

Article



On the Asymptotic Behavior of Advanced Differential Equations with a Non-Canonical Operator

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Abstract: In this paper, we aim to study the oscillatory behavior of a class of even-order advanced differential equations with a non-canonical operator. In addition, we present results on the asymptotic behavior of this type of equations and provide an example that illustrates our main results.

Keywords: oscillation; even-order; advanced differential equations; asymptotic behavior

1. Introduction

In recent decades, many authors have studied problems of a number of different classes of advanced differential equations including the asymptotic and oscillatory behavior of their solutions, see [1–8] and the references cited therein. For some more recent oscillation results, see [9–20]. The interest in studying advanced differential equations is also caused by the fact that they appear in models of several areas in science. In [21–23], singular systems of differential equations are used to study the dynamics and stability properties of electrical power systems. Some additional mathematical background on this can be found in [24]. Systems of differential equations with delays are used to study additional properties of electrical power systems in [25,26]. Non-linear advanced differential equations can be used to describe complex dynamical networks, see [27–29], and bring new insight to their stability. Furthermore, this type of equations can be also used in the modeling of dynamical networks of interacting free-bodies, see [30]. Finally, properties of advanced differential equations are used in the study of singular differential equations of fractional order, see [31,32]. Several other examples in Physics can be found in [33]. In this paper, we consider an even-order non-linear advanced differential equation with a non-canonical operator of the following type:

$$L_{y} + q(v) g(y(\eta(v))) = 0, \quad L_{y} := \left(a(v) \left(y^{(\kappa-1)}(v)\right)^{\beta}\right)',$$
(1)

where $v \ge v_0$, κ is even and β is a quotient of odd positive integers. The operator L_y is said to be in canonical form if $\int_{v_0}^{\infty} a^{-1/\beta}(s) ds = \infty$; otherwise, it is called noncanonical. Throughout this work, we suppose that:

C1: $a \in C^1([v_0,\infty),\mathbb{R})$, a(v) > 0, $a'(v) \ge 0$,

C2: $q, \eta \in C\left([v_0, \infty), \mathbb{R}\right), q\left(v\right) \ge 0, \eta\left(v\right) \ge v, \lim_{v \to \infty} \eta\left(v\right) = \infty,$

C3: $g \in C(\mathbb{R}, \mathbb{R})$ such that $g(x) / x^{\beta} \ge k > 0$, for $x \ne 0$ and under the condition

$$\zeta(v) = \int_{v_0}^{\infty} \frac{1}{a^{1/\beta}(s)} \mathrm{d}s < \infty.$$
⁽²⁾

Definition 1. The function $y \in C^{\kappa-1}[v_y, \infty)$, $v_y \ge v_0$, is called a solution of (1), if $(y^{(\kappa-1)}(v))^{\beta} \in C^1[v_y, \infty)$, for $a \in C^1([v_0, \infty), \mathbb{R})$, a(v) > 0 and y(v) satisfies (1) on $[v_y, \infty)$.

Definition 2. Let

$$D = \{(v,s) \in \mathbb{R}^2 : v \ge s \ge v_0\}$$
 and $D_0 = \{(v,s) \in \mathbb{R}^2 : v > s \ge v_0\}$

A kernel function $H_i \in C(D, \mathbb{R})$ is said to belong to the function class \Im , written by $H \in \Im$, if, for i = 1, 2, J

- (*i*) $H_i(v,s) > 0$, on D_0 and $H_i(v,s) = 0$ for $v \ge v_0$ with $(v,s) \notin D_0$;
- (ii) $H_i(v,s)$ has a continuous and nonpositive partial derivative $\partial H_i/\partial s$ on D_0 and there exist functions $\tau, \vartheta \in C^1([v_0, \infty), (0, \infty))$ and $h_i \in C(D_0, \mathbb{R})$ such that

$$\frac{\partial}{\partial s}H_{1}(v,s) + \frac{\tau'(s)}{\tau(s)}H_{1}(v,s) = h_{1}(v,s)H_{1}^{\beta/(\beta+1)}(v,s)$$
(3)

and

$$\frac{\partial}{\partial s}H_{2}(v,s) + \frac{\vartheta'(s)}{\vartheta(s)}H_{2}(v,s) = h_{2}(v,s)\sqrt{H_{2}(v,s)}.$$
(4)

Next we will discuss the results in [34–36]. Actually, our purpose in this article is to complement and improve these results. Agarwal et al. in [34,35] studied the even-order nonlinear advanced differential equations

$$\left(\left(y^{(\kappa-1)}\left(v\right)\right)^{\beta}\right)' + q\left(v\right)y^{\beta}\left(\eta\left(v\right)\right) = 0.$$
(5)

By means of the Riccati transformation technique, the authors established some oscillation criteria of (5). Grace and Lalli [36] investigated the second-order neutral Emden–Fowler delay dynamic equations

$$y^{(\kappa)}(v) + q(v)y(\eta(v)) = 0,$$
(6)

and established some new oscillation for (5) under the condition

$$\int_{v_0}^{\infty} \frac{1}{a^{1/\beta}(s)} \mathrm{d}s = \infty.$$
(7)

To prove this, we apply the previous results to the equation

$$y^{(\kappa)}(v) + \frac{q_0}{v^{\kappa}} y(\lambda v) = 0, \ v \ge 1.$$
(8)

if we set $\kappa = 4$ and $\lambda = 2$, then by applying conditions in [34–36] on Equation (8), we find the results in [35] improves those in [36]. Moreover, the those in [34] improves results in [35,36]. Thus, the motivation in our paper is to complement and improve results in [34–36]. We will use the following methods:

- Integral averaging technique.
- Riccati transformations technique.
- Method of comparison with second-order differential equations.

We will also use the following lemmas from (1):

Lemma 1 ([3]). If $y^{(i)}(v) > 0$, $i = 0, 1, ..., \kappa$, and $y^{(\kappa+1)}(v) < 0$, then

$$\frac{y\left(v\right)}{v^{\kappa}/\kappa!} \geq \frac{y'\left(v\right)}{v^{\kappa-1}/\left(\kappa-1\right)!}.$$

Lemma 2 ([19]). Suppose that $y \in C^{\kappa}([v_0, \infty), (0, \infty))$, $y^{(\kappa)}$ is of a fixed sign on $[v_0, \infty)$, $y^{(\kappa)}$ not identically zero and there exists a $v_1 \ge v_0$ such that

$$y^{(\kappa-1)}(v) y^{(\kappa)}(v) \le 0,$$

for all $v \ge v_1$. If we have $\lim_{v\to\infty} y(v) \ne 0$, then there exists $v_{\theta} \ge v_1$ such that

$$y(v) \geq \frac{\theta}{(\kappa-1)!} v^{\kappa-1} \left| y^{(\kappa-1)}(v) \right|,$$

for every $\theta \in (0, 1)$ and $v \geq v_{\theta}$.

Lemma 3 ([2]). Let β be a ratio of two odd numbers, V > 0 and U are constants. Then

$$Ux - Vx^{(eta+1)/eta} \leq rac{eta^{eta}}{(eta+1)^{eta+1}} rac{U^{eta+1}}{V^{eta}}, \ V > 0.$$

Lemma 4. Suppose that y is an eventually positive solution of (1). Then, there exist three possible cases:

$$\begin{array}{ll} (\mathbf{S}_1) & y\left(v\right) > 0, \ y'\left(v\right) > 0, \ y''\left(v\right) > 0, \ y^{(\kappa-1)}\left(v\right) > 0, \ y^{(\kappa)}\left(v\right) < 0, \\ (\mathbf{S}_2) & y\left(v\right) > 0, \ y^{(r)}(v) > 0, \ y^{(r+1)}(v) < 0 \ for \ all \ odd \ integer \\ & r \in \{1, 3, ..., \kappa - 3\}, \ y^{(\kappa-1)}(v) > 0, \ y^{(\kappa)}(v) < 0, \\ (\mathbf{S}_3) & y\left(v\right) > 0, \ y^{(\kappa-2)}\left(v\right) > 0, \ y^{(\kappa-1)}\left(v\right) < 0, \ L_y \le 0, \end{array}$$

for $v \ge v_1$, where $v_1 \ge v_0$ is sufficiently large.

2. Oscillation Criteria

Theorem 1. Assume that (2) holds. If the differential equations

$$\left(\frac{(\kappa-2)!a^{\frac{1}{\beta}}(v)}{(\theta v^{\kappa-2})^{\beta}}\left(y'(v)\right)^{\beta}\right)' + kq(v)y^{\beta}(v) = 0, \quad \forall \theta \in (0,1),$$
(9)

$$y''(v) + y(v) \frac{1}{(\kappa - 4)!} \int_{v}^{\infty} (\varsigma - v)^{\kappa - 4} \left(\frac{1}{a(\varsigma)} \int_{\varsigma}^{\infty} q(s) \, \mathrm{d}s\right)^{1/\beta} \mathrm{d}\varsigma = 0, \tag{10}$$

and

$$\left(a\left(v\right)\left(y'\left(v\right)\right)^{\beta}\right)' + y^{\beta}\left(v\right)kq\left(v\right)\left(\frac{\zeta\left(\eta\left(v\right)\right)}{\zeta\left(v\right)}\right)^{\beta}\left(\frac{\theta_{1}}{(\kappa-2)!}\eta^{\kappa-2}\left(v\right)\right)^{\beta} = 0, \quad \theta_{1} \in (0,1)$$
(11)

are oscillatory for every constant θ , $\theta_1 \in (0, 1)$, then every solution of (1) is either oscillatory or satisfies $\lim_{v \to \infty} y(v) = 0$.

Proof. Assume to the contrary that *y* is a positive solution of (1). Then, we can suppose that y(v) and $y(\eta(v))$ are positive for all $v \ge v_1$ sufficiently large. From Lemma 4, we have three possible cases (**S**₁), (**S**₂) and (**S**₃). Let case (**S**₁) hold. Using Lemma 2, we find

$$y'(v) \ge \frac{\theta}{(\kappa-2)!} v^{\kappa-2} y^{(\kappa-1)}(v), \qquad (12)$$

for every $\theta \in (0, 1)$ and for all large *v*. We set

$$\varphi(v) := \tau(v) \left(\frac{a(v) \left(y^{(\kappa-1)}(v) \right)^{\beta}}{y^{\beta}(v)} \right),$$
(13)

and observe that $\varphi\left(v\right)>0$ for $v\geq v_{1}$, where $\tau\in C^{1}\left(\left[v_{0},\infty\right),\left(0,\infty\right)\right)$ and

$$\begin{split} \varphi'(v) &= \tau'(v) \frac{a(v) \left(y^{(\kappa-1)}(v)\right)^{\beta}}{y^{\beta}(v)} + \tau(v) \frac{\left(a \left(y^{(\kappa-1)}\right)^{\beta}\right)'(v)}{y^{\beta}(v)} \\ &-\beta \tau(v) \frac{y^{\beta-1}(v) y'(v) a(v) \left(y^{(\kappa-1)}(v)\right)^{\beta}}{y^{2\beta}(v)}. \end{split}$$

Using (12) and (13), we obtain

$$\varphi'(v) \leq \frac{\tau'_{+}(v)}{\tau(v)}\varphi(v) + \tau(v)\frac{\left(a\left(v\right)\left(y^{(\kappa-1)}\left(v\right)\right)^{\beta}\right)'}{y^{\beta}(v)} \\ -\beta\tau(v)\frac{\theta}{(\kappa-2)!}v^{\kappa-2}\frac{a\left(v\right)\left(y^{(\kappa-1)}\left(v\right)\right)^{\beta+1}}{y^{\beta+1}(v)} \\ \leq \frac{\tau'(v)}{\tau(v)}\varphi(v) + \tau(v)\frac{\left(a\left(v\right)\left(y^{(\kappa-1)}\left(v\right)\right)^{\beta}\right)'}{y^{\beta}(v)} \\ -\frac{\beta\theta v^{\kappa-2}}{(\kappa-2)!\left(\tau(v)a\left(v\right)\right)^{\frac{1}{\beta}}}\varphi(v)^{\frac{\beta+1}{\beta}}.$$

$$(14)$$

From (1) and (14), we obtain

$$\varphi'\left(v\right) \leq \frac{\tau'\left(v\right)}{\tau\left(v\right)}\varphi\left(v\right) - k\tau\left(v\right)\frac{q\left(v\right)y^{\beta}\left(\eta\left(v\right)\right)}{y^{\beta}\left(v\right)} - \frac{\beta\theta v^{\kappa-2}}{\left(\kappa-2\right)!\left(\tau\left(v\right)a\left(v\right)\right)^{\frac{1}{\beta}}}\varphi\left(v\right)^{\frac{\beta+1}{\beta}}.$$

Note that y'(v) > 0 and $\eta(v) \ge v$, thus, we find

$$\varphi'(v) \leq \frac{\tau'(v)}{\tau(v)}\varphi(v) - k\tau(v)q(v) - \frac{\beta\theta v^{\kappa-2}}{(\kappa-2)!(\tau(v)a(v))^{\frac{1}{\beta}}}\varphi(v)^{\frac{\beta+1}{\beta}}.$$
(15)

If we set $\tau(v) = k = 1$ in (15), then we find

$$\varphi'(v) + \frac{\beta \theta v^{\kappa-2}}{(\kappa-2)! a^{\frac{1}{\beta}}(v)} \varphi(v)^{\frac{\beta+1}{\beta}} + q(v) \le 0.$$

From [37], we can see that Equation (9) is non-oscillatory, which is a contradiction. Let case (S_2) hold. If we set

$$\psi(v) := \vartheta(v) \frac{y'(v)}{y(v)},$$

we see that $\psi(v) > 0$ for $v \ge v_1$, where $\vartheta \in C^1([v_0, \infty), (0, \infty))$. By differentiating $\psi(v)$, we find

$$\psi'(v) = \frac{\vartheta'(v)}{\vartheta(v)}\psi(v) + \vartheta(v)\frac{y''(v)}{y(v)} - \frac{1}{\vartheta(v)}\psi(v)^2.$$
(16)

Now, by integrating (1) from *v* to *m* and using y'(v) > 0, we get

$$a(m)\left(y^{(\kappa-1)}(m)\right)^{\beta} - a(v)\left(y^{(\kappa-1)}(v)\right)^{\beta} = -\int_{v}^{m} q(s) g(y(\eta(s))) ds.$$

By virtue of y'(v) > 0 and $\eta(v) \ge v$, we get

$$a(m)\left(y^{(\kappa-1)}(m)\right)^{\beta} - a(v)\left(y^{(\kappa-1)}(v)\right)^{\beta} \leq -ky^{\beta}(v)\int_{v}^{u}q(s)\,ds.$$

Letting $m \to \infty$, we see that

$$a(v)\left(y^{(\kappa-1)}(v)\right)^{\beta} \ge ky^{\beta}(v)\int_{v}^{\infty}q(s)\,\mathrm{d}s$$

and so

$$y^{(\kappa-1)}(v) \ge y(v) \left(\frac{k}{a(v)} \int_{v}^{\infty} q(s) \,\mathrm{d}s\right)^{1/\beta}$$

Integrating again from *v* to ∞ , κ – 4 times, we get

$$y''(v) + \frac{y(v)}{(\kappa-4)!} \int_{v}^{\infty} (\varsigma - v)^{\kappa-4} \left(\frac{k}{a(\varsigma)} \int_{\varsigma}^{\infty} q(s) \,\mathrm{d}s\right)^{1/\beta} \mathrm{d}\varsigma \le 0.$$
(17)

From (16) and (17), we obtain

$$\psi'(v) \le \frac{\vartheta'(v)}{\vartheta(v)}\psi(v) - \frac{\vartheta(v)}{(\kappa-4)!}\omega(s) - \frac{1}{\vartheta(v)}\psi(v)^2,$$
(18)

where

$$\omega(s) = \int_{v}^{\infty} (\zeta - v)^{\kappa - 4} \left(\frac{k}{a(\zeta)} \int_{\zeta}^{\infty} q(s) \, \mathrm{d}s \right)^{1/\beta} \mathrm{d}\zeta.$$

If we now set $\vartheta(v) = k = 1$ in (18), then we obtain

$$\psi'(v) + \psi^2(v) + rac{1}{(\kappa - 4)!} \omega(s) \zeta \le 0.$$

From [37], we see Equation (10) is non-oscillatory, which is a contradiction.

Let case (**S**₃) hold. By recalling that $a(v)(y^{(\kappa-1)}(v))^{\beta}$ is non-increasing, we obtain

$$a^{1/\beta}(s) y^{(\kappa-1)}(s) \le a^{1/\beta}(v) y^{(\kappa-1)}(v), s \ge v \ge v_1$$

Dividing the latter inequality by $a^{1/\beta}(s)$ and integrating the resulting inequality from v to u, we get

$$y^{(\kappa-2)}(u) \le y^{(\kappa-2)}(v) + a^{1/\beta}(v) y^{(\kappa-1)}(v) \int_{v}^{u} a^{-1/\beta}(s) \,\mathrm{ds}.$$

Letting $u \to \infty$, we obtain

$$0 \le y^{(\kappa-2)}(v) + a^{1/\beta}(v) y^{(\kappa-1)}(v) \zeta(v).$$

Thus,

$$\frac{-a^{1/\beta}(v)y^{(\kappa-1)}(v)\zeta(v)}{y^{(\kappa-2)}(v)} \le 1.$$
(19)

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Furthermore, we get

 $\left(\frac{y^{(\kappa-2)}(v)}{\zeta(v)}\right)' \ge 0,\tag{20}$

due to (19). Now define

$$\phi(v) = \frac{a(v) \left(y^{(\kappa-1)}(v)\right)^{\beta}}{\left(y^{(\kappa-2)}(v)\right)^{\beta}},$$
(21)

we see that $\phi(v) < 0$ for $v \ge v_1$, and

$$\phi'(v) = \frac{\left(a(v)\left(y^{(\kappa-1)}(v)\right)^{\beta}\right)'}{\left(y^{(\kappa-2)}(v)\right)^{\beta}} - \frac{\beta a(v)\left(y^{(\kappa-1)}(v)\right)^{\beta+1}}{\left(y^{(\kappa-2)}(v)\right)^{\beta+1}}.$$

It follows from (1) and (19) that

$$\phi'(v) = \frac{-kq(v)y^{\beta}(\eta(v))}{(y^{(\kappa-2)}(v))^{\beta}} - \frac{\beta\phi^{\beta/\beta+1}(v)}{a^{1/\beta}(v)}$$

From Lemma 2, we find

$$y(v) \ge \frac{\theta_1}{(\kappa-2)!} v^{\kappa-2} y^{(\kappa-2)}(v).$$

$$(22)$$

Thus, we have

$$\phi'\left(\upsilon\right) = \frac{-kq\left(\upsilon\right)y^{\beta}\left(\eta\left(\upsilon\right)\right)}{\left(y^{\left(\kappa-2\right)}\left(\eta\left(\upsilon\right)\right)\right)^{\beta}} \frac{\left(y^{\left(\kappa-2\right)}\left(\eta\left(\upsilon\right)\right)\right)^{\beta}}{\left(y^{\left(\kappa-2\right)}\left(\upsilon\right)\right)^{\beta}} - \frac{\beta\phi^{\beta/\beta+1}\left(\upsilon\right)}{a^{1/\beta}\left(\upsilon\right)}$$

From (22), we obtain

$$\phi'(v) \leq -kq(v) \left(\frac{\theta_1 \eta^{\kappa-2}(v)}{(\kappa-2)!}\right)^{\beta} \left(\frac{\zeta(\eta(v))}{\zeta(v)}\right)^{\beta} - \frac{\beta \phi^{\beta/\beta+1}(v)}{a^{1/\beta}(v)}.$$
(23)

From [37], we can see that Equation (11) is non-oscillatory, which is a contradiction. Theorem 1 is proved. \Box

Remark 1. It is well known (see [15]) that if

$$\int_{v_0}^{\infty} \frac{1}{a(v)} \mathrm{d}v < \infty, \text{ and } \liminf_{v \to \infty} \left(\int_{v_0}^{v} \frac{1}{a(s)} \mathrm{d}s \right)^{-1} \int_{v}^{\infty} \left(\int_{v_0}^{v} \frac{1}{a(s)} \mathrm{d}s \right)^2 q(s) \, \mathrm{d}s > \frac{1}{4},$$

then Equations (9)–(11) with $\beta = 1$ are oscillatory.

Based on the above results and Theorem 1, we can easily obtain the following Hille and Nehari type oscillation criteria for (1) with $\beta = 1$.

Theorem 2. Let $\beta = k = 1$ and assume that (2) holds. If for $\theta, \theta_1 \in (0, 1)$

$$\liminf_{v \to \infty} \left(\int_{v_0}^{v} \frac{\theta s^{\kappa-2}}{(\kappa-2)! a\left(s\right)} \mathrm{d}s \right)^{-1} \int_{v}^{\infty} \left(\int_{v_0}^{v} \frac{\theta s^{\kappa-2}}{(\kappa-2)! a\left(s\right)} \mathrm{d}s \right)^2 q\left(s\right) \mathrm{d}s > \frac{1}{4}, \tag{24}$$

with

$$\int_{v_{0}}^{\infty}\frac{\theta v^{\kappa-2}}{\left(\kappa-2\right)!a\left(v\right)}\mathrm{d}v<\infty,$$

and if

$$\liminf_{v \to \infty} v \int_{v_0}^{v} \frac{1}{(\kappa - 4)!} \int_{v}^{v} (\varsigma - v)^{\kappa - 4} \left(\frac{1}{a(\varsigma)} \int_{\varsigma}^{v} q(s) \, \mathrm{d}s \right)^{1/\beta} \mathrm{d}\varsigma \mathrm{d}v > \frac{1}{4},\tag{25}$$

$$\liminf_{v \to \infty} \left(\int_{v_0}^{v} \frac{1}{a(s)} \mathrm{d}s \right)^{-1} \int_{v}^{\infty} \left(\int_{v_0}^{v} \frac{1}{a(s)} \mathrm{d}s \right)^{2} \frac{\theta_1 \zeta\left(\eta\left(s\right)\right) \eta^{\kappa-2}\left(s\right) q\left(s\right)}{\zeta\left(s\right)\left(\kappa-2\right)!} \mathrm{d}s > \frac{1}{4}, \tag{26}$$

then every solution of (1) is either oscillatory or satisfies $\lim_{v \to \infty} y(v) = 0$.

In the next theorem, we employ the integral averaging technique to establish a Philos-type oscillation criteria for (1):

Theorem 3. Let (2) holds. If there exist positive functions $\tau, \vartheta \in C^1([v_0, \infty), \mathbb{R})$ such that

$$\limsup_{v \to \infty} \frac{1}{H_1(v, v_1)} \int_{v_1}^{v} (H_1(v, s) k\tau(s) q(s) - \pi(s)) ds = \infty,$$
(27)

$$\limsup_{v \to \infty} \frac{1}{H_2(v, v_1)} \int_{v_1}^{v} \left(H_2(v, s) \frac{\vartheta(s)}{(\kappa - 4)!} \varpi(s) - \frac{\vartheta(s) h_2^2(v, s)}{4} \right) \mathrm{d}s = \infty,$$
(28)

and,

$$\limsup_{v \to \infty} \frac{1}{H_3(v, v_1)} \int_{v_1}^{v} \left(H_3(v, s) \, kq\left(s\right) \left(\frac{\theta_1 \eta^{\kappa-2}\left(s\right)}{(\kappa-2)!}\right)^{\beta} \zeta^{\beta}\left(\eta\left(s\right)\right) - \tilde{\pi}\left(s\right) \right) \mathrm{d}s = \infty,$$

where

$$\pi\left(s\right) = \frac{h_{1}^{\beta+1}\left(v,s\right)H_{1}^{\beta}\left(v,s\right)}{\left(\beta+1\right)^{\beta+1}}\frac{\left(\left(\kappa-2\right)!\right)^{\beta}\tau\left(s\right)a\left(s\right)}{\left(\theta s^{\kappa-2}\right)^{\beta}}$$

and

$$\tilde{\pi}\left(s\right) = \frac{\beta^{\beta+1}H_{3}\left(v,s\right)}{\left(\beta+1\right)^{\beta+1}} \frac{1}{a^{1/\beta}\left(s\right)\zeta\left(s\right)}.$$

Then every solution of (1) *is either oscillatory or satisfies* $\lim_{v \to \infty} y(v) = 0$.

Proof. Assume to the contrary that *y* is a positive solution of (1). Then, we can suppose that y(v) and $y(\eta(v))$ are positive for all $v \ge v_1$ sufficiently large. From Lemma 4, we have three possible cases (**S**₁), (**S**₂) and (**S**₃). Assume that (**S**₁) holds. From Theorem 1, we get that (15) holds. Multiplying (15) by $H_1(v, s)$ and integrating the resulting inequality from v_1 to v we find that

$$\begin{split} \int_{v_1}^{v} H_1\left(v,s\right) k\tau\left(s\right) q\left(s\right) \mathrm{d}s &\leq \quad \varphi\left(v_1\right) H_1\left(v,v_1\right) + \int_{v_1}^{v} \left(\frac{\partial}{\partial s} H_1\left(v,s\right) + \frac{\tau'\left(s\right)}{\tau\left(s\right)} H_1\left(v,s\right)\right) \varphi\left(s\right) \mathrm{d}s \\ &- \int_{v_1}^{v} \frac{\beta \theta s^{\kappa-2}}{\left(\kappa-2\right)! \left(\tau\left(s\right) a\left(s\right)\right)^{\frac{1}{\beta}}} H_1\left(v,s\right) \varphi^{\frac{\beta+1}{\beta}}\left(s\right) \mathrm{d}s. \end{split}$$

From (3), we get

$$\int_{v_{1}}^{v} H_{1}(v,s) k\tau(s) q(s) ds \leq \varphi(v_{1}) H_{1}(v,v_{1}) + \int_{v_{1}}^{v} h_{1}(v,s) H_{1}^{\beta/(\beta+1)}(v,s) \varphi(s) ds - \int_{v_{1}}^{v} \frac{\beta \theta s^{\kappa-2}}{(\kappa-2)! (\tau(s) a(s))^{\frac{1}{\beta}}} H_{1}(v,s) \varphi^{\frac{\beta+1}{\beta}}(s) ds.$$
(29)

Using Lemma 3 with $V = \beta \theta s^{\kappa-2} / \left((\kappa - 2)! (\tau(s) a(s))^{\frac{1}{\beta}} \right) H_1(v,s), U = h_1(v,s) H_1^{\beta/(\beta+1)}(v,s)$

And $y = \varphi(s)$, we get

$$\begin{split} & h_{1}\left(v,s\right)H_{1}^{\beta/(\beta+1)}\left(v,s\right)\varphi\left(s\right) - \frac{\beta\theta s^{\kappa-2}}{\left(\kappa-2\right)!\left(\tau\left(s\right)a\left(s\right)\right)^{\frac{1}{\beta}}}H_{1}\left(v,s\right)\varphi^{\frac{\beta+1}{\beta}}\left(s\right) \\ & \leq \quad \frac{h_{1}^{\beta+1}\left(v,s\right)H_{1}^{\beta}\left(v,s\right)}{\left(\beta+1\right)^{\beta+1}}\frac{\left((\kappa-2)!\right)^{\beta}\tau\left(s\right)a\left(s\right)}{\left(\theta s^{\kappa-2}\right)^{\beta}}, \end{split}$$

which, with (29) gives

$$\frac{1}{H_{1}\left(v,v_{1}\right)}\int_{v_{1}}^{v}\left(H_{1}\left(v,s\right)k\tau\left(s\right)q\left(s\right)-\pi\left(s\right)\right)\mathrm{d}s\leq\varphi\left(v_{1}\right),$$

which contradicts (27). Assume that (S_2) holds. From Theorem 1, we get that (18) holds. Multiplying (18) by $H_2(v,s)$ and integrating the resulting inequality from v_1 to v, we obtain

$$\begin{split} \int_{v_1}^{v} H_2\left(v,s\right) \frac{\vartheta\left(s\right)}{(\kappa-4)!} \varpi\left(s\right) \mathrm{d}s &\leq \psi\left(v_1\right) H_2\left(v,v_1\right) \\ &+ \int_{v_1}^{v} \left(\frac{\partial}{\partial s} H_2\left(v,s\right) + \frac{\vartheta'\left(s\right)}{\vartheta\left(s\right)} H_2\left(v,s\right)\right) \psi\left(s\right) \mathrm{d}s \\ &- \int_{v_1}^{v} \frac{1}{\vartheta\left(s\right)} H_2\left(v,s\right) \psi^2\left(s\right) \mathrm{d}s. \end{split}$$

Thus, from (4), we obtain

$$\int_{v_1}^{v} H_2(v,s) \frac{\vartheta(s)}{(\kappa-4)!} \omega(s) \, ds \le \psi(v_1) H_2(v,v_1) + \int_{v_1}^{v} h_2(v,s) \sqrt{H_2(v,s)} \psi(s) \, ds - \int_{v_1}^{v} \frac{1}{\vartheta(s)} H_2(v,s) \, \psi^2(s) \, ds \le \psi(v_1) H_2(v,v_1) + \int_{v_1}^{v} \frac{\vartheta(s) h_2^2(v,s)}{4} ds$$

and so

$$\frac{1}{H_{2}(v,v_{1})}\int_{v_{1}}^{v}\left(H_{2}(v,s)\frac{\vartheta(s)}{(\kappa-4)!}\varpi(s)-\frac{\vartheta(s)h_{2}^{2}(v,s)}{4}\right)ds \leq \psi(v_{1}),$$

which contradicts (28). Assume that (S_3) holds. Using (19) and (21), we see that

$$-\phi(v)\zeta^{\beta}(v) \le 1 \tag{30}$$

due to (30). Multiplying this inequality by $\zeta^{\beta}(v)$ and integrating the resulting inequality from v_1 to v, we get

$$\zeta^{\beta}(v) \phi(v) - \zeta^{\beta}(v_{1}) \phi(v_{1}) + \beta \int_{v_{1}}^{v} a^{-1/\beta}(s) \zeta^{\beta-1}(s) \phi(s) ds$$

$$\leq -\int_{v_{1}}^{v} kq(s) \left(\frac{\theta_{1}\eta^{\kappa-2}(s)}{(\kappa-2)!}\right)^{\beta} \zeta^{\beta}(\eta(s)) ds - \beta \int_{v_{1}}^{v} \frac{\phi^{\beta/\beta+1}(s)}{a^{1/\beta}(s)} \zeta^{\beta}(s) ds.$$
(31)

Multiplying (31) by $H_3(v,s)$, we find that

$$\begin{split} \int_{v_1}^{v} H_3(v,s) \, kq\,(s) \left(\frac{\theta_1 \eta^{\kappa-2}\,(s)}{(\kappa-2)!}\right)^{\beta} \zeta^{\beta}\,(\eta\,(s)) \, \mathrm{d}s &\leq \zeta^{\beta}\,(v_1)\,\phi\,(v_1)\,H_3\,(v,v_1) - \zeta^{\beta}\,(v)\,\phi\,(v)\,H_3\,(v,v_1) \\ &+ \int_{v_1}^{v} \beta a^{-1/\beta}\,(s)\,\zeta^{\beta-1}\,(s)\,\phi\,(s)\,H_3\,(v,s)\,\mathrm{d}s \\ &- \int_{v_1}^{v} \frac{\beta \phi^{\beta/\beta+1}\,(s)}{a^{1/\beta}\,(s)} \zeta^{\beta}\,(s)\,H_3\,(v,s)\,\mathrm{d}s. \end{split}$$

Using Lemma 3 with $V = \zeta^{\beta}(s) H_3(v,s) / a^{1/\beta}(s)$, $U = a^{-1/\beta}(s) \zeta^{\beta-1}(s) H_3(v,s)$ and $y = \phi(s)$, we get

$$\beta a^{-1/\beta}(s) \zeta^{\beta-1}(s) \phi(s) H_{3}(v,s) - \frac{\beta \phi^{\beta/\beta+1}(s)}{a^{1/\beta}(s)} \zeta^{\beta}(s) H_{3}(v,s)$$

$$\leq \frac{\beta^{\beta+1} H_{3}(v,s)}{(\beta+1)^{\beta+1}} \frac{1}{a^{1/\beta}(s) \zeta(s)}$$

and easily, we find that

$$\frac{1}{H_{3}\left(v,v_{1}\right)}\int_{v_{1}}^{v}\left(H_{3}\left(v,s\right)kq\left(s\right)\left(\frac{\theta_{1}\eta^{\kappa-2}\left(s\right)}{\left(\kappa-2\right)!}\right)^{\beta}\zeta^{\beta}\left(\eta\left(s\right)\right)-\tilde{\pi}\left(s\right)\right)\mathrm{d}s\leq\zeta^{\beta}\left(v_{1}\right)\phi\left(v_{1}\right)+1,$$

which contradicts (27). This completes the proof. \Box

Example 1. We consider the equation

$$\left(v^{5}y^{\prime\prime\prime}(v)\right)' + vq_{0}y(3v) = 0, \ v \ge 1,$$
(32)

where $q_0 > 0$ is a constant. Note that $\beta = 1$, $\kappa = 4$, $a(v) = v^5$, $q(v) = vq_0$ and $\eta(v) = 3v$. If we set k = 1, then condition (24) becomes

$$\begin{aligned} \liminf_{v \to \infty} \left(\int_{v_0}^v \frac{\theta s^{\kappa-2}}{(\kappa-2)! a\left(s\right)} \mathrm{d}s \right)^{-1} \int_v^\infty \left(\int_{v_0}^v \frac{\theta s^{\kappa-2}}{(\kappa-2)! a\left(s\right)} \mathrm{d}s \right)^2 q\left(s\right) \mathrm{d}s \\ = \liminf_{v \to \infty} \left(4v^2 \right) \int_v^\infty \frac{q_0}{16s^3} \mathrm{d}s = \liminf_{v \to \infty} \left(4v^2 \right) \left(\frac{q_0}{32v^2} \right) \\ = \frac{q_0}{8} > \frac{1}{4}, \end{aligned}$$

while condition (25) becomes

$$\liminf_{v \to \infty} v \int_{v_0}^{v} \frac{1}{(\kappa - 4)!} \int_{v}^{v} (\varsigma - v)^{\kappa - 4} \left(\frac{1}{a(\varsigma)} \int_{\varsigma}^{v} q(s) \, \mathrm{d}s \right)^{1/\beta} \mathrm{d}\varsigma \mathrm{d}v = \liminf_{v \to \infty} v \left(\frac{q_0}{4v} \right)$$
$$= \frac{q_0}{4} > \frac{1}{4},$$

and hence condition (26) is satisfied. Therefore, from Theorem 2, all solutions of Equation (32) are oscillatory if $q_0 > 2$.

Remark 2. One can easily see that the results obtained in [18,19] cannot be applied to conditions in Theorem 2, so our results are new.

Remark 3. We can generalize our results by studying the equation in the form

$$\left(a(v)\left(y^{(\kappa-1)}(v)\right)^{\beta}\right)' + \sum_{i=1}^{j} q_i(v) y^{\beta}(\eta_i(v)) = 0, \text{ where } v \ge v_0, j \ge 1.$$

For this we leave the results to researchers interested .

3. Conclusions

In this article we studied we provided three new Theorems on the oscillatory and asymptotic behavior of a class of even-order advanced differential equations with a non-canonical operator in the form of (1).

For researchers interested in this field, and as part of our future research, there is a nice open problem which is finding new results in the following cases:

$$\begin{aligned} & (\mathbf{S}_1) \quad y(v) > 0, \ y'(v) > 0, \ y^{(\kappa-2)}(v) > 0, \ y^{(\kappa-1)}(v) \le 0, \ \left(a(v)\left(y^{(\kappa-1)}(v)\right)^{\beta}\right)' \le 0, \\ & (\mathbf{S}_2) \quad y(v) > 0, \ y^{(r)}(v) < 0, \ y^{(r+1)}(v) > 0, \ \forall r \in \{1, 3, ..., \kappa - 3\}, \\ & \text{ and } y^{(\kappa-1)}(v) < 0, \ \left(a(v)\left(y^{(\kappa-1)}(v)\right)^{\beta}\right)' \le 0. \end{aligned}$$

For all this there is some research in progress.

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