



2D MHD Mixed Convection in a Zigzag Trapezoidal Thermal Energy Storage System Using NEPCM

Aissa Abderrahmane ¹, Obai Younis ², Mohammad Al-Khaleel ^{3,4,*}, Houssem Laidoudi ⁵, Nevzat Akkurt ⁶, Kamel Guedri ⁷ and Riadh Marzouki ^{8,9}

- ¹ Laboratoire de Physique Quantique de la Matière et Modélisation Mathématique (LPQ3M), University of Mascara, Mascara 29000, Algeria
- ² Department of Mechanical Engineering, College of Engineering in Wadi Addwasir, Prince Sattam Bin Abdulaziz University, Al-Kharj 11942, Saudi Arabia
- ³ Department of Mathematics, Khalifa University, Abu Dhabi 127788, United Arab Emirates
- ⁴ Department of Mathematics, Yarmouk University, Irbid 21163, Jordan
- ⁵ Laboratory of Sciences and Marine Engineering (LSIM), Oran 31000, Algeria
- ⁶ Department of Mechanical Engineering, Munzur University, 62000 Tunceli, Turkey
- ⁷ Mechanical Engineering Department, College of Engineering and Islamic Architecture, Umm Al-Qura University, P.O. Box 5555, Makkah 21955, Saudi Arabia
- ⁸ Chemistry Department, College of Science, King Khalid University, Abha 61413, Saudi Arabia
- ⁹ Chemistry Department, Faculty of Sciences of Sfax, University of Sfax, Sfax 3038, Tunisia
 - * Correspondence: mohammad.alkhaleel@ku.ac.ae

Abstract: In a magnetic field, two-dimensional (2D) mixed convection is investigated within a zigzagged trapezoidal chamber. The lower side of the trapezoidal chamber is irregular, in particular, a zigzagged wall with different zigzag numbers N. The fluid particles move in the room due to the motion of the upper wall, while the porosity-enthalpy approach represents the melting process. The thermal parameters of the fluid are enhanced by what is called a nano-encapsulated phase change material (NEPCM) consisting of polyurethane as the shell and a nonadecane as the core, while water is used as the base fluid. In order to treat the governing equations, the well-known Galerkin finite element method (GFEM) is applied. In addition, the heat transfer (HT) irreversibility and the fluid friction (FF) irreversibility are compared in terms of the average Bejan number. The main results show that the melt band curve behaves parabolically at smaller values of Reynolds number (Re) and larger values of Hartmann number (Ha). Moreover, minimizing the wave number is better in order to obtain a higher heat transfer rate.

Keywords: magneto-hydrodynamics; NEPCM; 2D flow; melting process; GFEM; zigzag trapezoidal vessels

1. Introduction

Phase-change materials (PCM) in thermal systems are gaining tremendous attraction for energy storage systems, heat exchangers, air-conditioning systems, green buildings, etc. The advantage of the PCM is that the energy is saved in the materials as latent energy during the phase transition process, resulting in the temperature being almost constant in the system. This process is also referred to as latent heat storage material, which is carried out in evaporators in cooling systems [1]. From the literature, it has been found that the PC properties of the different materials vary due to temperature variation. On behalf of the construction of the multiscale structure, the PCMs are categorized as organic, inorganic, and eutectics [2]. Cui et al. [3] reviewed the PCM to understand its application in building better. They found the best phase change material applications when the room air temperature was as low as 4.2 °C. Aftab et al. [4] used PCM-integrated latent heat storage systems for sustainable energy advancements. Choudhari et al. [5] employed a fin structure to develop a PCM battery thermal management system. Numerous pieces of research



Citation: Abderrahmane, A.; Younis, O.; Al-Khaleel, M.; Laidoudi, H.; Akkurt, N.; Guedri, K.; Marzouki, R. 2D MHD Mixed Convection in a Zigzag Trapezoidal Thermal Energy Storage System Using NEPCM. *Nanomaterials* **2022**, *12*, 3270. https://doi.org/10.3390/ nano12193270

Academic Editors: Florian Ion Tiberiu Petrescu and Gang Shi

Received: 25 August 2022 Accepted: 15 September 2022 Published: 20 September 2022

Publisher's Note: MDPI stays neutral with regard to jurisdictional claims in published maps and institutional affiliations.



Copyright: © 2022 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). have been communicated using the PCM to upgrade the battery thermal management systems [6–10].

It is challenging to implement PCM in a two-phase state, in particular, force convection heat transfer in the solid-liquid phase. Miansari et al. [11] addressed the fluid motion in the melting process and illustrated that the fluid motion is high in the thicker tank, while motion became slower in the thinner tank. Hosseinizadeh et al. [12] encapsulated nanoenhanced phase change materials (NEPCMs) inside the core and shell of paraffin. They found that the NEPCM has a higher impact when used in the energy storage system compared to PCM. Alehosseini et al. [13] illustrated that NEPCMs have thermal and mechanical stability, which make better phase-change materials. Some researchers have seen the thermal properties of the fluid while using NEPCM and found that the heat transfer rate was enhanced by 23.6% [14]. Ghalambaz et al. [15,16] simulated the mixed and natural convection heat transfer through an inclined porous cavity in the presence of NEPCM. Nadezhda et al. [17] investigated the entropy generation of the partially triangular open cavity in the presence of natural convection of a nanofluid. Seyf et al. [18] discussed the heat transfer in the microtube by using the phase-change heat transfer of NEPCMs. For the phase changer, many studies have extensively targeted heat transfer in tubes, micro conduits, and ducts. Varol [19] and Alshuraiaan [20] investigated the free convective heat transfer in a trapezoidal cavity without considering the NEPCM. Hossain et al. [21] investigated the magnetohydrodynamic (MHD) natural convective flow in a trapezoidal cavity filled by porous medium using the Galerkin weighted residual numerical approach (GWRNA). They obtained strong convection flow for high Ra (Rayleigh number) and Da (Darcy number) at Ha = 50 (Hartmann number). In the present model, we have considered the MHD mixed convection in the trapezoidal cavity for the dispersion of chemical contaminants through the fluids flow. Typically, cavities are carried out in industries such as triangular, square, trapezoidal, etc.

Khalil et al. [22] optimized the MHD free convective in the porous trapezoidal cavity using a multi-objective optimum technique. They found that heat transfer and energy enhancements (HTE and EE) are increased as the number of bottom wall waves increases. Hussain et al. [23] studied the entropy generation of natural convective flow inside the sinusoidal corrugated in the presence of the magnetic field. Extensive research and experimental works have been reported on both mixed and natural convective heat transfer in different cavities under different physical effects for different applications [24–30]. Few stabilization studies have been introduced for the heat transfer process in the cavities models. Galerkin Finite Element (GFE) is one of the most useful methods evolved from the FEM. For illustration, Chen et al. [31] used a class of the Galerkin method in order to analyze the two-dimensional natural convective porous medium flow. Al-Kouz et al. [32] studied entropy optimization of MHD convection in a cavity of odd shape using the GFE method. Generally, the usage of NEPCM in heat transfer applications results in enhancing heart transfer rates; however, special attention must be paid regarding nanoparticle accumulation as it would result in negative impacts, as reported by Qiu et al. [33] and Feng et al. [34].

According to the above literature survey, no literature has been reported regarding the MHD mixed convective flow within a trapezoidal cavity using the GFE method. In the present article, NEPCM is considered in the zigzag trapezoidal thermal storage system in order to enhance heat transfer and energy storage. The Galerkin Finite Element (GFEM) is applied to simulate the model's numerical solutions. In addition, a comparative study of the heat transfer irreversibility and the fluid friction irreversibility in terms of the average Bejan number has been presented. The outcomes of the model will be applicable in various energy systems, particularly for the storage of energy using the NEPCMs.

The current numerical work aims to respond to the following queries in the problem:

- 1. Does the velocity of the flow increase with Re along with the temperature distribution?
- 2. What is the role of the zigzag number on the flow velocity and the heat capacity?
- 3. Does increasing the value of Ha affect the heat transfer inside the cavity?

2. Formulation, Properties, and Problem Settings

The geometry consists of a trapezoidal prism cavity (see Figure 1). The upper cavity wall is moving at a constant velocity of U, while the rest of the walls are stationary. Mixed convection of liquid-NEPCM takes place in the cavity due to movement of the upper wall (forced convection) and the temperature difference between side walls and the lower zigzag wall (natural convection). The side walls of the chamber have a constant temperature of T_c , whereas the temperature of the bottom zigzag wall is given by a constant temperature of T_h . The PCM serves as the core of the nano-encapsulated PCM, while the encapsulating material serves as the shell. At the fusion temperature, the core goes through a phase shift. Within its fusion temperature, the core of NEPCM nanoparticles can undergo a phase shift to a liquid state and absorb or release energy in the form of latent heat.



Figure 1. Physical model geometry.

2.1. Mathematical Modeling and Equations

As mentioned before, the buoyancy force (due to the temperature difference of the walls) and the moving upper wall produce the natural and forced convection currents, respectively. As a result of such a combination, one gets mixed convection. It is necessary to make some assumptions in order to predict the heat transfer behavior of NEPCMs when they are exposed to mixed convection. In this context, it is assumed that the distribution of NEPCMs throughout the combination of a base fluid and NEPCMs is consistent and stable. The thermophysical characteristics are unaffected by temperature in any way, with the exception of the density, which may be approximated using the Boussinesq approximation.

In addition, it is presumed that the nanoparticles and the base fluid are in a state of local thermal equilibrium that the flow is steady, two-dimensional, and incompressible.

In the sequel, we present the continuity, momentum, and energy equations in the x-y coordinates [35]:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \tag{1}$$

$$\rho_m \left(u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} \right) = -\frac{\partial P}{\partial x} + \mu_m \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right)$$
(2)

$$\rho_m \left(u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} \right) = -\frac{\partial P}{\partial y} + \mu_m \left(\frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} \right) + (\rho \beta)_m (T - T_C)g - \frac{\sigma_{nf} B_0^2 v}{\rho_{nf}}$$
(3)

$$\left(u\frac{\partial}{\partial x}((\rho C_P)_m T) + v\frac{\partial}{\partial y}((\rho C_P)_m T)\right) = k_m \left(\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2}\right).$$
(4)

The velocity components in the *x* and *y* directions are given by *u* and *v* in the above equations. In addition, the pressure, temperature, and acceleration due to gravity in the y-direction are given by *P*, *T*, and g, respectively. The last term in Equation (3) is the buoyant Boussinesq source term. The NEPCMs act through the advection term that appears in the energy conservation equation (Equation (4)). The impact of the NEPCM nanoparticles, as well as their effects on the overall performance of the problem, will be discussed later.

Boundary conditions are required for the above coupled equations. As mentioned earlier, the walls of the cavity have a constant temperature of Tc, whereas the temperature of the lower wall is given by a constant temperature of Th. We define the following dimensionless parameters, which leads in return to the dimensionless governing equations given below.

$$(X,Y) = \begin{pmatrix} x \\ L \end{pmatrix}, \begin{pmatrix} y \\ L \end{pmatrix}, (U,V) = \begin{pmatrix} u \\ L \end{pmatrix}, \begin{pmatrix} v \\ L \end{pmatrix}$$
(5)

$$\theta = \frac{T - T_C}{T_H - T_C}, p = \frac{P}{\rho_{bf} r^2 \omega^2}$$
(6)

$$\frac{\partial U}{\partial X} + \frac{\partial V}{\partial Y} = 0 \tag{7}$$

$$\left(\frac{\rho_m}{\rho_{bf}}\right)\left(U\frac{\partial U}{\partial X} + V\frac{\partial U}{\partial Y}\right) = -\frac{\partial p}{\partial X} + \frac{1}{Re_{bf}}\left(\frac{\mu_m}{\mu_{bf}}\right)\left(\frac{\partial^2 U}{\partial X^2} + \frac{\partial^2 U}{\partial Y^2}\right)$$
(8)

$$\left(\frac{\rho_m}{\rho_{bf}}\right)\left(U\frac{\partial V}{\partial X} + V\frac{\partial V}{\partial Y}\right) = -\frac{\partial p}{\partial Y} + \frac{1}{Re_{bf}}\left(\frac{\mu_m}{\mu_{bf}}\right)\left(\frac{\partial^2 V}{\partial X^2} + \frac{\partial^2 V}{\partial Y^2}\right) + \frac{Gr_{bf}}{R_{bf}^2}\left(\frac{(\rho\beta)_m}{(\rho\beta)_{bf}}\right)\theta - \frac{\sigma_m}{\sigma_{bf}}Ha^xV \tag{9}$$

$$\left(U\frac{\partial}{\partial X}(Cr\theta) + V\frac{\partial}{\partial Y}(Cr\theta)\right) = \frac{1}{Re_{bf}Pr_{bf}}\left(\frac{k_m}{k_{bf}}\right)\left(\frac{\partial^2\theta}{\partial X^2} + \frac{\partial^2\theta}{\partial Y^2}\right)$$
(10)

The dimensionless numbers are expressed as follows [36]:

$$v_{bf} = \frac{\mu_{bf}}{\rho_{bf}}, \alpha_{bf} = \frac{k_{bf}}{(\rho c_p)_{bf}}, Pr = \frac{v_{bf}}{\alpha_{bf}}, Re_{bf} = \frac{\omega r^2}{v_{bf}},$$

$$Ha = LB \sqrt{\frac{\sigma_{nf}}{\mu_{nf}}} \text{ and } Gr_{bf} = \frac{g\beta_{bf}\Delta Tr^3}{v_{bf}^2},$$
(11)

where v_{bf} , α_{bf} , Pr, Re_{bf} , and Gr_{bf} represent the kinematic viscosity, thermal diffusion, Prandtl, Reynolds, and Grashof numbers, respectively [34]. Moreover, the dimensionless constant in the energy equation is the ratio between the heat capacity of the mixture and that of the base fluid (*Cr*) [32].

$$Cr = \frac{(\rho C_P)_m}{(\rho C_P)_{bf}} = 1 - \varphi + \lambda \varphi + \frac{\varphi}{\chi} f$$
(12)

The subscripts *m* and *bf* in Equation (12) stand for the nanoliquid mixture and base liquid, respectively. The specific heat capacity and density of liquid $(\rho C_P)_{bf}$ are assumed to be constants. Nevertheless, for certain materials of the NEPCMs and the base liquid, the values of the specific heat capacity of the mixture are a function of the amount of latent heat of the NEPCMs core, the volume fraction of the NEPCMs, and the temperature (*f*). The variability relative to the volume fraction is shown by φ , whereas the variability associated with the latent heat of the NEPCM core is shown by χ . The quantity λ is calculated by [35]:

$$\lambda = \frac{\left(C_{p_{cl}} + lC_{p_s}\right)\rho_c\rho_s}{(\rho_s + l\rho_c)(\rho C_P)_{bf}}$$
(13)

Clearly, λ is the ratio of the heat capacity of NEPCMs in the liquid phase to the base liquid. The subscripts *c*, *l*, and *s* stand for core, liquid, and sell, respectively. The core-shell weight ratio is given by *l* the coefficient and equals 0.447. Moreover, the quantity χ in Equation (12) is calculated using the following [36].

$$\chi = \frac{C_{P,bf}}{h_{sf}/T_{Mr}} \frac{\rho_{bf}(\rho_s + l\rho_c)}{(\rho_s \rho_c)}$$
(14)

$$f = \frac{\pi}{2} \sin\left(\frac{\pi}{\delta}\left(\theta - \theta_f + \frac{\delta}{2}\right)\right) \times \begin{cases} 0 \text{ if } \theta < \theta_f - \frac{\delta}{2} \\ 1 \text{ if } \theta_f - \frac{\delta}{2} < \theta < \theta_f + \frac{\delta}{2} \\ 0 \text{ if } \theta < \theta_f + \frac{\delta}{2} \end{cases}$$
(15)

where $\delta = \frac{T_{Mr}}{\Delta T}$, and $\theta_f = \frac{T_f - T_C}{\Delta T}$. The parameters δ and θ_f stand for the dimensionless internal fusion and fusion temperature, respectively. Note that δ represents the thickness of the melting zone. Going back to Equation (15), one can see that if the nanoliquid temperature is larger than the NEPCM core melting temperature $\left(T > T_f + \frac{T_{Mr}}{2}\right)$ or smaller than the core solidification temperature $\left(T < T_f - \frac{T_{Mr}}{2}\right)$, then the last term of Equation (12) becomes zero, and therefore the values of Cr are reduced. The melting temperature and the melting range also affect the values of *Cr*. Clearly, the heat transfer can be improved by changing the values of *Cr*. The substantial improvement in specific heat capacity when utilizing NEPCMs as opposed to conventional nanoparticles is the main advantage here. The coefficient of f appears in Equation (12) and is the amplitude of the change in specific heat capacity in the two-phase state. In comparison with Equation (15), it can be seen that the latent heat of the NEPCMs core hsf also has effects on the quantity *Cr*.

The rate of heat transfer can be computed using the Nusselt number as [35,36].

$$Nu_{loc} = \frac{k_m}{k_{bf}} \frac{\partial \theta}{\partial n}, Nu_{Ave} = \frac{1}{\pi D} \int Nu_{loc}(n) dn.$$
(16)

where Nu_{loc} and Nu_{Ave} represent the local and the average Nusselt number, respectively, and *D* represents the diameter of the inner cylinder.

We mention here that the studied work was achieved in a permanent simulation (steady state) because the physical phenomena of the fluid in the studied range produces phenomena that are not related to the time Criterion.

2.2. The Thermophysical Properties

The NEPCM consists of an n-nonadecane core and a polyurethane shell. For thermal energy storage, the n-nonadecane is a favorable organic phase change material due to the fact that melting and solidification occur at a temperature of 31 °C. In addition, the latent heat of fusion of 156.07 kJ/kg and the latent heat of solidification of 164.99 kJ/kg. Hence, it can be used in buildings in order to store thermal energy. The operating temperature range and the amount of storable energy make such PCM feasible to be utilized in buildings [37]. On the other hand, the polyurethane shell has high strength, elevated elastic properties, smooth surface [36], and low Crystallinity [35]; therefore, it has been taken into consideration by many researchers. Moreover, such shell has no health or environmental issues related to the release of formaldehyde. Some previous experimental studies helped in extracting the thermophysical properties of water and the aforementioned NEPCM. The preparation of these nanoparticles for different applications was also studied and presented [36]. Considering the thermophysical properties of core and shell, the thermophysical properties of NEPCMs can be described as follows [35]:

$$\rho_n = \frac{(1+l)\rho_c \rho_s}{\rho_s + l\rho_c}.$$
(17)

where n indicates the subscription of nanoparticles. The specific heat capacity of the nucleus can be presented in terms of a sinusoidal function whose argument is a function of temperature; the melting temperature T_f , and the phase change temperature range T_{Mr} , as described earlier. Its amplitude can also be seen as a function of the latent heat of the core h_{sf} , and the specific heat capacity of the core in the liquid state C_p . In addition, the specific heat capacity of the NEPCM core is C_p , in the phase change temperature range [35]:

$$C_{p,c} = C_{pc,l} + \left\{ \frac{\pi}{2} \left(\frac{h_{sf}}{T_{Mr}} - C_{pc,l} \right) \left(\sin \pi \frac{T - (T_f - T_{Mr}2)}{T_{Mr}} \right) \right\} \\ \times \left\{ \begin{array}{c} 0 \text{ if } T < T_f - \frac{T_{Mr}}{2} \\ 1 \text{ if } T_f - \frac{T_{Mr}}{2} < T < T_f + \frac{T_{Mr}}{2} \\ 0 \end{array} \right.$$
(18)
if $T < T_f + \frac{T_{Mr}}{2}$

Below, we introduce the specific heat capacity and thermal expansion coefficient of NEPCM [37]:

$$C_{p,n} = \frac{(C_{p,c} + lC_s)\rho_c\rho_s}{(\rho_s + l\rho_c)\rho_n}$$
(19)

$$\beta_n = \beta_c + \left(\frac{\beta_s - \beta_c}{2}\right) \left(1 - \frac{l\rho_s}{\rho_c}\right) \tag{20}$$

The thermophysical properties of liquid, NEPCM nanoparticles, NEPCM volume fraction, and the water-NEPCM mixture are given in Table 1. The thermophysical equations of the mixture, including density, specific heat capacity, and thermal expansion coefficient, are given by the following equations [35,38].

$$\rho_m = (1 - \varphi)\rho_{bf} + \varphi \tag{21}$$

$$\rho_n \text{ and } C_{P,m} = \frac{(1-\varphi)\rho_f C_{P,bf} + \varphi \rho_n C_{P,n}}{\rho_m}$$
(22)

$$\beta_m = \frac{(1-\varphi)\rho_{bf}\beta_{bf} + \varphi\rho_n\beta_n}{\rho_m} \tag{23}$$

Table 1. The thermophysical properties of the materials that are used at 303 K [32-40].

Material Property	$ ho\left(rac{kg}{m^3} ight)$	$C_p\left(\frac{J}{kgK}\right)$	$m{eta} imes 10^{-5} \left(rac{1}{K} ight)$	$k\left(\frac{W}{mK}\right)$	$\mu imes 10^{-6} \left(rac{kg}{ms} ight)$
Water (Base fluid).	955.6	4180	21	0.615	797
Nonadecane (core)	786	1317	50	0.19	-
Polyurethane (shell)	721	2037	17.3	0.025	-

In all equations, the subscripts m, n, c, and s stand for the mixture, NEPCM, core, and shell, respectively. The performed simulation was conducted with a constant volume fraction of 0.035. Hence, as indicated by [35], the thermal conductivity and dynamic viscosity of the mixture at 303 K are equal to 122×10^{-5} kg/m.s and 0.7 W/m.K, respectively. The thermophysical properties of the materials that are used are given in Table 1.

Validation: The results of the previous study were confirmed using numerical analysis. Despite the fact that there are different numerical techniques that can be employed to solve the differential equations appear in different applications (see, e.g. [41–44]), the Galerkin Finite Element has been applied here to simulate the model's numerical solutions. Figure 2 shows a comparison between the outcomes from the model employed in the current investigation and those shown in Ghalambaz et al. [35].



Figure 2. Comparison of current work with that of Mohammad Ghalambaz et al. reprinted/adapted with permission from Ref. [35]. 2022, Elsevier.

Table 2 shows the average Nusselt number results for various grid sizes for the greatest values of interest parameters and various power-law fluid indices (Ha = 0, Re= 500). In the following computations, several elements of 23,522 were used.

Table 2. Grid sensitivity check (Ha = 0, φ = 0.04, n = 1, Re = 100).

	959	2049	3376	5165	23,522	88,044	
ψ_{\max}	0.0407	0.039	0.036	0.035	0.035	0.035	
Nuavg	14.12	0.039	14.07	14.04	14.02	14.02	

3. Results and Discussion

We present in this section the results that were obtained in order to clarify the thermal and dynamic composition of a homogeneous suspension of NEPCM particles and water. This new type of suspension is studied in a chamber with a lid and trapezoidal geometry with a zigzag bottom. The questions to be studied are:

- The value of Reynolds number (Re) is in the range of 1–500. This number (Re) expresses the velocity of the movement of the top of the container.
- The value of Hartmann's number (Ha) is in the range of 0–100. This element (Ha) determines the magnetic field acting on the space from the outside.
- Finally, the value of the Darcy number (Da) was studied in the range of $10^{-5}-10^{-2}$. The permeability of the medium is determined by this number (Da).
- The shape of the soil changed by changing the number of ripples in the range of 1–4.

We find that the suspension (water + NEPCM particles) moves due to two factors, namely the horizontal movement of the top of the chamber and the force of thermal buoyancy (Gr = 1000). Understanding and deducing the dynamic behavior of the suspension flow and its direct influence on the heat transfer process is primarily done by analyzing the contours of the isotherms, heat capacity and trajectories. The purpose of adding NEPCM particles is to enhance the thermal properties of the water (heat capacity and conductivity).

To understand the effect of the velocity of motion of the wall, Figure 3 shows the incoming changes in the trajectories, dimensionless temperature, and heat capacity as a function of Reynolds number for Ha = 0, Da = 10^{-2} , and N = 2. Figure 3 shows that the motion of the wall Creates a circular flow in space with its center near the moving wall. On the other hand, if we increase the value of Re, the velocity spread of the space increases, which also causes the vortex to move to the right side of the space. It can be concluded that the velocity of the rotating flow in the space increases as Re increases. For the dimensionless temperature, we find that there is an accumulation of isotherms near the bottom as Re increases. This accumulation is an indication of the high value of the temperature gradient. That is, the thermal activity interacts positively with respect to Re. Moreover, it is observed that the concentration of isotherms on the right-hand side of the soil is better than on the left-hand side, which means that the heat transfer is better on the right side of the soil. The heat capacity contours reflect the same behavior derived from the isotherms and streamlines. First, we note that the red line indicates the regions where the change in the physical state of the NEPCM particles occurs. We note that the red line follows the motion of the flow. The high-temperature gradient makes the line narrow and close to the bottom on the right side.

Figure 4 shows the effects of soil shape on suspension stirring and thermal pattern. Therefore, Figure 4 demonstrates the effects of increasing the number of zigzag lines on the streamlines, isotherms, and heat capacity for Ha = 0, Da = 10^{-2} , and Re = 10. When examining the streamlines, it is clear that the greater the number of zigzag lines on the bottom, the slower the velocity of the suspension flow. This is normal because the more zigzag lines the bottom has, the more it hinders the flow motion, which leads to a decrease in its velocity. As the velocity gradually decreases, the dimensionless temperatures also show a decrease in isothermal density near the ground, i.e., a decay in heat transfer in the form of the zigzag number. The same observations can be derived from the heat capacity diagrams. That is, the larger the zigzag number, the greater the distance of the red bar from the ground is.

Figure 5 shows the effect of magnetic field strength (Ha = 0–100) on the motion of the suspension and its thermal pattern for N = 2, Da = 10^{-2} , and Re = 10. As the suspension moves in a magnetic field, a force, the Lorentz force, is generated that impedes the motion of the suspension in the chamber. Therefore, the streamlines show a decrease in the velocity of the suspension flow as the number Ha grows. In addition, the size of the main vortex formed decreases as the value of Ha increases. The isotherms and heat capacity contours also confirm that the velocity decreases because the temperature gradient near the bottom decreases, and the distance between the bottom and the heat capacity line increases with increasing Ha.



Figure 3. Effect of Reynolds number of the isotherm, heat capacity, and streamlines for Ha = 0, $Da = 10^{-2}$, and N = 2.



Figure 4. Effect of the number of zigzags on the isotherm, heat capacity, and streamlines for Ha = 0, $Da = 10^{-2}$, and Re = 10.



Figure 5. Effect of the number of Ha on the isotherm, heat capacity, and streamlines for $Da = 10^{-2}$ and N = 2.

Figure 6 explains the evolution of the individual streamlines, heat capacity, and dimensionless temperature as a function of the Darcy number for N = 2, Re = 10, and Ha = 0. The value of the Darcy number expresses the permeability of the space, i.e., the larger this value, the better the permeability is. The streamlines show that the higher the Da number, the higher the velocity of the suspended particles and the more pronounced the development of the vortex is. The reason, of course, is the permeability of the space. That is, the greater the permeability, the more positive effect on the levitation is. In the contours of heat capacity and isotherms, there is a rise in the temperature gradient near the bottom, and the red line approaches the bottom because of the growth in the flow velocity inside the medium.

Figure 7 shows the evolution of the Nu number of the hot soil surface as a function of Re and Da for N = 2 and Ha = 0. It can be noted that the higher the value of Re or Da, the higher the value of the Nu number is. The reason for this increase is that the higher the value of Da or Re, the better the flow velocity is, which increases the heat transfer, i.e., the value of the Nusselt number increases.

The effects of the Ha and Re numbers on Nu for N = 2 and $Da = 10^{-2}$ are shown in Figure 8. It can be observed that the higher the value of the Ha number, the lower the value of Nu is, which indeed, results from the Lorentz force, that prevents the movement of the suspension and, as a result, causes a reduction in the thermal activity between the fluid and the hot wall.

In Figure 9, we show the bottom wall shape's effect on the suspension flow's thermal performance. As the number N illustrates the effect of the number of zigzag lines of the soil on the number of Nu for $Da = 10^{-2}$ and without the MHD effect (Ha = 0). It can be observed that the higher the number of bottom ripples, the lower the value of the Nusselt number is, i.e., the thermal activity decreases, which is due to the shape of the bottom, i.e., the increase in ripples makes it more difficult for the flow to stick to the bottom, reducing the heat transfer and thus decreasing the value of the Nusselt number.



Figure 6. Effect of the number of Da on the isotherm, heat capacity, and streamlines for Ha = 0 and N = 2.



Figure 7. Variation of Nu versus Re and Da for N = 2 and Ha = 0.



Figure 8. Variation of Nu versus Re and Ha for N = 2 and $Da = 10^{-2}$.



Figure 9. Variation of Nu versus Re and N for Ha = 0 and Da = 10^{-2} .

Through our analysis of all figures, we can adopt a general idea, which is that the better and easier the fluid movement, the greater the thermal transfer.

4. Conclusions

In this work, we simulate numerically the mixed convection of NEPCM confined in a trapezoidal lidded cavity saturated with porous media and subject to a magnetic field. While the upper cavity wall is insulated and moves at a constant speed, we let the lower one be zigzagged and maintained at a cold temperature. In addition, the two sloping side walls are maintained at a hot temperature. The effects of Reynolds number (Re = 1, 10, 100, 500), zigzag number of the lower wall (N = 1, 2, 3, 4), Hartmann number (Ha = 0, 25, 50, 100), and Darcy number (Da = 10^{-5} , 10^{-4} , 10^{-3} , 10^{-2}) are presented and analyzed. From the study results, we demonstrate the following conclusions:

- Higher values of Re increase the velocity of the flow, resulting in better temperature distribution.
- An increase in the zigzag number of the lower wall hinders the flow velocity; this leads to a limitation of the temperature distribution and a decrease in the heat capacity.
- Increasing the value of Ha (increasing the intensity of the magnetic field) had a negative effect on the heat transfer inside the cavity because the magnetic field inhibits the flow velocity inside the cavity.
- Heat transfer inside the cavity lid improved at higher values of Darcy number because the permeability of the cavity increased with increasing Da.
- At the highest Re, an increase in Da resulted in a 100% increase in Nu, while an increase in N and Ha resulted in a 45% and 38% decrease in Nu, respectively.

Author Contributions: Conceptualization, A.A. and H.L.; methodology, H.L.; software, A.A.; validation, M.A.-K. and N.A.; formal analysis, K.G. and O.Y.; investigation, N.A. and R.M.; resources, R.M.; writing—original draft preparation, A.A. and M.A.-K.; writing—review and editing, A.A., K.G. and M.A.-K.; visualization, O.Y. and R.M.; supervision, A.A.; project administration, A.A.; funding acquisition, M.A.-K. All authors have read and agreed to the published version of the manuscript.

Funding: Not applicable.

Data Availability Statement: Not applicable.

Acknowledgments: The authors extend their appreciation to the Deanship of Scientific Research at King Khalid University for funding this work through General Research Project under grant number (RGP.1/322/43). The authors would like to thank the Deanship of Scientific Research at Umm Al-Qura University for supporting this work by Grant Code: (22UQU4331317DSR90). Authors also extend their appreciation to Mathematics Department at Khalifa University for supporting this work.

Conflicts of Interest: The authors declare no conflict of interest.

Nomenclatures

Greek symbols		Greek symbols	
C _r	the ratio of heat capacity of mixture to water	α	thermal diffusivity (m^2/s)
C_p	specific heat capacity (J/kgK)	δ	thickness of melting zone
f	fusion function	θ	non-dimensional temperature
8	gravitational acceleration (m/s^2)	λ	the ratio of heat capacity of NEPCM to water
Gr	Grashof number	Н	length of the cavity
h _{sf}	latent heat of core (J/kg)	μ	dynamics viscosity (kg/ms)
k	thermal conductivity (W/mK)	υ	kinematic viscosity (m^2/s)
L	length of the cavity (m)	ϑ_s	volume ratio of core to shell
1	weight ratio of core to shell	ρ	density (kg/m^3)
Nu _{Avg}	average Nusselt number	ω	angular velocity of inner cylinder (rad/s)
p	pressure (Pa)	χ	the intensity of stored energy in the core
Pr	Prandtl number	Subscript	
Re	Reynolds number	f	fusion
Т	temperature (K)	п	nanoparticle
T _{Mr}	melting temperature range (K)	т	mixture of water-NEPCM
T_f	fusion temperature (K)	1	liquid phase
V	velocity vector (m/s)	S	sell of the NEPCM

References

- Mourad, A.; Abderrahmane, A.; Younis, O.; Marzouki, R.; Alazzam, A. Numerical Simulations of Magnetohydrodynamics Natural Convection and Entropy Production in a Porous Annulus Bounded by Wavy Cylinder and Koch Snowflake Loaded with Cu–Water Nanofluid. *Micromachines* 2022, 13, 182. [CrossRef] [PubMed]
- Salimpour, M.R.; Kalbasi, R.; Lorenzini, G. Constructal multi-scale structure of PCM-based heat sinks. *Contin. Mech. Thermodyn.* 2017, 29, 477–491. [CrossRef]
- 3. Cui, Y.; Xie, J.; Liu, J.; Wang, J.; Chen, S. A review on phase change material application in building. *Adv. Mech. Eng.* 2017, *9*, 1687814017700828. [CrossRef]
- Aftab, W.; Usman, A.; Shi, J.; Yuan, K.; Qin, M.; Zou, R. Phase change material-integrated latent heat storage systems for sustainable energy solutions. *Energy Environ. Sci.* 2021, 14, 4268–4291. [CrossRef]
- Choudhari, V.; Dhoble, A.; Panchal, S. Numerical analysis of different fin structures in phase change material module for battery thermal management system and its optimization. *Int. J. Heat Mass Transf.* 2020, *163*, 120434. [CrossRef]
- Bhatti, M.M.; Bég, O.A.; Abdelsalam, S.I. Computational Framework of Magnetized MgO–Ni/Water-Based Stagnation Nanoflow Past an Elastic Stretching Surface: Application in Solar Energy Coatings. *Nanomaterials* 2022, 12, 1049. [CrossRef]
- Bhatti, M.; Öztop, H.F.; Ellahi, R.; Sarris, I.E.; Doranehgard, M. Insight into the investigation of diamond (C) and Silica (SiO₂) nanoparticles suspended in water-based hybrid nanofluid with application in solar collector. *J. Mol. Liq.* 2022, 357, 119134. [CrossRef]
- Abderrahmane, A.; Qasem, N.A.A.; Younis, O.; Marzouki, R.; Mourad, A.; Shah, N.A.; Chung, J.D. MHD Hybrid Nanofluid Mixed Convection Heat Transfer and Entropy Generation in a 3-D Triangular Porous Cavity with Zigzag Wall and Rotating Cylinder. *Mathematics* 2022, 10, 769. [CrossRef]
- Rasool, G.; Saeed, A.M.; Lare, A.I.; Abderrahmane, A.; Guedri, K.; Vaidya, H.; Marzouki, R. Darcy-Forchheimer Flow of Water Conveying Multi-Walled Carbon Nanoparticles through a Vertical Cleveland Z-Staggered Cavity Subject to Entropy Generation. *Micromachines* 2022, 13, 744. [CrossRef]
- Bhatti, M.M.; Bég, O.A.; Ellahi, R.; Abbas, T. Natural Convection Non-Newtonian EMHD Dissipative Flow Through a Microchannel Containing a Non-Darcy Porous Medium: Homotopy Perturbation Method Study. *Qual. Theory Dyn. Syst.* 2022, 21, 1–27. [CrossRef]
- 11. Miansari, M.; Nazari, M.; Toghraie, D.; Akbari, O.A. Investigating the thermal energy storage inside a double-wall tank utilizing phase-change materials (PCMs). *J. Therm. Anal. Calorim.* **2020**, *139*, 2283–2294. [CrossRef]
- 12. Masoumi, H.; Khoshkhoo, R.H.; Mirfendereski, S. Experimental and numerical investigation of melting/solidification of nanoenhanced phase change materials in shell & tube thermal energy storage systems. *J. Energy Storage* **2022**, 47, 103561. [CrossRef]
- 13. Alehosseini, E.; Jafari, S.M. Micro/nano-encapsulated phase change materials (PCMs) as emerging materials for the food industry. *Trends Food Sci. Technol.* **2019**, *91*, 116–128. [CrossRef]
- 14. Alehosseini, E.; Jafari, S.M. Nanoencapsulation of phase change materials (PCMs) and their applications in various fields for energy storage and management. *Adv. Colloid Interface Sci.* 2020, 283, 102226. [CrossRef] [PubMed]
- Ghalambaz, M.; Mehryan, S.; Mashoofi, N.; Hajjar, A.; Chamkha, A.J.; Sheremet, M.; Younis, O. Free convective meltingsolidification heat transfer of nano-encapsulated phase change particles suspensions inside a coaxial pipe. *Adv. Powder Technol.* 2020, *31*, 4470–4481. [CrossRef]
- 16. Ghalambaz, M.; Groşan, T.; Pop, I. Mixed convection boundary layer flow and heat transfer over a vertical plate embedded in a porous medium filled with a suspension of nano-encapsulated phase change materials. *J. Mol. Liq.* **2019**, 293, 111432. [CrossRef]
- 17. Bondareva, N.S.; Sheremet, M.A.; Oztop, H.F.; Abu-Hamdeh, N. Entropy generation due to natural convection of a nanofluid in a partially open triangular cavity. *Adv. Powder Technol.* **2017**, *28*, 244–255. [CrossRef]
- 18. Reza Seyf, H.; Wilson, M.R.; Zhang, Y.; Ma, H.B. Flow and heat transfer of nanoencapsulated phase change material slurry past an unconfined square cylinder. *J. Heat Transf.* **2014**, *136*, 051902. [CrossRef]
- 19. Varol, Y. Natural convection in divided trapezoidal cavities filled with fluid saturated porous media. *Int. Commun. Heat Mass Transf.* 2010, *37*, 1350–1358. [CrossRef]
- Alshuraiaan, B. Natural convection in a trapezoidal porous filled cavity with a flexible wall. J. Porous Media 2020, 23, 851–864. [CrossRef]
- 21. Hossain, M.S.; Alim, M.; Andallah, L.S. Numerical simulation of MHD Natural convection flow within porous trapezoidal cavity with heated triangular obstacle. *Int. J. Appl. Comput. Math.* **2020**, *6*, 1–27. [CrossRef]
- 22. Khalil, W.H.; Azzawi, I.D.; Al-Damook, A. The optimisation of MHD free convection inside porous trapezoidal cavity with the wavy bottom wall using response surface method. *Int. Commun. Heat Mass Transf.* **2022**, *134*, 106035. [CrossRef]
- 23. Hussain, S.H. Analysis of heatlines and entropy generation during double-diffusive MHD natural convection within a tilted sinusoidal corrugated porous enclosure. *Eng. Sci. Technol. Int. J.* **2016**, *19*, 926–945. [CrossRef]
- 24. Garoosi, F.; Talebi, F. Numerical analysis of conjugate natural and mixed convection heat transfer of nanofluids in a square cavity using the two-phase method. *Adv. Powder Technol.* **2017**, *28*, 1668–1695. [CrossRef]
- 25. Raji, A.; Hasnaoui, M. Mixed convection heat transfer in a rectangular cavity ventilated and heated from the side. *Numer. Heat Transf. Part A Appl.* **1998**, *33*, 533–548. [CrossRef]
- 26. Yang, L.; Du, K. A comprehensive review on the natural, forced, and mixed convection of non-Newtonian fluids (nanofluids) inside different cavities. *J. Therm. Anal. Calorim.* **2020**, 140, 2033–2054. [CrossRef]

- 27. Raizah, Z.; Aly, A.M. Effect of dual-rotation on MHD natural convection of NEPCM in a hexagonal-shaped cavity based on time-fractional ISPH method. *Sci. Rep.* **2021**, *11*, 1–18. [CrossRef]
- Zidan, A.; Nayak, M.; Karimi, N.; Dogonchi, A.S.; Chamkha, A.J.; Ben Hamida, M.B.; Galal, A.M. Thermal management and natural convection flow of nano encapsulated phase change material (NEPCM)-water suspension in a reverse T-shaped porous cavity enshrining two hot corrugated baffles: A boost to renewable energy storage. J. Build. Eng. 2022, 53, 104550. [CrossRef]
- 29. Ebrahimi, D.; Yousefzadeh, S.; Akbari, O.A.; Montazerifar, F.; Rozati, S.A.; Nakhjavani, S.; Safaei, M.R. Mixed convection heat transfer of a nanofluid in a closed elbow-shaped cavity (CESC). *J. Therm. Anal. Calorim.* **2021**, *144*, 2295–2316. [CrossRef]
- Aly, A.M.; Alsedias, N.; Galal, A.M. The conformable fractal systems of natural convection in an annulus suspended by NEPCM. Int. Commun. Heat Mass Transf. 2022, 134, 106023. [CrossRef]
- Chen, J.; Zhang, X.; Zhang, P.; Deng, J. Variational multiscale element free Galerkin method for natural convection with porous medium flow problems. *Int. J. Heat Mass Transf.* 2017, 107, 1014–1027. [CrossRef]
- Abderrahmane, A.; Qasem, N.A.A.; Mourad, A.; Al-Khaleel, M.; Said, Z.; Guedri, K.; Younis, O.; Marzouki, R. Enhancing the Melting Process of Shell-and-Tube PCM Thermal Energy Storage Unit Using Modified Tube Design. *Nanomaterials* 2022, 12, 3078. [CrossRef] [PubMed]
- Qiu, L.; Zhu, N.; Feng, Y.; Michaelides, E.E.; Żyła, G.; Jing, D.; Zhang, X.; Norris, P.M.; Markides, C.N.; Mahian, O. A review of recent advances in thermophysical properties at the nanoscale: From solid state to colloids. *Phys. Rep.* 2020, 843, 1–81. [CrossRef]
- Feng, G.; Feng, Y.; Qiu, L.; Zhang, X. Evaluation of thermal performance for bionic porous ceramic phase change material using micro-computed tomography and lattice Boltzmann method. *Int. J. Therm. Sci.* 2022, 179, 107621. [CrossRef]
- Ghalambaz, M.; Chamkha, A.J.; Wen, D. Natural convective flow and heat transfer of Nano-Encapsulated Phase Change Materials (NEPCMs) in a cavity. Int. J. Heat Mass Transf. 2019, 138, 738–749. [CrossRef]
- 36. Ghalambaz, M.; Mehryan, S.; Zahmatkesh, I.; Chamkha, A. Free convection heat transfer analysis of a suspension of nano– encapsulated phase change materials (NEPCMs) in an inclined porous cavity. *Int. J. Therm. Sci.* **2020**, 157, 106503. [CrossRef]
- Abderrahmane, A.; Al-Khaleel, M.; Mourad, A.; Laidoudi, H.; Driss, Z.; Younis, O.; Guedri, K.; Marzouki, R. Natural Convection within Inversed T-Shaped Enclosure Filled by Nano-Enhanced Phase Change Material: Numerical Investigation. *Nanomaterials* 2022, 12, 2917. [CrossRef]
- Arıcı, M.; Tütüncü, E.; Yıldız, Ç.; Li, D. Enhancement of PCM melting rate via internal fin and nanoparticles. Int. J. Heat Mass Transf. 2020, 156, 119845. [CrossRef]
- Golab, E.; Goudarzi, S.; Kazemi-Varnamkhasti, H.; Amigh, H.; Ghaemi, F.; Baleanu, D.; Karimipour, A. Investigation of the effect of adding nano-encapsulated phase change material to water in natural convection inside a rectangular cavity. *J. Energy Storage* 2021, 40, 102699. [CrossRef]
- Sadr, A.N.; Shekaramiz, M.; Zarinfar, M.; Esmaily, A.; Khoshtarash, H.; Toghraie, D. Simulation of mixed-convection of water and nano-encapsulated phase change material inside a square cavity with a rotating hot cylinder. *J. Energy Storage* 2022, 47, 103606. [CrossRef]
- 41. Wu, S.-L.; Al-Khaleel, M.D. Semi-discrete Schwarz waveform relaxation algorithms for reaction diffusion equations. *BIT Numer. Math.* **2014**, *54*, 831–866. [CrossRef]
- Wu, S.-L.; Al-Khaleel, M. Convergence analysis of the Neumann–Neumann waveform relaxation method for time-fractional RC circuits. Simul. Model. Pract. Theory 2016, 64, 43–56. [CrossRef]
- 43. Wu, S.-L.; Al-Khaleel, M.D. Optimized waveform relaxation methods for RC circuits: Discrete case. *ESAIM: Math. Model. Numer. Anal.* **2016**, *51*, 209–223. [CrossRef]
- Al-Khaleel, M.D.; Gander, M.J.; Ruehli, A.E. Optimized waveform relaxation solution of RLCG transmission line type circuits. In Proceedings of the 2013 9th International Conference on Innovations in Information Technology, Al Ain, United Arab Emirates, 17–19 March 2013; pp. 136–140.