1. The above proposition holds under more general conditions than those we impose. Namely, the domain Q_v is only required to be Lipschitz and no conditions whatsoever are imposed on the geometry of the set Γ_v^ε .

Next, we turn our attention to the solution operator $S_z^{(v)} := S_z^V|_{\Gamma_v^\varepsilon}$ and the corresponding harmonic lift $\Pi^{(v)} := \Pi_V|_{\Gamma_v^\varepsilon}$. The two are clearly related by the formula

$$S_z^{(v)} = (1 - z(A_0^{(v)})^{-1})^{-1}\Pi^{(v)}.$$

In order to bound the norm of $\Pi^{(v)}$, we can follow, e.g., the following approach. First, consider the corresponding Zaremba problem on Q_v^0 . We proceed by relating the norm of the corresponding Poisson operator to the least Steklov eigenvalue of the bi-Laplacian, following the blueprint of [43], based in turn on Fichera’s duality principle; see [44]. Since the boundary of Q_v^0 is non-smooth, in doing so we follow the generalisations developed in [45,46], with obvious modifications required when passing from the Dirichlet to Zaremba setup. The estimate for the said Steklov eigenvalue is then taken from the norm of the compact embedding of $H^2(Q_v^0)$ to the traces of normal derivatives on the contact plates; see, e.g., [47]. Rescaling back to Q_v , we obtain the following auxiliary result.

Lemma 1. *There exists $C > 0$ such that $\|\Pi^{(v)}\| \leq C$ for all ε .*

By Proposition 3, the above lemma yields the following estimate for the solution operator $S_z^{(v)}$.

Lemma 2. *For $\varepsilon \in (0, 1)$, uniformly in $z \in K_\sigma^\varepsilon$ (and, in particular, for $z \in K_\sigma$) one has*

$$S_z^{(v)} = (1 + O(|z|\varepsilon|\log \varepsilon|^\gamma))\Pi^{(v)} = \Pi^{(v)} + O(|z|\varepsilon|\log \varepsilon|^\gamma),$$

where the error bounds are understood in the uniform operator norm topology.

Our next step is the analysis of the “part” of the DN map Λ^V pertaining to the vertex domain Q_v . We will denote by Λ_v^V its self-adjoint restriction $\Lambda^V|_{L^2(\Gamma_v^\varepsilon)} =: \Lambda_v^V$.

First, we note that the spectrum of Λ_v^V (which can be termed as the Steklov spectrum of the sloshing problem pertaining to $A_0^{(v)}$; see [48]) is discrete and accumulates to negative infinity. The point $\lambda_1 = 0$ is the least (by absolute value) Steklov eigenvalue with $\psi_v = |\Gamma_v^\varepsilon|^{-1/2}\mathbf{1}_{\Gamma_v^\varepsilon}$ being the corresponding eigenvector. For the second eigenvalue λ_2 , one has the following estimate; see, e.g., [32] and references therein.

Lemma 3. *There exists $C > 0$ such that $|\lambda_2| \geq C\varepsilon^{-1}$.*

Introduce the N -dimensional orthogonal projection

$$\mathcal{P} := \sum_v \langle \cdot, \psi_v \rangle \psi_v,$$

define $\mathcal{P}_\perp := 1 - \mathcal{P}$ and consider the operator $\mathcal{P}_\perp M(z) \mathcal{P}_\perp$, which is well defined since $\mathcal{P} \text{ dom } \Lambda \subset \text{dom } \Lambda$. By a straightforward estimate for sesquilinear forms, see [11] (Section 3.2), and taking into account (4) applied to $M_V(z)$ and combined with Proposition 3 as well as Lemmata 1, 3, one has the following statement.

Lemma 4. *There exists $C > 0$ such that for all $\varepsilon \in (0, 1)$ one has*

$$\|(\mathcal{P}_\perp M(z) \mathcal{P}_\perp)^{-1}\| \leq C\varepsilon, \quad z \in K_\sigma^\varepsilon,$$

where the operator $\mathcal{P}_\perp M(z)\mathcal{P}_\perp$ is considered as a linear (unbounded) operator in $\mathcal{P}_\perp \mathcal{H}$.

We conclude this section by noting that, alternatively, one can derive the bound on the Poisson operator $\Pi^{(v)}$ claimed in Lemma 1 by employing the scaling property of the Dirichlet-to-Neumann map Λ_v^V , similar to the argument of [32] referenced above, combined with a standard estimate on the solutions to the classical Neumann problem.

4. Norm-Resolvent Asymptotics

We will make use of the Kreĭn formula (7) to obtain a norm-resolvent asymptotics of the family A_ε . In doing so, we will compute the asymptotics of $M^{-1}(z)$ based on a Schur–Frobenius-type inversion formula, having first rewritten $M(z)$ as a 2×2 operator matrix relative to the orthogonal decomposition of the Hilbert space $\mathcal{H} = \mathcal{P}\mathcal{H} \oplus \mathcal{P}_\perp \mathcal{H}$. In the study of operator matrices, we rely upon the material of [49]; see also references therein.

The operator $M(z)$ admits the block matrix representation

$$M(z) = \begin{pmatrix} \mathbb{A} & \mathbb{B} \\ \mathbb{E} & \mathbb{D} \end{pmatrix} \text{ with } \mathbb{A}, \mathbb{B}, \mathbb{E} \text{ bounded.}$$

For the inversion of $M(z)$, we then use the Schur–Frobenius inversion formula [49] (Theorem 2.3.3)

$$\begin{pmatrix} \mathbb{A} & \mathbb{B} \\ \mathbb{E} & \mathbb{D} \end{pmatrix}^{-1} = \begin{pmatrix} \mathbb{A}^{-1} + \overline{\mathbb{A}^{-1}\mathbb{B}\mathcal{S}^{-1}}\mathbb{E}\mathbb{A}^{-1} & -\overline{\mathbb{A}^{-1}\mathbb{B}\mathcal{S}^{-1}} \\ -\overline{\mathcal{S}^{-1}}\mathbb{E}\mathbb{A}^{-1} & \overline{\mathcal{S}^{-1}} \end{pmatrix} \text{ with } \mathcal{S} := \mathbb{D} - \mathbb{E}\mathbb{A}^{-1}\mathbb{B}. \tag{8}$$

Note that by Proposition 1, one has $\Im M(z) = (\Im z)S_z^*S_z$. Moreover, since $S_z = (1 - zA_0^{-1})^{-1}\Pi$, one has

$$S_z^*S_z = \Pi^*(1 - zA_0^{-1})^{-1}(1 - \bar{z}A_0^{-1})^{-1}\Pi,$$

and therefore, for some constants $c_1, c_2 > 0$,

$$\langle S_z^*S_z\mathcal{P}\phi, \mathcal{P}\phi \rangle_{\mathcal{H}} = \|(1 - \bar{z}A_0^{-1})^{-1}\Pi\mathcal{P}\phi\|^2 \geq c_1\|\Pi\mathcal{P}\phi\|^2 \geq c_1\|\Pi_V\mathcal{P}\phi\|^2 \geq c_2\|\mathcal{P}\phi\|_{\mathcal{H}}^2$$

for all $\phi \in \mathcal{H}, z \in K_\sigma^\varepsilon$, where we have used the fact that the operator A_0 is bounded below by a positive constant. It follows that $\mathbb{A}^{-1} = (\mathcal{P}M(z)\mathcal{P})^{-1}$ is bounded.

Proceeding exactly as in [11] based on the estimate provided by Lemma 4 which now reads

$$\|\mathbb{D}^{-1}\| \leq C\varepsilon,$$

we use $\mathcal{S}^{-1} = (I - \mathbb{D}^{-1}\mathbb{E}\mathbb{A}^{-1}\mathbb{B})^{-1}\mathbb{D}^{-1}$ to obtain $\mathcal{S}^{-1} = O(\varepsilon)$.

Returning to (8), one obtains

$$M(z)^{-1} = \begin{pmatrix} \mathbb{A} & \mathbb{B} \\ \mathbb{E} & \mathbb{D} \end{pmatrix}^{-1} = \begin{pmatrix} \mathbb{A}^{-1} & 0 \\ 0 & 0 \end{pmatrix} + O(\varepsilon) \tag{9}$$

with a uniform estimate for the remainder term. Comparing our result with (6) of Proposition 2 with $\beta_0 := \mathcal{P}_\perp$ and $\beta_1 := \mathcal{P}$, one arrives at the following.

Theorem 1. *There exists $C > 0$ such that for all $\varepsilon \in (0, 1)$ and $z \in K_\sigma^\varepsilon$ (in particular, for all $z \in K_\sigma$) one has the estimate*

$$\|(A_\varepsilon - z)^{-1} - (A_{\beta_0, \beta_1} - z)^{-1}\| \leq C|z|^2\varepsilon$$

for a universal constant C and $\beta_0 = \mathcal{P}_\perp, \beta_1 = \mathcal{P}$, where the operator A_{β_0, β_1} is defined in Proposition 2.

Proof. The proof is identical to that of [11] (Theorem 3.1); we include it here for the sake of completeness. For the resolvent $(A_\varepsilon - z)^{-1}$, the formula (7) is applicable, in which for $M(z)^{-1}$ we use (9). As for the resolvent $(A_{\mathcal{P}_\perp, \mathcal{P}} - z)^{-1}$, Proposition 2 with $\beta_0 = \mathcal{P}_\perp$, $\beta_1 = \mathcal{P}$ is clearly applicable. Moreover, for this choice of β_0, β_1 , the operator

$$Q_{\mathcal{P}_\perp, \mathcal{P}}(z) = -(\overline{\mathcal{P}_\perp + \mathcal{P}M(z)})^{-1}\mathcal{P}$$

in (6) is easily computable (e.g., by the Schur–Frobenius inversion formula of [49], see (8)). The mentioned computation is facilitated by the fact that $\mathcal{P}_\perp + \mathcal{P}M(z)$ is triangular ($\mathbb{A} = \mathcal{P}M(z)\mathcal{P}$, $\mathbb{B} = \mathcal{P}M(z)\mathcal{P}_\perp$, $\mathbb{E} = 0$, $\mathbb{D} = I$ in (8)) with respect to the decomposition $\mathcal{H} = \mathcal{P}\mathcal{H} \oplus \mathcal{P}_\perp\mathcal{H}$, yielding

$$Q_{\mathcal{P}_\perp, \mathcal{P}}(z) = -\mathcal{P}(\mathcal{P}M(z)\mathcal{P})^{-1}\mathcal{P}, \tag{10}$$

and the claim follows. \square

Already the estimate of Theorem 1 establishes norm-resolvent convergence of the family A_ε to an operator which by (10) is a relative (i.e., with respect to the difference of resolvents) finite-rank perturbation of the decoupled operator A_0 . However, in the case $z \in K_\sigma$, which we will assume henceforth, it is possible to obtain a further simplification of this answer, relating the leading-order asymptotic term to a self-adjoint operator on the limiting metric graph. This procedure follows the blueprint of our paper [11]. We next briefly outline the related argument. For the case of z not constrained to a compact, a similar argument yields a sequence of dimensionally reduced models, as mentioned in the Introduction.

Note first that $(A_0 - z)^{-1} = (A_0^V - z)^{-1} \oplus (A_0^E - z)^{-1}$ is easily analysed. Indeed, by Proposition 3 one has

$$(A_0^V - z)^{-1} = O(\varepsilon |\log \varepsilon|^\gamma). \tag{11}$$

Furthermore, the operator $(A_0^E - z)^{-1}$, by separation of variables, is $O(\varepsilon^2)$ -close to the Dirichlet Laplacian on the space

$$H_G := \oplus_v L^2([0, l_e] \times \mathbf{1}_\varepsilon),$$

where $\mathbf{1}_e := \varepsilon^{-(d-1)/2}\mathbf{1}$ is the normalised constant function in the variable transverse to the edge e .

We remark that this is the only place where we use the assumption about the geometric shape of the edge parts Q_e of the thin structure. This can be generalised to the setup of [1], allowing for curvature and non-uniform thickness, leading to Laplace–Beltrami operators on the edges of the limiting graph.

The operator $(A_0 - z)^{-1}$ is therefore close, uniformly in $z \in K_\sigma$, to an operator that is unitary equivalent to the resolvent of A_0^G , where A_0^G is the Dirichlet-decoupled graph Laplacian pertaining to the graph G . The related error estimates are the same as in (11). The finite-dimensional second term on the right-hand side of (10) is therefore expected to encode the matching conditions at the vertices of the limiting graph G . In order to see this, one passes over to the generalised resolvent $R_\varepsilon(z) := P_E(A_\varepsilon - z)^{-1}P_E$, which is shown to admit the following asymptotics.

Theorem 2. *The operator family $R_\varepsilon(z)$ admits the following asymptotics in the operator-norm topology for $z \in K_\sigma$:*

$$R_\varepsilon(z) - R_{\text{eff}}(z) = O(\varepsilon),$$

where $R_{\text{eff}}(z)$ is the solution operator for the following spectral BVP on the edge domain Q_E :

$$\begin{aligned} -\Delta u - zu &= f, & f &\in L^2(Q_E), \\ \beta_0(z)\Gamma_0^E u + \beta_1\Gamma_1^E u &= 0, \end{aligned} \tag{12}$$

with $\beta_0(z) = \mathcal{P}_\perp - \mathcal{P}B(z)\mathcal{P}$, $B(z) := -M_V(z)$ and $\beta_1 = \mathcal{P}$.

The boundary condition in (12) can be written in the more conventional form

$$\mathcal{P}_\perp u|_{\Gamma^\varepsilon} = 0, \quad \mathcal{P}\partial_n u = \mathcal{P}B(z)\mathcal{P}u|_{\Gamma^\varepsilon}.$$

Equivalently,

$$R_\varepsilon(z) - (A_{\mathcal{P}_\perp - \mathcal{P}B(z)\mathcal{P}, \mathcal{P}}^E - z)^{-1} = O(\varepsilon),$$

where $A_{\mathcal{P}_\perp - \mathcal{P}B(z)\mathcal{P}, \mathcal{P}}^E$, for any fixed z , is the operator in $L^2(Q_E)$ defined by Proposition 2 relative to the triple $(\mathcal{H}, \Pi_E, \Lambda^E)$, where the term “triple” is understood in the sense of [24]. This operator is maximal anti-dissipative for $z \in \mathbb{C}_+$ and maximal dissipative for $z \in \mathbb{C}_-$; see [40].

The proof of the theorem follows immediately from Theorem 1; see [11] (Theorem 3.6), together with the observation that

$$\mathcal{P}M(z)\mathcal{P} = \mathcal{P}M_E(z)\mathcal{P} + \mathcal{P}M_V(z)\mathcal{P}.$$

The next step of our argument is to introduce the truncated (reduced) boundary space $\check{\mathcal{H}}$ in order to make all the ingredients finite-dimensional. In view of clarity, in what follows we consistently supply the (finite-dimensional) “truncated” spaces and operators pertaining to them by the breve overscript.

We put $\check{\mathcal{H}} := \mathcal{P}\mathcal{H}$ (noting that in our setup $\check{\mathcal{H}}$ is N -dimensional, where N is the number of vertices; see Section 2). Introduce the truncated Poisson operator on $\check{\mathcal{H}}$ by $\check{\Pi}_E := \Pi_E|_{\check{\mathcal{H}}}$ and the truncated DN map $\check{\Lambda}^E := \mathcal{P}\Lambda^E|_{\check{\mathcal{H}}}$. Then, the following statement holds.

Proposition 4 ([11], Theorem 3.7). 1. The formula

$$R_{\text{eff}}(z) = (A_0^E - z)^{-1} - \check{S}_z^E (\check{M}_E(z) - \mathcal{P}B(z)\mathcal{P})^{-1} (\check{S}_z^E)^* \tag{13}$$

holds, where \check{S}_z^E is the solution operator of the problem

$$\begin{aligned} -\Delta u_\phi - zu_\phi &= 0, \quad u_\phi \in \text{dom } A_0^E \dot{+} \text{ran } \check{\Pi}_E, \\ \Gamma_0^E u_\phi &= \phi, \quad \phi \in \check{\mathcal{H}}, \end{aligned}$$

and \check{M}_E is the M -operator defined in accordance with (3), (5) relative to the triple $(\check{\mathcal{H}}, \check{\Pi}_E, \check{\Lambda}^E)$.

2. The “effective” generalised resolvent $R_{\text{eff}}(z)$ is represented as the generalised resolvent of the problem

$$\begin{aligned} -\Delta u - zu &= f, \quad f \in L^2(Q_E), \quad u \in \text{dom } A_0^E \dot{+} \text{ran } \check{\Pi}_E, \\ \mathcal{P}\partial_n u|_{\Gamma^\varepsilon} &= \mathcal{P}B(z)\mathcal{P}u|_{\Gamma^\varepsilon}. \end{aligned}$$

3. The triple $(\check{\mathcal{H}}, \check{\Gamma}_0^E, \check{\Gamma}_1^E)$ is the classical boundary triple [50,51] for the operator A_{max} defined by the differential expression $-\Delta$ on the domain $\text{dom } A_{\text{max}} = \text{dom } A_0^E \dot{+} \text{ran } \check{\Pi}_E$. Here, $\check{\Gamma}_0^E$ and $\check{\Gamma}_1^E$ are defined on $\text{dom } A_{\text{max}}$ as the operator of the boundary trace on Γ and $\mathcal{P}\partial_n u$, respectively.

We now consider the operator $\mathcal{P}B(z)\mathcal{P}$ in (13); since $B = -M_V$ by definition, we invoke the estimates derived in Section 3 to obtain

$$\mathcal{P}B\mathcal{P} = -\mathcal{P}\Lambda^V\mathcal{P} - z\mathcal{P}\Pi_V^*\Pi_V\mathcal{P} + O(\varepsilon|\log \varepsilon|^\gamma) = -z\check{\Pi}_V^*\check{\Pi}_V + O(\varepsilon|\log \varepsilon|^\gamma),$$

with a uniform estimate for the remainder term. Here, the truncated Poisson operator $\check{\Pi}_V$ is introduced as $\check{\Pi}_V := \Pi_V|_{\check{\mathcal{H}}}$ relative to the same truncated boundary space as above, $\check{\mathcal{H}} = \mathcal{P}\mathcal{H}$. As a result, we obtain

$$R_{\text{eff}}(z) - R_{\text{hom}}(z) = O(\varepsilon|\log \varepsilon|^\gamma),$$

with

$$R_{\text{hom}}(z) := (A_0^E - z)^{-1} - \check{S}_z^E(\check{M}_E(z) + z\check{\Gamma}_V^* \check{\Gamma}_V)^{-1} (\check{S}_z^E)^*. \tag{14}$$

By a classical result of [40] (see also [38,39]), the operator $R_{\text{eff}}(z)$ is a generalised resolvent, so it defines a z -dependent family of closed densely defined operators in $L^2(Q_E)$, which are maximal anti-dissipative for $z \in \mathbb{C}_+$ and maximal dissipative for $z \in \mathbb{C}_-$. Writing the resolvent $(A_\varepsilon - z)^{-1}$ in the matrix form relative to the orthogonal decomposition $L^2(Q) = P_E L^2(Q) \oplus P_V L^2(Q) = L^2(Q_E) \oplus L^2(Q_V)$ then yields the following result.

Theorem 3. *The resolvent $(A_\varepsilon - z)^{-1}$ admits the following asymptotics in the uniform operator-norm topology:*

$$(A_\varepsilon - z)^{-1} = \mathcal{R}_{\text{eff}}(z) + O(\varepsilon |\log \varepsilon|^\gamma),$$

where the operator $\mathcal{R}_{\text{eff}}(z)$ has the following representation relative to the decomposition $L^2(Q_E) \oplus L^2(Q_V)$:

$$\mathcal{R}_{\text{eff}}(z) = \begin{pmatrix} R_{\text{eff}}(z) & \left(\mathfrak{K}_{\bar{z}} [R_{\text{eff}}(\bar{z}) - (A_0^E - \bar{z})^{-1}] \right)^* \check{\Gamma}_V^* \\ \check{\Gamma}_V \mathfrak{K}_z [R_{\text{eff}}(z) - (A_0^E - z)^{-1}] & \check{\Gamma}_V \mathfrak{K}_z \left(\mathfrak{K}_{\bar{z}} [R_{\text{eff}}(\bar{z}) - (A_0^E - \bar{z})^{-1}] \right)^* \check{\Gamma}_V^* \end{pmatrix}. \tag{15}$$

Here, $\mathfrak{K}_z := \Gamma_0^E|_{\mathfrak{N}_z}$ with $\mathfrak{N}_z := \text{ran } S_z^E \mathcal{P}$, $z \in \mathbb{C}_\pm$, and the generalised resolvent $R_{\text{eff}}(z)$ is defined by (14).

The above theorem provides us with the simplest possible leading-order term $\mathcal{R}_{\text{eff}}(z)$ of the asymptotic expansion for $(A_\varepsilon - z)^{-1}$. However, it is not yet obvious whether it is the resolvent of some self-adjoint operator in the space $L^2(Q_E) \oplus \check{\Gamma}_V \mathcal{H} \subset L^2(Q)$. It turns out that this is indeed so, which is seen via the following explicit construction.

Put $L^2(G) := \oplus_e L^2(0, l_e)$, $H^2(G) := \oplus_e H^2(0, l_e)$. For all $u \in H^2(G)$, denote by u_{ev} the limit of $u_e(x) := u|_e(x)$ at the vertex v . Let $H_{\text{eff}} := L^2(G) \oplus \mathbb{C}^N$, and set

$$\text{dom } \mathcal{A}_{\text{hom}} = \{ (u, \beta)^\top \in H_{\text{eff}} : u \in H^2(G), u_{ev} = u_{e'v} =: u_v \text{ for any } v \text{ and } e, e' \text{ incident to } v, \text{ and } \beta = \kappa u_V \}, \tag{16}$$

where u_V is the N -dimensional vector of $\{u_v\}_{v \in V}$ and κ is the diagonal matrix

$$\kappa := \text{diag}\{|Q_v^0|^{1/2}\}. \tag{17}$$

The action of the operator is set by

$$\mathcal{A}_{\text{hom}} \begin{pmatrix} u \\ \beta \end{pmatrix} = \begin{pmatrix} -u'' \\ -\kappa^{-1} \partial_n u|_V \end{pmatrix}, \quad \begin{pmatrix} u \\ \beta \end{pmatrix} \in \text{dom } \mathcal{A}_{\text{hom}}, \tag{18}$$

where $\partial_n u|_V$ is the N -dimensional vector $\{\sum_{e \sim v} \partial_n u_e|_v\}_{v \in V}$, i.e., the vector for which each element is represented by the sum of edge-inward normal derivatives of the function u over all the edges incident to the vertex v . We write $e \sim v$ if and only if the edge e is incident to the vertex v .

The main result of the present work, which is obtained by computing explicitly the resolvent of (16)–(18) and comparing it with (15) (see details of a similar computation in [13]), is formulated next.

Theorem 4. *The resolvent $(A_\varepsilon - z)^{-1}$ admits the following estimate in the uniform operator norm topology, uniform in $z \in K_\sigma$:*

$$(A_\varepsilon - z)^{-1} - \Theta (\mathcal{A}_{\text{hom}} - z)^{-1} \Theta^* = O(\varepsilon |\log \varepsilon|^\gamma),$$

where Θ is a partial isometry from H_{eff} onto $L^2(Q)$, acting as follows:

- For every edge $e \in G, e = [0, l_e]$, it embeds $u \in H^2(e)$ into $L^2(Q_e)$ as $u(x) \times \varepsilon^{-(d-1)/2} \mathbf{1}(y)$, where y is the variable in the direction transverse to that of x ;
- For every vertex $v \in G$, it embeds the value u_v , i.e., the common value of $u \in H^2(G)$ at the vertex v , into $L^2(Q_v)$ as $\varepsilon^{-(d-1)/2} u_v \mathbf{1}$.

5. Analysis of Vertex Matching Conditions

In the present section, we continue our study of the operator (16)–(18) associated with an arbitrary metric graph G , with a view to analyse its spectral structure. We will show that the matching conditions at graph vertices associated with the spectral problem for the mentioned operator, albeit closely resembling δ -type conditions with coupling constants linear in the spectral parameter z , can be in fact represented (up to a unitary gauge) by δ' -type matching conditions for all $z \neq 0$ (while for $z = 0$ they coincide with the classical Kirchhoff condition). We will assume throughout that this graph contains no loops, in line with the assumptions imposed on the thin network studied above. We will further assume without loss of generality that the graph G is connected and that the matrix κ is invertible.

Since the operator \mathcal{A}_{hom} can be viewed as a self-adjoint out-of-space extension (see, e.g., [52] and references therein) of a symmetric differential operator on the metric graph G , it is amenable to the classical boundary triples theory; see [53,54]. We recall that for a closed and densely defined symmetric operator \mathcal{A} on a separable Hilbert space H with domain $\text{dom } \mathcal{A}$, a boundary triple is defined as follows.

Definition 2 ([55]). *A triple $(\mathcal{K}, \Gamma_0, \Gamma_1)$ consisting of an auxiliary Hilbert space \mathcal{K} and linear mappings Γ_0, Γ_1 defined everywhere on $\text{dom } \mathcal{A}^*$ is called a boundary triple for \mathcal{A}^* if the following conditions are satisfied:*

- (1) *The abstract Green’s formula is valid*

$$(\mathcal{A}^* \vec{u}, \vec{v})_H - (\vec{u}, \mathcal{A}^* \vec{v})_H = (\Gamma_1 \vec{u}, \Gamma_0 \vec{v})_{\mathcal{K}} - (\Gamma_0 \vec{u}, \Gamma_1 \vec{v})_{\mathcal{K}}, \quad \vec{u}, \vec{v} \in \text{dom } \mathcal{A}^* \quad (19)$$

- (2) *For any $Y_0, Y_1 \in \mathcal{K}$ there exist $\vec{u} \in \text{dom } \mathcal{A}^*$, such that $\Gamma_0 \vec{u} = Y_0, \Gamma_1 \vec{u} = Y_1$. In other words, the mapping $\vec{u} \mapsto \Gamma_0 \vec{u} \oplus \Gamma_1 \vec{u}$ from $\text{dom } \mathcal{A}^*$ to $\mathcal{K} \oplus \mathcal{K}$ is surjective.*

It can be shown (see [55]) that a boundary triple for \mathcal{A}^* exists, although it is not unique.

Definition 3. *Let $\mathcal{T} = (\mathcal{K}, \Gamma_0, \Gamma_1)$ be a boundary triple of \mathcal{A}^* . The Weyl function of \mathcal{A}^* corresponding to \mathcal{T} and denoted by $M(z), z \in \mathbb{C} \setminus \mathbb{R}$, is an analytic operator-function with a positive imaginary part for $z \in \mathbb{C}_+$ (i.e., an operator R -function) with values in the algebra of bounded operators on \mathcal{K} such that*

$$M(z) \Gamma_0 \vec{u} = \Gamma_1 \vec{u} \quad \forall \vec{u} \in \ker(\mathcal{A}^* - zI).$$

For $z \in \mathbb{C} \setminus \mathbb{R}$ one has $(M(z))^ = (M(\bar{z}))$ and $\Im(z) \Im(M(z)) > 0$.*

A comparison with the assertion 4 of Proposition 1 shows that the Weyl function $M(z)$ in the context of the boundary triples theory is intimately related to the object introduced in Definition 1. The overall setup leading to its construction is, however, different and is based on the explicit choice of the boundary operators Γ_0 and Γ_1 .

Definition 4. *An extension \mathcal{A} of a closed densely defined symmetric operator \mathcal{A} is called almost solvable and is denoted by $\mathcal{A} = A_B$, if there exist a boundary triple $(\mathcal{K}, \Gamma_0, \Gamma_1)$ for \mathcal{A}^* and a bounded operator $B : \mathcal{K} \rightarrow \mathcal{K}$ defined everywhere in \mathcal{K} such that*

$$\vec{u} \in \text{dom } A_B \iff \Gamma_1 \vec{u} = B \Gamma_0 \vec{u}.$$

This definition implies that $\text{dom } A_B \subset \text{dom } \mathcal{A}^*$ and A_B is a restriction of \mathcal{A}^* to the linear set $\text{dom } A_B := \{ \vec{u} \in \text{dom } \mathcal{A}^* : \Gamma_1 \vec{u} = B \Gamma_0 \vec{u} \}$. In this context, the operator B plays

the rôle of a parameter for the family of extensions $\{A_B \mid B : \mathcal{K} \rightarrow \mathcal{K}\}$. It can be shown (see [36] for references) that the resolvent set of A_B is non-empty (i.e., A_B is maximal), both A_B and $(A_B)^* = A_{B^*}$ are restrictions of \mathcal{A}^* to their domains, and A_B and B are self-adjoint or dissipative simultaneously.

Under the additional assumption that \mathcal{A} is *simple* (or, in other words, completely non-self-adjoint), that is, it has no reducing self-adjoint “parts”, the spectrum of A_B coincides, counting multiplicities, with the set of points $z_0 \in \mathbb{C}$ into which $(M(z_0) - B)^{-1}$ does not admit analytic continuation. In the general case, however, the spectrum is a union of the “zeroes” of the operator-valued function $M(z) - B$ introduced above and the spectrum of the self-adjoint “part” of the symmetric operator \mathcal{A} in its Wold decomposition [56].

Our immediate aim is to construct a convenient boundary triple for the operator \mathcal{A}_{eff} . In doing so, we rely upon the framework (in a particular case of a loop-graph with exactly one vertex) of the paper [57].

We define the symmetric operator \mathcal{A} as follows (cf. (16)):

$$\text{dom } \mathcal{A} = \left\{ (u, \beta)^\top \in H_{\text{eff}} : u \in H^2(G), u_{ev} = u_{e'v} =: u_v \text{ for any } v \right. \\ \left. \text{and } e, e' \text{ incident to } v, \partial_n u|_V = 0, \text{ and } \beta = \kappa u_V \right\},$$

where u_V is, as above, the N -dimensional vector of $\{u_v\}_{v \in V}$, κ is the diagonal matrix (17), and $\partial_n u|_V$ is the N -dimensional vector of $\{\sum_{e \sim v} \partial_n u_e|_v\}_{v \in V}$. The action of the operator \mathcal{A} is set by (18). The operator thus defined is clearly symmetric in $L^2(G) \oplus \mathbb{C}^N$; its adjoint \mathcal{A}^* is defined by the same expression (18) on the domain

$$\text{dom } \mathcal{A}^* = \left\{ (u, \beta)^\top \in H_{\text{eff}} : u \in H^2(G), u_{ev} = u_{e'v} =: u_v \text{ for any } v \right. \\ \left. \text{and } e, e' \text{ incident to } v \right\}.$$

We have the following lemma.

Lemma 5. *Let $\mathcal{K} = \mathbb{C}^N$, $\Gamma_0(u, \beta)^\top := \partial_n u|_V$, and $\Gamma_1(u, \beta)^\top := \kappa^{-1}\beta - u_V$. The triple $(\mathcal{K}, \Gamma_0, \Gamma_1)$ is a boundary triple for the operator \mathcal{A}^* . The operator \mathcal{A}_{eff} is a self-adjoint almost solvable extension of \mathcal{A} , corresponding to the matrix $B = 0$ with respect to this boundary triple.*

The *proof* of the above lemma is obtained via integration by parts; for details, see [57].

We shall further require two ordinary differential operators on the metric graph G together with their boundary triples. Consider A_{max}^δ to be the operator generated by the negative Laplacian $-\Delta$ on $L^2(G)$, defined on the domain

$$\text{dom } A_{\text{max}}^\delta = \{u \in H^2(G) : u_v := u_{ev} = u_{e'v} \text{ for any } v \text{ and } e, e' \text{ incident to } v\}.$$

In the paper [58], it is shown that this is a natural choice of a maximal operator if one seeks to consider the so-called δ -type matching conditions at the graph vertices, i.e., matching conditions of the type

$$u \text{ continuous at every vertex } v \text{ and } \partial_n u|_V = \eta^2 u_V, \tag{20}$$

where η is a diagonal matrix. The conditions (20) reduce to Kirchhoff, or standard, matching conditions under the choice $\eta = 0$. The natural boundary triple for A_{max}^δ is $(\mathcal{K}, \Gamma_0^\delta, \Gamma_1^\delta)$, where $\mathcal{K} = \mathbb{C}^N$, $\Gamma_0^\delta u = u_V$, and $\Gamma_1^\delta u = \partial_n u|_V$. The corresponding Weyl function (“ M -matrix”) admits the form [58] $M_\delta(z) := \{m_{v'v}^\delta\}_{v, v' \in V}$, where

$$m_{v'v}^\delta(k) = \begin{cases} -k \sum_{e \sim v} \cot kl_e, & v = v', \\ k \sum_{e \sim v, e \sim v'} (\sin kl_e)^{-1}, & v \neq v'. \end{cases} \tag{21}$$

Here, $k := \sqrt{z}$ is such that $\Im k \geq 0$.

Next, we consider the *magnetic* Laplacian on the graph G subject to δ' -type matching at the vertices. Namely, we assume that for all edges e the action of the operator is described as

$$-\left(\frac{d}{dx} + i\tau_e\right)^2, \tag{22}$$

where τ_e is an edgewise-constant magnetic potential. In order to introduce δ' -type matching, consider the co-normal derivatives

$$\partial_n^\tau u_e|_v = \begin{cases} u'_e|_v + i\tau_e u_e|_v & \text{if } v \text{ is the left endpoint of } e, \\ -(u'_e|_v + i\tau_e u_e|_v) & \text{otherwise.} \end{cases}$$

We say that the magnetic Laplacian on G is subject to δ' -type matching at the vertices if

$$\partial_n^\tau u|_v := \partial_n^\tau u_e|_v = \partial_n^\tau u_{e'}|_v \text{ for any } e, e' \sim v, \text{ and } \Sigma u|_V = \eta^2 \partial_n^\tau u|_V \quad \forall v \in V, \tag{23}$$

where $\partial_n^\tau u|_V$ and $\Sigma u|_V$ denote the vectors $\{\partial_n^\tau u|_v\}_{v \in V}$ and $\{\sum_{e \sim v} u_e|_v\}_{v \in V}$, respectively.

By an argument similar to that of [58], see also [13], one easily checks that $(\mathcal{K}, \hat{\Gamma}_0, \hat{\Gamma}_1)$ is a boundary triple for the magnetic Laplacian $A_{\max}^{\delta'}$ defined on

$$\text{dom } A_{\max}^{\delta'} = \{u \in H^2(G) : \partial_n^\tau u|_v := \partial_n^\tau u_e|_v = \partial_n^\tau u_{e'}|_v \text{ for any } v \text{ and } e, e' \text{ incident to } v\},$$

if $\mathcal{K} = \mathbb{C}^N$, $\hat{\Gamma}_0 u := \{\partial_n^\tau u|_v\}_{v \in V}$, and, finally, $\hat{\Gamma}_1 u := -\{\sum_{e \sim v} u_e|_v\}_{v \in V}$. The corresponding M -matrix $\hat{M}(z) := \{\hat{m}_{vv'}\}_{v, v' \in V}$ admits the form [58]

$$\hat{m}_{vv'}(k) = \begin{cases} -k^{-1} \sum_{e \sim v} \cot kl_e, & v = v', \\ -k^{-1} \sum_{e \sim v, e \sim v'} \exp(i\sigma_e(v, v')\tau_e l_e) (\sin kl_e)^{-1}, & v \neq v'. \end{cases}$$

Here $k = \sqrt{z}$ such that $\Im k \geq 0$ and $\sigma_e(v, v') = 1$ if e is directed from v to v' , $\sigma_e(v, v') = -1$ otherwise.

We will now fix the values of the magnetic potential as follows: $\tau_e := \pi/l_e$. The operator $A_{\max}^{\delta'}$ corresponding to this choice will be henceforth denoted by \hat{A}_{\max} . Its M -matrix \hat{M} relative to the triple $(\mathcal{K}, \hat{\Gamma}_0, \hat{\Gamma}_1)$ admits the form

$$\hat{m}_{vv'}(k) = \begin{cases} -k^{-1} \sum_{e \sim v} \cot kl_e, & v = v', \\ k^{-1} \sum_{e \sim v, e \sim v'} (\sin kl_e)^{-1}, & v \neq v', \end{cases}$$

which coincides with (21) up to the factor z^{-1} . We remark that the operator \hat{A}_{\max} and any of its self-adjoint restrictions can be unitary transformed into a regular (non-magnetic) Laplacian on the same graph G by a standard gauge transform; this will, however, be reflected in the corresponding change of matching conditions at the vertices. Motivated by applications to electromagnetic wave propagation [16], here we prefer to proceed with the magnetic setup.

Returning to the analysis of the operator \mathcal{A}_{eff} , we now have the following lemma.

Lemma 6. *Relative to the boundary triple of Lemma 5, the Weyl M -matrix of the operator \mathcal{A}_{eff} admits the form*

$$M(z) = -M_\delta^{-1}(z) - \frac{1}{z} \kappa^{-2} = -\frac{1}{z} (\hat{M}^{-1}(z) + \kappa^{-2}).$$

Proof. In view of Lemma 5, consider the vector $(u, \beta)^\top \in \text{dom } \mathcal{A}^*$ such that

$$\mathcal{A}^* \begin{pmatrix} u \\ \beta \end{pmatrix} = z \begin{pmatrix} u \\ \beta \end{pmatrix}.$$

By the definition of \mathcal{A}^* , one equivalently has $-u'' = zu$ on G and $-\kappa^{-1}\partial_n u|_V = z\beta$. Abbreviating $\Gamma_0(u, \beta)^\top =: \alpha$, one therefore has, in view of the definition of Γ_0 (see Lemma 5):

$$\beta = -\frac{1}{z}\kappa^{-1}\alpha.$$

Furthermore, taking now into account the definition of the second boundary operator Γ_1 , one has

$$\Gamma_1(u, \beta)^\top = \kappa^{-1}\beta - u_V = -\frac{1}{z}\kappa^{-2}\alpha - u_V.$$

Note that the function u , by the definition of \mathcal{A}^* , must be continuous at every $v \in V$ and therefore belongs to the domain of the operator A_{\max}^δ . Therefore, $u \in \ker(A_{\max}^\delta - z)$ and thus one has $\Gamma_1^\delta u = M_\delta(z)\Gamma_0^\delta u$. By construction, one has $\Gamma_0^\delta u = u|_V$ and $\Gamma_1^\delta u = \partial_n u|_V$, whence $u_V = M_\delta^{-1}(z)\alpha$. Ultimately,

$$M(z)\alpha = \Gamma_1(u, \beta)^\top = -\frac{1}{z}\kappa^{-2}\alpha - M_\delta^{-1}(z)\alpha,$$

as claimed. \square

Let $\mathcal{E}(\cdot)$ denote the orthogonal operator spectral measure of the self-adjoint operator \mathcal{A}_{eff} and $E(\cdot)$ the orthogonal spectral measure of the self-adjoint operator $\hat{A}_{-\kappa^2}$, where the latter is defined as the almost solvable extension of $(A_{\max}^{\delta'})^*$ corresponding to the parameterising operator $B = -\kappa^2$. In other words, it is the magnetic Laplacian on the graph G subject to the condition $\tau_e = \pi/l_e$ for all $e \in E$ and δ' -type matching conditions (23) at the graph vertices with $\eta = \kappa$. We have the following statement.

Theorem 5. *For any $c_0 > 0$, the operators $\mathcal{A}_{\text{eff}}\mathcal{E}(c_0, \infty)$ and $\hat{A}_{-\kappa^2}E(c_0, \infty)$ are unitary equivalent.*

Proof. In the generic case when \mathcal{A} is simple (in particular when all l_e are rationally independent), the claim follows immediately from Lemma 6. Indeed, in this case no positive z can be an eigenvalue of the operator \mathcal{A} and the same applies to the operator $(A_{\max}^{\delta'})^*$. Therefore, any reducing self-adjoint “part” of either symmetric operator can only be zero. In the general case, if $z_0 > 0$ is an eigenvalue of \mathcal{A} , then the corresponding eigenfunction solves $-u'' = z_0u$ on G subject to $-\kappa^{-1}\partial_n u|_V = z_0\beta$, $u_V = \kappa^{-1}\beta$, $\partial_n u|_V = 0$ and therefore $u_V = \partial_n u|_V = 0$. Thus, any eigenfunction must be of the form $c_e \sin \sqrt{z_0}x$ on each e ; moreover, one must also have $\sin \sqrt{z_0}l_e = 0$ for all e such that $c_e \neq 0$ and $\partial_n u|_V = 0$. For each of these, the function that is edgewise transformed as $c_e \sin \sqrt{z_0}x \mapsto c_e \exp(-i\tau_e x) \cos \sqrt{z_0}x$ is shown to be an eigenfunction of $(A_{\max}^{\delta'})^*$ corresponding to the eigenvalue z_0 . The same argument applied in the opposite direction completes the proof. \square

Remark 2. *The analysis of the unitary equivalence linking $\mathcal{A}_{\text{eff}}\mathcal{E}(c_0, \infty)$ and $\hat{A}_{-\kappa^2}E(c_0, \infty)$ is an exciting possible development from the point of view of classical functional analysis, since it appears to be a natural graph-based generalisation of the classical Hilbert transform. This can be seen, in particular, from the explicit calculation in the case of an infinite chain graph [14].*

The above theorem shows that the spectral analysis of the operator \mathcal{A}_{eff} in relation to its non-zero spectrum reduces to that of the magnetic graph Laplacian with (non-trivial) δ' -type matching condition at the graph vertices, which could come as a surprise given that the eigenvalue problem for \mathcal{A}_{eff} yields, for the first component of the eigenvector $(u, \beta)^\top$, the equation $-u'' = zu$ on G subject to $\partial_n u|_V = -z\kappa^2 u|_V$, which on the face of it is a δ -type matching condition, albeit with coupling constants proportional to the spectral parameter z . It would seem natural, therefore, for the self-adjoint operator \mathcal{A}_{eff} to be a H^{-1} singular perturbation of the graph Laplacian with standard boundary conditions. Instead, our result shows that it is in fact a more singular H^{-2} perturbation. The underlying reasons of this peculiar behaviour are discussed in detail in [59], where the relationship of operators of

the class considered with those in the area of zero-range potentials with internal structure, as introduced by B. S. Pavlov [60,61], is explained. We also point out that the assertion of Theorem 5 has been observed in a particular case of the cycle graph with one vertex (the loop) in [14] in the context of a high-contrast homogenisation problem on the real line.

We conclude the spectral analysis of the operator \mathcal{A}_{eff} by a brief discussion of its kernel and the comparison of the latter with that of the operator $\hat{A}_{-\kappa^2}$. It turns out that, unlike what happens with its non-zero spectrum, the kernel $\ker \mathcal{A}_{\text{eff}}$ of the operator \mathcal{A}_{eff} is that of the Kirchhoff graph Laplacian.

Note first that $\ker \mathcal{A}_{\text{eff}}$ necessarily belongs to the non-simple (i.e., self-adjoint) part of \mathcal{A} . Indeed, for $(u, \beta)^\top \in \ker \mathcal{A}_{\text{eff}}$ it has to satisfy

$$-u'' = 0 \text{ on } G, \quad \partial_n u|_V = 0, \quad u_V = \kappa^{-1}\beta,$$

where the boundary conditions are equivalent to $\Gamma_0(u, \beta)^\top = 0, \Gamma_1(u, \beta)^\top = 0$. On the other hand, this is precisely the condition for u to be in the kernel of a graph Laplacian with Kirchhoff matching conditions. One therefore infers from [5] that $\dim \ker \mathcal{A}_{\text{eff}} = 1$ (in our case of connected graphs), and the elements of $\ker \mathcal{A}_{\text{eff}}$ are constants on G .

The kernel of $\hat{A}_{-\kappa^2}$ is spanned by functions u such that

$$-u'' = 0 \text{ on } G \text{ and } \Sigma u|_V = \kappa^2 \partial_n^\tau u|_V.$$

It is easily checked that a non-trivial solution to this problem could exist only if $\partial_n^\tau u|_V = 0, \Sigma u|_V = 0$ or, in other words, if $\hat{\Gamma}_0 u = \hat{\Gamma}_1 u = 0$. This means that, precisely as in the case of \mathcal{A}_{eff} , the kernel of $\hat{A}_{-\kappa^2}$ necessarily belongs to the non-simple (i.e., self-adjoint) part of the symmetric operator $(A_{\text{max}}^{\delta'})^*$. The question of its existence and dimension admits a simple answer in terms of the graph topology. It is clear that it is trivial in the case when G is a tree; in general, its dimension is shown to be equal to the cyclomatic number χ of the graph G . In particular, this yields unitary equivalence of \mathcal{A}_{eff} and $\hat{A}_{-\kappa^2}$ in the case where G contains exactly one cycle.

Remark 3. *The result of the present section seems to have been overlooked in a number of non-classical papers dealing with Sturm–Liouville problems on an interval with boundary conditions depending on a spectral parameter; see, e.g., [62–67].*

Conjecture 1. *The above discussion raises the question of which definition of δ' -type interaction on a graph is motivated physically, i.e., whether it is the one emerging from the analysis of thin networks as the operator \mathcal{A}_{hom} , see (16)–(18), or the traditional (see [5]) definition (22)–(23). At first sight, the difference between the two operators is insignificant: it is only in their kernels. However, it can happen to be of paramount importance if, e.g., one considers an ε -periodic graph with δ' -type matching conditions, in which case the homogenisation procedure [11,12] will lead to drastically different outcomes for the two related setups, as it relies upon a “threshold effect” [68] in the behaviour of the least eigenvalue of the operator on the fundamental cell for small quasimomenta.*

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References

1. Post, O. *Spectral Analysis on Graph-Like Spaces*; Lecture Notes in Mathematics 2039; Springer: Berlin/Heidelberg, Germany, 2012.
2. Exner, P.; Post, O. Convergence of spectra of graph-like thin manifolds. *J. Geom. Phys.* **2005**, *54*, 77–115. [\[CrossRef\]](#)
3. Kuchment, P.; Zeng, H. Convergence of spectra of mesoscopic systems collapsing onto a graph. *J. Math. Anal. Appl.* **2001**, *258*, 671–700. [\[CrossRef\]](#)
4. Kuchment, P.; Zeng, H. Asymptotics of spectra of Neumann Laplacians in thin domains. *Contemp. Math.* **2003**, *327*, 199–213.
5. Berkolaiko, G.; Kuchment, P. Introduction to Quantum Graphs; In *Mathematical Surveys and Monographs*; American Mathematical Society: Providence, RI, USA, 2012.
6. Mikhailova, A.B.; Pavlov, B.S.; Prokhorov, L.V. Intermediate Hamiltonian via Glazman’s splitting and analytic perturbation for meromorphic matrix-functions. *Math. Nachr.* **2007**, *280*, 1376–1416. [\[CrossRef\]](#)
7. Pavlov, B.S. Krein formula with compensated singularities for the ND-mapping and the generalized Kirchhoff condition at the Neumann Schrödinger junction. *Russ. J. Math. Phys.* **2008**, *15*, 364–388. [\[CrossRef\]](#)
8. Kadic, M.; Bückmann, T.; Stenger, N.; Thiel, M.; Wegener, M. On the practicability of pentamode mechanical metamaterials. *Appl. Phys. Lett.* **2012**, *100*, 191901. [\[CrossRef\]](#)
9. Milton, G.W.; Cherkaev, A.V. Which Elasticity Tensors are Realizable? *J. Eng. Mater. Technol.* **1995**, *117*, 483. [\[CrossRef\]](#)
10. Cherednichenko, K.; Kiselev, A.V.; Velčić, I.; Žubrinić, J. Effective behaviour of critical-contrast PDEs: Micro-resonances, frequency conversion, and time dispersive properties. II. *arXiv* **2023**, arXiv:2307.01125.
11. Cherednichenko, K.D.; Ershova, Y.Y.; Kiselev, A.V. Effective behaviour of critical-contrast PDEs: Micro-resonances, frequency conversion, and time dispersive properties. I. *Commun. Math. Phys.* **2020**, *375*, 1833–1884. [\[CrossRef\]](#)
12. Cherednichenko, K.D.; Ershova, Y.Y.; Kiselev, A.V. Time-dispersive behaviour as a feature of critical contrast media. *SIAM J. Appl. Math.* **2019**, *79*, 690–715. [\[CrossRef\]](#)
13. Cherednichenko, K.; Ershova, Y.; Kiselev, A.; Naboko, S. Unified approach to critical-contrast homogenisation with explicit links to time-dispersive media. *Trans. Mosc. Math. Soc.* **2019**, *80*, 295–342. [\[CrossRef\]](#)
14. Cherednichenko, K.D.; Kiselev, A.V. Norm-resolvent convergence of one-dimensional high-contrast periodic problems to a Kronig-Penney dipole-type model. *Commun. Math. Phys.* **2017**, *349*, 441–480. [\[CrossRef\]](#)
15. Rubinstein, J.; Schatzman, M. Variational Problems on multiply connected thin strips I: Basic estimates and convergence of the Laplacian spectrum. *Arch. Ration. Mech. Anal.* **2001**, *160*, 271–308. [\[CrossRef\]](#)
16. Kiselev, A.V.; Ryadovkin, K.S. Phase transition in a periodic tubular structure. *arXiv* **2023**, arXiv:2303.00872; to appear in: *SIAM J. Appl. Math.* (2024).
17. Zarembo, S. Sur un problème mixte relatif à l’équation de Laplace (French). In *Bulletin International de l’Académie des Sciences de Cracovie. Classe des Sciences Mathématiques et Naturelles, Serie A: Sciences Mathématiques*; Uniwersytet Jagielloński: Krakow, Poland, 1910; pp. 313–344.
18. Brown, R.M. The mixed problem for Laplace’s equation in a class of Lipschitz domains. *Commun. Partial Diff. Eqns.* **1994**, *19*, 1217–1233. [\[CrossRef\]](#)
19. Taylor, J.L.; Ott, K.A.; Brown, R.M. The mixed problem in Lipschitz domains with general decompositions of the boundary. *Trans. Am. Math. Soc.* **2013**, *365*, 2895–2930. [\[CrossRef\]](#)
20. Grieser, D. Spectra of graph neighborhoods and scattering. *Proc. Lond. Math. Soc.* **2008**, *97*, 718–752. [\[CrossRef\]](#)
21. Birman, M.S. On the averaging procedure for periodic operators in a neighborhood of an edge of an internal gap. *St. Petersburg Math. J.* **2004**, *15*, 507–513. [\[CrossRef\]](#)
22. Craster, R.V.; Kaplunov, J.; Pichugin, A.V. High-frequency homogenization for periodic media. *Proc. R. Soc. Lond. Ser. A* **2010**, *466*, 2341–2362. [\[CrossRef\]](#)
23. Friedlander, L. On the density of states of periodic media in the large coupling limit. *Commun. Part. Diff. Eq.* **2002**, *27*, 355–380. [\[CrossRef\]](#)
24. Ryzhov, V. Linear operators and operator functions associated with spectral boundary value problems. *Oper. Adv. Appl.* **2020**, *276*, 576–626.
25. Schechter, M. A generalization of the problem of transmission. *Ann. Scuola Norm. Sup. Pisa* **1960**, *14*, 207–236.
26. Behrndt, J.; Langer, M. Boundary value problems for elliptic partial differential operators on bounded domains. *J. Func. Anal.* **2007**, *243*, 536–565. [\[CrossRef\]](#)
27. Brown, M.; Marletta, M.; Naboko, S.; Wood, I. Boundary triples and M -functions for non-selfadjoint operators with applications to elliptic PDEs and block operator matrices. *J. Lond. Math. Soc.* **2008**, *77*, 700–718. [\[CrossRef\]](#)
28. Birman, M.S. On the self-adjoint extensions of positive definite operators. *Math. Sb.* **1956**, *38*, 431–450.
29. Kreĭn, M.G. Theory of self-adjoint extensions of semibounded hermitian operators and applications II. *Mat. Sb.* **1947**, *21*, 365–404.
30. Višik, M.I. On general boundary problems for elliptic differential equations (Russian). *Trudy Moskov. Mat. Obšč.* **1952**, *1*, 187–246.
31. Denzler, J. Bounds for the heat diffusion through windows of given area. *J. Math. Anal. Appl.* **1998**, *217*, 405–422. [\[CrossRef\]](#)
32. Agranovich, M.S. On a mixed Poincaré-Steklov type spectral problem in a Lipschitz domain. *Russ. J. Math. Phys.* **2006**, *13*, 239–244. [\[CrossRef\]](#)
33. Grubb, G. The mixed boundary value problem, Krein resolvent formulas and spectral asymptotic estimates. *J. Math. Anal. Appl.* **2011**, *382*, 339–363. [\[CrossRef\]](#)

34. Ryzhov, V. Weyl-Titchmarsh function of an abstract boundary value problem, operator colligations, and linear systems with boundary control. *Complex Anal. Oper. Theory* **2009**, *3*, 289–322. [[CrossRef](#)]
35. Behrndt, J.; Rohleder, J. Spectral analysis of selfadjoint elliptic differential operators, Dirichlet-to-Neumann maps, and abstract Weyl functions. *Adv. Math.* **2015**, *285*, 1301–1338. [[CrossRef](#)]
36. Cherednichenko, K.D.; Kiselev, A.V.; Silva, L.O. Scattering theory for non-selfadjoint extensions of symmetric operators. *Oper. Theory Adv. Appl.* **2020**, *276*, 194–230.
37. Abels, H.; Grubb, G.; Wood, I.G. Extension theory and Kreĭn-type resolvent formulas for non-smooth boundary value problems. *J. Func. Anal.* **2014**, *266*, 4037–4100. [[CrossRef](#)]
38. Neumark, M. Spectral functions of a symmetric operator (Russian). *Bull. Acad. Sci. URSS. Sér. Math.* **1940**, *4*, 277–318.
39. Neumark, M. Positive definite operator functions on a commutative group (Russian). *Bull. Acad. Sci. URSS Sér. Math.* **1943**, *7*, 237–244.
40. Strauss, A.V. Generalised resolvents of symmetric operators (Russian). *Izv. Akad. Nauk SSSR Ser. Mat.* **1954**, *18*, 51–86.
41. Girouard, A.; Polterovich, I. Spectral geometry of the Steklov problem (survey article). *J. Spec. Theory* **2017**, *7*, 321–359. [[CrossRef](#)] [[PubMed](#)]
42. Denzler, J. Windows of given area with minimal heat diffusion. *Trans. Am. Math. Soc.* **1999**, *351*, 569–580. [[CrossRef](#)]
43. Kuttler, J.R. Remarks on a Stekloff eigenvalue problem. *SIAM J. Numer. Anal.* **1972**, *9*, 1–5. [[CrossRef](#)]
44. Fichera, G. Su un principio di dualità per talune formole di maggiorazione relative alle equazioni differenziali (Italian). *Atti Accad. Naz. Lincei Rend. Cl. Sci. Fis. Mat. Nat.* **1956**, *19*, 411–418.
45. Auchmuty, G. The S.V.D. of the Poisson kernel. *J. Fourier Anal. Appl.* **2017**, *23*, 1517–1536. [[CrossRef](#)]
46. Bucur, D.; Ferrero, A.; Gazzola, F. On the first eigenvalue of a fourth order Steklov problem. *Calc. Var.* **2009**, *35*, 103–131. [[CrossRef](#)]
47. Ferrero, A.; Gazzola, F.; Weth, T. On a fourth order Steklov eigenvalue problem. *Analysis* **2005**, *25*, 315–332. [[CrossRef](#)]
48. Levitin M.; Parnovski, L.; Polterovich, I.; Scher, D. Sloshing, Steklov and corners: Asymptotics of sloshing eigenvalues. *J. Anal. Math.* **2022**, *146*, 65–125. [[CrossRef](#)]
49. Tretter, C. *Spectral Theory of Block Operator Matrices and Applications*; Imperial College Press: London, UK, 2008.
50. Derkach, V.A.; Malamud M.M. Generalised resolvents and the boundary value problems for Hermitian operators with gaps. *J. Funct. Anal.* **1991**, *95*, 1–95. [[CrossRef](#)]
51. Gorbachuk, V.I.; Gorbachuk, M.L. *Boundary Value Problems for Operator Differential Equations*; Kluwer Academic Publishers: Dordrecht, The Netherlands, 1991.
52. Štraus, A.V. Functional models and generalized spectral functions of symmetric operators. *St. Petersburg Math. J.* **1999**, *10*, 733–784.
53. Behrndt, J.; Hassi, S.; de Snoo, H. *Boundary Value Problems, Weyl Functions, and Differential Operators*; Springer: Berlin/Heidelberg, Germany, 2020.
54. Schmüdgen, K. *Unbounded Self-Adjoint Operators on Hilbert Space*; Graduate Texts in Mathematics 265; Springer: Dordrecht, The Netherlands, 2012.
55. Kočubeĭ, A.N. Extensions of symmetric operators and of symmetric binary relations. *Math. Notes* **1975**, *17*, 25–28. [[CrossRef](#)]
56. Gorbachuk, M.I.; Gorbachuk, V.I. *M.G. Kreĭn's Lectures on Entire Operators*; Springer: Dordrecht, The Netherlands, 2012.
57. Cherednichenko, K.; Ershova, Y.; Naboko, S. Functional model for generalised resolvents and its application to time-dispersive media. *arXiv* **2021**, arXiv:2111.05387.
58. Ershova, Y.; Karpenko, I.I.; Kiselev, A.V. Isospectrality for graph Laplacians under the change of coupling at graph vertices. *J. Spectr. Theory* **2016**, *6*, 43–66. [[CrossRef](#)]
59. Cherednichenko, K.D.; Ershova, Y.Y.; Kiselev, A.V.; Ryzhov, V.A.; Silva, L.O. Asymptotic analysis of operator families and applications to resonant media. *Oper. Theory Adv. Appl.* **2023**, *291*, 239–311.
60. Pavlov, B.S. A model of zero-radius potential with internal structure (Russian). *Teoret. Mat. Fiz.* **1984**, *59*, 345–353.
61. Pavlov, B.S. The theory of extensions, and explicitly soluble models. *Russ. Math. Surv.* **1987**, *42*, 127–168. [[CrossRef](#)]
62. Evans, W.D. A non-selfadjoint differential operator in $L^2_{\sigma}[a, b]$. *Quart. J. Math.* **1970**, *21*, 371–383. [[CrossRef](#)]
63. Fulton, C.T. Two-point boundary value problems with eigenvalue parameter contained in the boundary conditions. *Proc. R. Soc. Edinb A.* **1977**, *77*, 293–308. [[CrossRef](#)]
64. Hinton, D.B. An expansion theorem for an eigenvalue problem in the boundary conditions. *Quart. J. Math.* **1977**, *30*, 33–42. [[CrossRef](#)]
65. Schneider, A. A note on eigenvalue problems with eigenvalue parameter in the boundary conditions. *Math. Z.* **1974**, *136*, 163–167. [[CrossRef](#)]
66. Shkalikov, A.A. Boundary problems for ordinary differential equations with parameter in the boundary conditions. *J. Soviet. Math.* **1983**, *33*, 1311–1342. [[CrossRef](#)]
67. Walter, J. Regular eigenvalue problems with eigenvalue parameter in the boundary condition. *Math. Z.* **1973**, *133*, 301–312. [[CrossRef](#)]
68. Birman, M.S.; Suslina, T.A. Second order periodic differential operators. Threshold properties and homogenisation. *St. Petersburg Math. J.* **2004**, *15*, 639–714. [[CrossRef](#)]

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