

# Article



# Study on Lightweight and Strengthening Effect of Carbon Nanotube in Highly Ordered Nanoporous Nickel: A Molecular Dynamics Study

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**Abstract:** The mechanical behavior of nanocomposites consisting of highly ordered nanoporous nickel (HONN) and its carbon nanotube (CNT)-reinforced composites (CNHONNs) subjected to a high temperature of 900 K is investigated via molecular dynamics (MD) simulations. The study indicates that, out-of-plane mechanical properties of the HONNs are generally superior to its in-plane mechanical properties. Whereas the CNT shows a significant strengthening effect on the out-of-plane mechanical properties of the CNHONN composites. Compared to pure HONNs, through the addition of CNTs from 1.28 wt‰ to 5.22 wt‰, the weight of the composite can be reduced by 5.83‰ to 2.33% while the tensile modulus, tensile strength, compressive modulus and compressive strength can be increased by 2.2% to 8.8%, 1% to 5.1%, 3.6% to 10.2% and 4.9% to 10.7%, respectively. The energy absorption capacity can also be improved due to the existence of CNTs. Furthermore, the MD simulations provide further insights into the deformation mechanism at the atomic scale, including fracture in tension, pore collapse in compression and local changes in lattice structures due to stacking faults.

**Keywords:** nanoporous nickel; carbon nanotube/nickel composite; molecular dynamics; lightweight; strengthening effect

## 1. Introduction

Being a type of electrochemical devices, solid oxide fuel cells (SOFCs) have been receiving attentions as one of the most promising power generation systems [1]. While Nickel (Ni), as an important material in electrochemical technology, is used widely as anode material in SOFCs [2–5]. Typical SOFC anodes contain 40–60 wt% nickel, which imparts favorable microstructure, electrical conductivity, thermal robustness and catalytic activity to the anode [4]. In contrast to disordered porous structures, materials with high surface area and highly ordered pore network would be expected to exhibit superior performance in electrochemical systems which rely on gas or liquid diffusion within the electrode material [6]. Nelson et al. [6] have successfully produced highly ordered nanoporous nickel (HONN) via the liquid crystal templated electrodeposition route. Moreover, the monolithic HONNs make electron transport within the material not to encounter grain boundary resistances, offering the material superior electronic conductivity and removing power losses [7].

Since SOFCs usually operate at a range of temperature from 600 to 1000  $^{\circ}$ C, thermal stresses in various components of SOFCs strongly impact the cell lifetime by inducing cracks, contact loss and structural instability [8]. It is desired that the HONNs have not only enough chemical and

electrochemical stability but also excellent mechanical properties to withstand stresses that arise during processing and service.

On the other hand, the discovery of carbon nanotubes (CNTs) [9] stimulated enormous studies on carbon nanomechanics because of their extraordinary strength [10], high electrical properties [11] and high thermal conductivity [12]. CNTs have been used in electrode materials of various cells [13–15] and reinforcement fillers of different structural materials [16–18]. Porosity have been assumed to be an important factor that affects the mechanical properties of anode materials of SOFCs [3,19,20] and therefore, mechanical properties of HONNs with different porosity and their CNT reinforced composites (CNHONNs) are of interest.

Most experimental studies on CNT-Ni composites were dedicated to the development and improvement of fabrication process for improving the mechanical properties [21–25] or to use it as functional materials [14,26–29]. Sun et al. [21] fabricated CNT-reinforced nickel nanocomposites by using an innovative electrochemical co-deposition process and found that the tensile strength of CNT/Ni composite is as high as 2 GPa, which is three times that of pure nickel. Hwang et al. [24] processed CNT/Ni composite by a molecular-level mixing (MLM) process followed by spark plasma sintering (SPS) and the result showed that the composite exhibits a yield strength of 710 MPa, about 3.7 times that of monolithic nickel. Borkar et al. [25] constructed two types of CNT/Ni nanocomposites by using simple mechanical dry milling of nickel and CNT powders or MLM process as precursors followed by SPS and found that both types of nanocomposites exhibits a substantial enhancement in microhardness and tensile yield strength as compared to SPS processed pure Ni. Song et al. [26] used CNTs as a functional additive to improve the electrochemical performance of pasted nickel-foam electrodes for rechargeable alkaline nickel batteries and found that as compared with the electrode without CNTs, the electrode with 3 wt% CNT addition exhibits a higher charging efficiency, greater specific discharge capacity, higher discharge voltage, better high-rate capability and superior cycling stability. Jiang et al. [14] synthesized CNT/Ni hybrid nanostructured arrays (NSAs) on a stainless steel substrate through a one-step chemical-vapor-deposition (CVD) method using nullaginite NSAs as starting materials and found that the composite can be used as high-performance electrode materials for electrochemical capacitors (ECs). And the composites displayed well-defined pseudo capacitive capabilities with a high areal capacitance and good rate capability owing to the unique interconnected hybrid structure inherited from the conductive CNT network and porous NSAs. What's more, the addition of CNTs can also improve the properties of ethanol biosensor [28], glucose biosensor [29].

With the advent of supercomputers and high performance numerical tools, molecular dynamics (MD) simulations have become effective means to predict the properties of nanocomposites. Particularly for the study of failure mechanisms, MD simulations have significant advantages due to the capability of capturing dynamic atomic behaviors [30,31]. As for CNT/Ni nanocomposites, a series of MD simulation or ab initio MD simulation works have concentrated on the growth mechanism of CNT on the surface of Ni nanoparticles [32–35] or using Ni as a coating on CNT surface to provide an effective channel for load transfer between the CNT and metals (Al [36], Au [37], Mg [38], Cu [39]). Investigations and analyses, however, of using CNT as a reinforced material in Ni matrix, especially HONNs, are still rare.

In this work, the in-plane and out-of-plane Young's modulus *E* and ultimate strength  $\sigma_u$  of HONNs with different porosity are studied by using MD simulations at a temperature of 900 K. The effect of different mass fractions of CNTs on mechanical properties is also investigated.

#### 2. Methodology

Representative volume elements (RVEs) of HONNs (see Figure 1a,b) and CNHONNs (see Figure 1c–f) are chosen for MD simulations to investigate the mechanical behavior of the RVE in the in-plane (*y* axis) and out-of-plane (*z* axis) direction. The *x*, *y* and *z* axes correspond to [010], [100], [001] directions, respectively. The lattice constant of Ni is 3.52 Å and the size of RVE is 140.8 Å × 140.8 Å (as shown in Figure 1a), with a through cylindrical pore in the center and a quarter

cylindrical pore in each corner which have identical radius with central pore (as shown in Figure 1b). In order to investigate the influence of relative density, 7 classes of samples with a center cylindrical pore of different radii are built and the simulated system contains 112,720~175,440 atoms, as listed in Table 1. The relative density  $\rho_r$  of the samples is calculated by  $\rho_r = M_H/M_B$ , where  $M_B$  is the sum of atomic mass of the bulk Ni and  $M_{\rm H}$  is the sum of atomic mass of HONNs. And its porosity can be calculated simply by  $(1-\rho_r)$ . In order to investigate the influence of CNTs, we take the 4th group of samples as the example by adding different mass fractions of CNTs (as shown in Figure 1c–f). In previous published studies, both zigzag [40] and armchair [36–39,41–43] CNTs were used to strengthen different types of material. And the (4, 4) CNT [36,37,42,43] was often used to simulate the mechanical properties of CNT/metal composites. When investigating the strengthening effect, the present work is limited to a smaller diameter of CNT, that is, (4, 4) CNT. The effect of the diameter and chirality of the CNTs will be the topic of further study. The locations of CNTs are in the center of ligament throughout z axis and distribute symmetrically around the central pore (as shown in Figure 1f). The mass fractions of CNTs  $\omega_{CNT}$  is calculated by  $\omega_{CNT} = M_C / (M_C + M_H)$ , where  $M_C$  is the sum of atomic mass of the CNT. The  $\omega_{\text{CNT}}$  in our simulations are 1.28% (Figure 1c), 2.58% (Figure 1d), 3.89% (Figure 1e) and 5.22% (Figure 1f) respectively. The rate of weight reduction due to the addition of CNT can be calculated by  $(M_{\rm H}-M_{\rm CNH})/M_{\rm H}$  and they are 5.82% (Figure 1c), 1.16% (Figure 1d), 1.75% (Figure 1e) and 2.33% (Figure 1f) respectively in our simulations, where *M*<sub>CNH</sub> is the sum of atomic mass of CNHONNs. Atomic images are visualized via Ovito [44].



**Figure 1.** (a) Perspective view and (b) In-plane (top) view (*x-y* plane) of representative volume elements (RVE) for highly ordered nanoporous nickel (HONN); and in-plane views for CNT-reinoforced HONNs(CNHONNs) with  $\omega_{\text{CNT}}$  of (c) 1.28%, (d) 2.58%, (e) 3.89% and (f) 5.22%.

Group	Radius of the Pore(Å)	<b>Relative Density</b>	Number of Atoms
1	42.24	0.43	112,720
2	40.48	0.48	123,920
3	38.72	0.52	134,800
4	36.96	0.57	146,320
5	35.2	0.61	155,680
6	33.44	0.65	165,840
7	31.68	0.68	175,440

Table 1. Parameters of RVEs of HONNs.

MD simulations are performed by means of the Large-scale Atomic/Molecular Massively Parallel Simulator (LAMMPS) [45]. The Ni-Ni interaction is commonly described by an embedded atom method (EAM) potential [46] and the binding interaction of C-C is described by the widely used adaptive intermolecular reactive empirical bond order (AIREBO) potential [47]. The Lennard-Jones (LJ) 12-6 potential is adopted to simulate interatomic forces between C and Ni atoms and long rang van der Waals (vdW) interaction between C atoms. The LJ parameters  $\varepsilon_{Ni-C}$  and  $\sigma_{Ni-C}$  are determined by Lorentz-Berthelot (LB) mixing rule [48]. The value of  $\varepsilon_{Ni-C}$  is estimated from geometric average ( $\varepsilon_{Ni-C} = \sqrt{\varepsilon_{Ni}\varepsilon_C}$ ) between  $\varepsilon_{Ni}$  and  $\varepsilon_C$  and the value of  $\sigma_{Ni-C}$  is estimated from the arithmetic average ( $(\sigma_{Ni}+\sigma_C)/2$ ) between  $\sigma_{Ni}$  and  $\sigma_C$ . For non-bonded C atoms Kutana and Giapis [49] recommended the LJ parameters  $\varepsilon_C = 0.00296$  ev and  $\sigma_C = 3.407$  Å and Foiles et al. [46] suggested the LJ parameters  $\varepsilon_{Ni} = 0.219877$  ev and  $\sigma_{Ni} = 2.30$  Å. Hence, according to the LB rule, the LJ parameters for the C-Ni interface are  $\varepsilon_{Ni-C} = 0.02551$  ev and  $\sigma_{Ni-C} = 2.8535$  Å. The equilibrium distance is  $0.858\sigma_{Ni-C}$  [50], So the distance between the CNTs and Ni matrix is taken as 2.45 Å (Figure 1c). The cutoff radius is chosen as 10.2 Å which is larger than  $3\sigma_{Ni-C}$  (8.56 Å), in order to avoid truncation errors and fit to the cutoff distance of the AIREBO potential [42].

The periodic boundary conditions are applied in all three directions. Prior to loading, all models are relaxed using NPT ensemble (the total number of particles, the system's volume and the system's pressure are constant) for 100 ps to minimize the total free energy. In previous MD studies [30,42,51], the Young's modulus was obtained from the stress-strain curves with the NVT ensemble (the total number of particles, the system's volume and the absolute temperature are constant). Our simulations find that for nanoporous nickel, deformation under the NVT ensemble rather than NPT ensemble is more reasonable especially in compression because no densification stage will be found under the NPT ensemble. So the NVT ensemble is adopted to simulate the compressive or tensile behavior by applying a constant strain rate of  $0.005 \text{ ps}^{-1}$ . Young's moduli *E* are calculated by fitting the data of stress-strain curves before a strain of 2% because no dislocations are generated before this strain. All simulations are carried out at the temperature of 900 K. The time step is reduced to 0.5 fs to prevent a blowout of the CNT atoms during fracture [42].

### 3. Results and Discussion

## 3.1. Stress-Strain Curves and Lattice Structure Evolution of HONNs

It is noted that the stress-strain curves for the groups 1 to 7 HONNs have the same trend and they just differ in the magnitude of the values. Without loss of generality, here we choose the 4th group of HONNs as examples to investigate the mechanical behavior and lattice structure evolution of the HONNs during the in-plane and out-of-plane tension and compression. Figure 2a,c show the stress-strain curves of the HONNs and Figure 2b,d show the corresponding variation of the fraction  $\eta$  of hexagonal close-packed (HCP) atoms with strain. The fraction  $\eta$  of HCP atoms is calculated by  $\eta = N_{\text{HCP}}/N_{\text{all}}$ , where  $N_{\text{HCP}}$  is the total number of HCP atoms and  $N_{\text{all}}$  is the total number of atoms in the sample. In order to study the relationship between stress-strain curves and the fraction  $\eta$  of HCP atoms, Figure 2a,b are divided by dash lines **A**, **B**, **C** and **D** (with the black dash lines **A** and **C** corresponding to the in-plane direction and the red dash lines **B** and **D** corresponding to the out-of-plane direction), respectively, while Figure 2c,d are divided by dash lines **E** to **J**, noticing that lines **F** and **G** are overlapped and indicated by blue line specially.



Figure 2. Stress-strain curves and the hexagonal close-packed (HCP) atom fraction  $\eta$  of group 4 for HONNs along different deformation direction.

It can be seen from Figure 2a that the in-plane and out-of-plane tensile curves have the same trend. Taking the in-plane tensile curve as example, the stress-strain curve of tension can be divided into two stages (as shown in Figure 2a): (It) initial linear elastic stage and (IIt) the strain-softening and rupture. At the beginning of tension, sample enters into the linear elastic stage and reaches the peak stress (defined as the ultimate tensile strength  $\sigma_{\rm ut}$ ) (Stage I<sub>t</sub>). After that the stress decreases as the strain is further increased and finally begins to rupture (Stage II<sub>t</sub>). Dash lines **A** and **B** correspond to the peak stresses for the in-plane and out-of-plane tension at strain of 0.08 and 0.091, respectively. Figure 2b shows that the fraction  $\eta$  of HCP atoms is at a lower level when the stress reaches ultimate tensile strength  $\sigma_{\rm ut}$ . After which  $\eta$  sharply rises and reaches the peak value quickly, leading to a dramatic strain-softening (see Figure 2a) because of the accumulation of stacking faults. Figure 3 gives the atomic configurations at different stages based on the common neighbor analysis (CNA) [52]. Figure 3a,d show the atomic snapshots with the strain corresponding to the  $\sigma_{\rm ut}$  and it can be found that few regions of the samples have transformed into plastic deformation from elastic deformation because a small amount of dislocations and stacking faults are generated. As the applied strain increases to a critical value of 0.236 (dash line C), the fraction  $\eta$  of HCP atoms of the in-plane tension reaches a peak value and the corresponding stress-strain curve exhibits a faster softening rate when the applied strain exceeds the critical value of 0.236. Figure 3b shows the atomic snapshot with strain of 0.236 for the in-plane tension, which is the result excluding the surface atoms. Cracks can be observed inside the ligament (marked in Figure 3b) and lots of stacking faults are formed around the cracks. As the applied strain further increases, rupture of the sample (marked in Figure 3c) occurs around the strain of 0.32. As for the out-of-plane tension in Figure 3e, a more homogeneous necking is observed compared to the in-plane tension in Figure 3b, so its stress-strain curve shows a more steadily softening rate. When the applied strain increases to the value of 0.423 (dash line **D**), the fraction  $\eta$  of HCP atoms changes into a nearly steady value from a negative slope and then the corresponding stress-strain curve exhibits a slightly faster softening rate due to the rupture of the samples (marked in Figure 3f).

As for compression, it can be seen from Figure 2c that the stress-strain curves of the in-plane and out-of-plane compressions also have the same trend. Taking the in-plane compression as example, the stress-strain curves of compression can be divided into four stages (as shown in Figure 2c): (I<sub>c</sub>) initial linear elastic stage, (II<sub>c</sub>) instability of pore, (III<sub>c</sub>) collapse propagation and (IV<sub>c</sub>) densification stage. At the beginning of compression, sample enters into the linear elastic stage and reaches the peak stress (defined as the ultimate compressive strength  $\sigma_{uc}$ ) (Stage I<sub>c</sub>). After reaching the peak stress, the pores become unstable and begin to collapse, which shows a dramatic strain-softening in stress-strain curve (Stage  $II_c$ ) and then the collapse begins to propagate, which causes a plateau stage in stress-strain curve (Stage III<sub>c</sub>). Finally, the strain-hardening effect becomes noticeable in the stress-strain curves which means the sample enters into densification stage (Stage IV<sub>c</sub>). The densification strain  $\varepsilon_d$  can be determined roughly by the intersection of tangents to the stress-strain curve for the cell wall collapse regime and densification regime [53]. Dash lines E and F correspond to the peak stresses for the out-of-plane and in-plane compressions at a strain of 0.049 and 0.068, respectively. Similar to the tension  $\eta$ -strain curves, the fraction  $\eta$  of HCP atoms has a very small value when the stress reaches the ultimate compressive strength  $\sigma_{uc}$ . From the atomic snapshots corresponding to  $\sigma_{uc}$  as shown in Figure 3g,j, it can be seen that, a few regions of the sample have entered into plastic deformation from elastic deformation. And then the sample enters into Stage II<sub>c</sub> (see Figure 2c). The end of Stage II<sub>c</sub> for the out-of-plane and in-plane compressions is divided by dash lines G and H at strain of 0.068 and 0.102, respectively. Figure 3d shows that the fractions  $\eta$  of HCP atoms rise sharply during Stage II<sub>c</sub>. The atomic snapshots at the strain of 0.068 and 0.102 shown in Figure 3h,k hints that more and more dislocations and stacking faults are generated during Stage II<sub>c</sub>. After that the fractions  $\eta$  of HCP atoms slowly increase through the Stage III<sub>c</sub> and enter into Stage IV<sub>c</sub> eventually. The atomic snapshots at the densification strain determined by tangent corresponding to dash lines I and J are shown in Figure 3i,l, in which the void is completely filled with the atoms and the RVE is nearly densified accompanied with lots of dislocations and stacking faults.



**Figure 3.** Atomic snapshots at specified macro loading strain: (**a**, **d**, **g** and **j**) peak stress, (**b**, **e**, **h** and **k**) stage of collapse or strain-softening, (**c**, **f**, **i** and **l**) densification or rupture of samples. Identification of atomic configurations is based on common neighbor analysis (CNA) and four types of atoms are painted in different colors: green for face-centered-cubic (FCC) atoms, blue for body-centered-cubic (BCC) atoms, red for hexagonal close-packed (HCP) atoms and gray for non-12-coordinated atoms.

## 3.2. Mechanical Properties of HONNs

In this section, we focus on Young's modulus *E* and ultimate strength  $\sigma_u$  of the HONNs. As shown in Figure 4a,b, they strongly depend on the relative density and loading direction and show approximately linear relationship in both Young's modulus *E* and ultimate strength  $\sigma_u$  with respect to the relative density. The Young's moduli *E* (as shown in Figure 4a) increase with the increase of the relative density and the out-of-plane Young's moduli *E* are obviously higher than the in-plane Young's moduli, which is similar to traditional honeycomb materials [54]. The compressive moduli are slightly bigger than the tensile moduli in the in-plane direction, while they are almost equal in the out-of-plane direction. Similar to the *E*, the ultimate strengths  $\sigma_u$  (as shown in Figure 4b) increase with the increase of the relative density. Figure 4b also shows that the ultimate strengths of out-of-plane tension are significantly higher than those in other three cases. It is noted that, although the out-of-plane compressive Young's moduli *E* are higher than both the in-plane tensile and compressive Young's moduli to the smaller strain (around 0.05) corresponding to the ultimate strengths.



**Figure 4.** (a) Young's modulus *E* and (b) ultimate strength  $\sigma_u$  of HONNs with different relative density.

As the relative density changes from 0.43 to 0.68, in the in-plane direction, the tensile modulus increases by 144.1% from 25.1 GPa to 61.28 GPa; the compressive modulus increases by 141% from 28.7 GPa to 61.17 GPa;  $\sigma_{ut}$  increases by 104.1% from 1.95 GPa to 3.69 GPa; and  $\sigma_{uc}$  increases by 155.3% from 1.79 GPa to 4.57 GPa. For the out-of-plane direction, the tensile modulus increases from 49.66 GPa to 90.06 GPa by 81.3%; the compressive modulus increases from 48.63 GPa to 91.06 GPa by 87.2%;  $\sigma_{ut}$  increases from 3.54 GPa to 6.24 GPa by 76.2%; and  $\sigma_{uc}$  increases from 1.96 GPa to 4.55 GPa by 132.1%. From Figure 4 it can be seen that the out-of-plane mechanical properties are generally superior to those of the in-plane direction. The main reason for this phenomenon is the in-plane direction has a weakness in ligament (marked in Figure 3b,c) which may induce stress concentration and lead to successive deformation and rupture of ligament. Compared with the in-plane direction, the samples have a more homogeneous deformation in the out-of-plane direction.

#### 3.3. Effect of CNTs on Mechanical Properties

Figure 5 plot the mechanical properties of Group 4 (as listed in Table 1) of the CNHONNs with different additions of CNTs  $\omega_{CNT}$ , suggesting that the effect of CNTs strongly depends on deforming direction. In the out-of-plane direction, both the tensile and compressive mechanical properties increase

with the increase of  $\omega_{\text{CNT}}$ . As compared to the HONNs with the same relative density, the increase of mass fraction  $\omega_{\text{CNT}}$  of CNTs ranging from 1.28% to 5.22% can make the increase of the tensile modulus ranging from 2.2% to 8.8%, the tensile strength ranging from 1% to 5.1%, the compressive modulus ranging from 3.6% to 10.2%, the compressive strength ranging from 4.9% to 10.7%. But in the in-plane direction, the mechanical properties of CNHONNs cannot be strengthened or even may be weakened by the addition of CNTs.



**Figure 5.** (a) Young's modulus *E* and (b) ultimate strength  $\sigma_u$  of CNHONNs with different  $\omega_{CNT}$  of CNTs.

Figure 6 plots the average stress of the CNTs embedded in the CNHONNs with a  $\omega_{CNT}$  of 5.22‰ with respect to strain. The stresses in the CNTs considerably depend on the loading directions. Compared to the in-plane tension and compression, Figure 6 shows that the CNTs can bear more loads during the out-of-plane loadings, resulting in a better reinforcement of CNHONNs in the out-of-plane direction due to the addition of CNTs. The interface between CNT and Ni is so weak that the load cannot be effectively transferred in the in-plane direction, resulting in defects at the interface, which is similar to the CNT-reinforced metallic glass nanocomposites [40]. This is why the in-plane mechanical properties of the CNHONNs are similar to those of the HONNs and sometimes they are lower than

those of HONNs, as evidenced by the result in Figure 5. The main reason for the difference in the in-pane and out-of-plane reinforcing effect of CNTs is the anisotropy of CNTs and the weak interfacial bonding between C and Ni.



**Figure 6.** Stress contribution of CNTs in the CNHONN in the case of  $\omega_{CNT}$  of 5.22% along different deformation direction.

Moreover, the stress-strain curves of CNHONN are shown in Figure 7. As for the in-plane tension (see Figure 7a), it can be seen that as the strain increases, the stresses increase linearly for all samples at the initial stage up to peak stresses (i.e., the ultimate strengths) around a strain of 0.08 and then drop sharply owing to the fracture of Ni matrix. The toughness of the CNHONNs, except the  $\omega_{CNT}$  of 3.89%, is slightly higher than that of the HONNs. As for the in-plane compression (see Figure 7b), the stress-strain curves of all samples almost overlapped and reach a peak stress around a strain of 0.068. For out-of-plane tension (see Figure 7c), similar to the in-plane tension, the stresses reach peak values around a strain of 0.088 and then decrease dramatically. Particularly, the area under the curve between the 0.12 and 0.31 of strain values largely increases with the increase of  $\omega_{CNT}$  indicating that the endurance capability of the CNHONNs can be enhanced by the addition of the CNTs, even after the nickel matrix fractures. At around a strain of 0.31, the second sharp decline occurs in the stress-strain curves of the CNHONNs due to the failure of the CNTs. As for the out-of-plane compression (see Figure 7d), the stresses reach peak values around the strain of 0.05 and the plateau stresses generally increase with the increase of  $\omega_{CNT}$ , which demonstrates that the energy absorption capacity of the CNHONNs can be improved by the addition of the CNTs.





Figure 7. Cont.



**Figure 7.** Stress-strain curves of CNHONNs along different deformation direction: (**a**) In-plane tension, (**b**) In-plane compression, (**c**) Out-of-plane tension and (**d**) Out-of-plane compression.

## 4. Conclusions

In the present study, MD simulations were conducted to investigate the mechanical behavior of HONNs with different relative densities and CNHONNs with different  $\omega_{CNT}$ . The stress-strain curves of HONN were obtained under the in-plane tension and compression, out-of-plane tension and compression. The variations of the curves were analyzed through MD snapshots which enable detailed analyses of atomic behaviors. The fraction of HCP atoms was also given to explain the deformation mechanism in the HONNs. The overall mechanical properties of the HONNs, such as Young's modulus *E* and ultimate strength  $\sigma_u$ , linearly increase with the increase of the relative density and in the out-of-plane direction they are generally superior to those in the in-plane direction. As for the effects of CNTs to the mechanical properties of CNHONNs, it strongly depends on loading directions. Our analysis revealed that, compared to the in-plane loadings, the CNTs can bear more loads during the out-of-plane loadings (i.e., the axial direction of the CNTs), resulting in a better reinforcement of CNHONNs in the out-of-plane direction due to the addition of CNTs. Furthermore, CNTs show a nice reinforced effect on energy absorption capacity in out-of-plane direction.

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