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A Dumbbell Shaped Piezoelectric Motor Driven by the First-Order Torsional and the First-Order Flexural Vibrations

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Abstract: A piezoelectric motor driven by the first-order torsional and first-order flexural (T/F) vibrations is designed, fabricated, and tested in this study. The actuating force is generated by the torsional vibration of the dumbbell-shaped vibrator, while the elliptical motion shape is adjusted with the flexural vibration. The rotor, pressed onto the vibrator's lateral surface, is frictionally driven with the vibrator. Here, the torsional vibration, the shear modes of piezoelectric ceramics, and the driving method may contribute to high torque and high output power. To test the feasibility of our proposal, first, a prototype of the T/F vibrator is built and its vibration properties are explored. As predicted, the torsional and flexural vibrations are excited on the vibrator. Then, the load characteristics of the piezoelectric motor are investigated. The maximal torque, the no-load rotation speed, and maximal output power are 4.3 Nm, 125 r/min, and 16.9 W, respectively. The results imply that using the first-order torsional and the first-order flexural vibrations is a feasible method to achieve high torque and high output power of piezoelectric motors.

Keywords: piezoelectric motor; torsional vibration; flexural vibration; vibrator

1. Introduction

Based on the inverse piezoelectric effect [1-4], piezoelectric motors convert electrical energy into mechanical energy and utilize frictional force to achieve actuation [5-10]. Compared to electromagnetic motors, they exhibit rapid response, self-locking, and absence of electromagnetic interfere [7,10-14]. Recently, piezoelectric motors have become increasingly in demand in optical instruments, automatic stages, and robotics' arm joints [1,7,12,15]. Vibrators are core components for piezoelectric motors, and they are generally composed by lead-zirconate-titanate (PZT) ceramics and metal vibrating bodies [1,16-19]. PZT ceramics basically work in three modes—i.e., thickness (d₃₃), extensional (d₃₁), and shear (d₁₅) modes [1,20], where the d₃₃ and d₃₁ modes have been commonly adopted in conventional piezoelectric motors. For example, Zhang et al. [11] clamped the PZT disks with the d33 modes with a frog-shaped vibrating body to form a biconically inspired piezoelectric motor. Cao et al. [10] bonded several PZT disks on a polymeric vibrating body and utilized the d₃₁ mode to drive the slider. Though satisfactory performances—e.g., compact structure and good



controllability—have been achieved with the aforementioned motors [10,11], their mechanical outputs still need improvement. In most cases, the electromechanical coupling factors are higher corresponding to d_{15} modes than to d_{33} and d_{31} modes [21]; this infers the probability to obtain large output of piezoelectric motors based on d₁₅ modes. Despite the probability, there are fewer reports about piezoelectric motors based on the d_{15} modes than those based on the d_{31} and d_{33} ones except the following cases: Yan et al. [2] and Dong et al. [3,4] bonded thin PZT disks, polarized along the length directions instead of the thickness directions, onto vibrating bodies and excited two torsional (T/T) vibrations to drive the micromotors. Though the motors had larger power densities than the conventional ones that have identical structures but work in the d_{31} and/or d_{33} modes, it was difficult to markedly improve the performance because of the large mechanical loss on the boundaries between the PZT disks and the vibrating bodies. Satonobu et al. [5] clamped several PZT disks between cylindrical vibrating bodies and excited torsional and longitudinal (T/L) modes to form a piezoelectric motor, where good performance was achieved because not only frictional loss was suppressed but also elliptical motion shapes could be easily adjusted. Here, the T/L motors should be sufficiently long to obtain optimal elliptical-motion shapes, but in most cases, the application fields—e.g., automatic stages and optical instruments [9,13]—cannot provide enough depth for the T/L motors. Considering these problems, it would be necessary to devise other types of piezoelectric motors based on d₁₅ modes.

In this study, we develop a vibrator working in the first-order torsional and the first-order flexural (T/F) vibrations to form a new piezoelectric motor. Here, the torsional vibration generates the driving force, while the flexural vibration adjusts the elliptical motion shapes. In addition, the cylindrical vibrator drives the rotor on its circumferential outer surface instead of the end surface. Thus, it would be interesting to conduct a trial on the T/F motors. First, we build a vibrator prototype and explore its vibration properties. Then, we measure the load characteristics of the T/F motor to verify the feasibility of our proposal.

2. Configuration and Working Principle

2.1. Design Procedure

Figure 1a depicts the configuration of the dumbbell-shaped T/F vibrator, which consists of two vibrating bodies, torsional PZT disks, and flexural PZT disks. The motor is designed with the following procedures:

(1) Determine the material of each part. Here, stainless steel (elastic modulus: 197 GPa, density: 7800 kg/m³, and Poisson's ratio: 0.29) was chosen for the vibrating bodies. The density of the PZT ceramics, provided with our agent, is 7700 kg/m³ and the elastic, piezoelectric, and dielectric matric are as follows:

$$\begin{bmatrix} c^{E} \end{bmatrix} = \begin{bmatrix} 78 & 42 & 20 & & & \\ 42 & 78 & 20 & & & \\ 20 & 20 & 72 & & & \\ & & & 26 & & \\ & & & & & 23 \\ & & & & & & 23 \end{bmatrix} (GPa)[d] = \begin{bmatrix} 0 & 0 & 0 & 0 & 0.70 & 0 \\ 0 & 0 & 0 & 0.70 & 0 & 0 \\ -0.19 & -0.19 & 0.45 & 0 & 0 & 0 \end{bmatrix} \times 10^{-9} (C/N)$$

and

$$\begin{bmatrix} \varepsilon^S \end{bmatrix} = \begin{bmatrix} 22.13 & 0 & 0 \\ 0 & 22.13 & 0 \\ 0 & 0 & 15.93 \end{bmatrix} \times 10^{-9} (F/m)$$

Figure 1b shows the dimensions and the poling directions of PZT disks. The flexural PZT disks (30 mm in outer diameter, 15 mm in inner diameter, and 3 mm in thickness) were divided into two parts with opposite poling directions, while the torsional ones (30 mm in outer diameter, 15 mm in inner diameter, and 4 mm in thickness) were poled in the circumferential direction. The PZT disks were clamped between the vibrating bodies with a stainless-steel screw.

(2) Arrange the PZT disks. Since the motor works in the 1st-order torsional and the 1st-order flexural modes, the torsional PZT disks were set in the middle part of the vibrator. Additionally, the central position of the flexural PZT disks located at the node of the flexural mode. By solving the Equation [21]:

$$\cos\left[\frac{4.73}{l}\left(\Delta l+\frac{l}{2}\right)\right] + \cosh\left[\frac{4.73}{l}\left(\Delta l+\frac{l}{2}\right)\right] - 1.02 \cdot \left\{\sin\left[\frac{4.73}{l}\left(\Delta l+\frac{l}{2}\right)\right] + \sinh\left[\frac{4.73}{l}\left(\Delta l+\frac{l}{2}\right)\right]\right\} = 0, \tag{1}$$

it can be found that the distance between the flexural mode's node and the vibrator's middle surface (Δl) should be equal to 0.275 times of the vibrator's entire length (l). Then, the number of PZT disks were determined according to the distance between each node and the PZT disks' thicknesses.

- (3) Decide the dominant diameters. Here, the diameters of the vibrating bodies adjacent to the PZT disks were set to 30 mm, as torsional PZT in this diameter is easy to produce.
- (4) Conduct modal degeneration. Through the finite element method (FEM; software ANSYS 14.0, ANSYS Incorporated, Canonsburg, PA, USA.), the lengths of the vibrating bodies' bilateral outer parts were adjusted to gather the resonance frequencies of the flexural and torsional vibrations. Here, the FEM model of the vibrator was meshed with hexahedron elements, whose types were SOLID227 and SOLID5 for, respectively, the vibrating bodies and the PZT disks. The electrodes were neglected as they are too thin for meshing. Figure 1a also gives the final values of each dimension. The vibrating bodies include two steps: the outer step with a circumferential surface are 45 mm in diameter and 18.4 mm in length, while the inner step 30 mm in diameter and 8.6 mm in length contacts the PZT disks.
- (5) Design the accessory components. First, two polymetric hemispheres were fixed near the vibrator's edges as the driving feet. Second, Figure 1c illustrates the configuration of the rotor, whose contact rim had the width and height of, respectively, 12 and 10 mm. Third, Figure 1d shows the assembly of the T/F motor. Observably, the rotor was pressed orthogonally to the vibrator by tightening two coil springs. Since the driving-feet distance was approximately equal to the rotor's diameter, the elliptical motions existing on two driving feet frictionally drove the rotor.

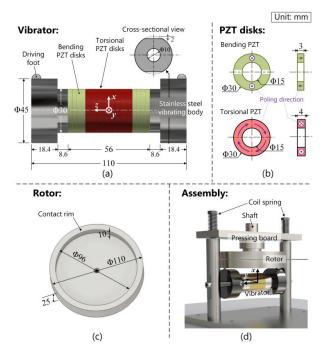


Figure 1. Schematic of the T/F motor. (a) The transducer's configuration, (b) poling direction of flexural and torsional PZT disks, (c) the rotor's structure, and (d) the assembly. A Cartesian coordinate (x, y, and z axes) is built on the middle surface of the transducer.

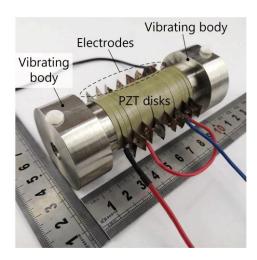


Figure 2 shows the photo of the vibrator prototype, which weights 0.89 kg.

Figure 2. Photograph of the transducer prototype.

2.2. Working Principle

Figure 3 illustrates the working principle of the T/F motor. The voltages applied to the flexural and torsional PZT disks were, respectively, $U_0 \sin(\omega t)$ and $U_0 \cos(\omega t)$, where U_0 , ω , and t represent the voltage amplitude, angular frequency, and time, respectively. The vibration velocities and vibration displacements are shown in Table 1, where A_f and A_t denote the flexural and torsional vibration amplitudes, respectively, and n represents an integer.

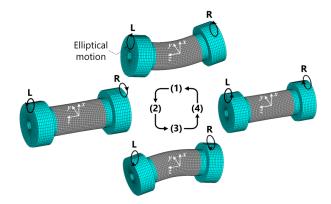


Figure 3. Working principle of the T/F motor. Deformations for the vibration modes are given and the arrows in each circle indicates the direction of each elliptical motion.

Table 1. Vibration velocities and vibration displacements at different time.

Time	Point	Vibration Velocity		Vibration Displacement	
		Flexural	Torsional	Flexural	Torsional
1 0 × 2(0 ⁰ ×	L	0	ωA_t	A_f	0
$\omega t = 0 + 360^{\circ} n$	R	0	$-\omega A_t$	A_f	0
$\omega t = 90^\circ + 360^\circ n$	L	ωA_f	0	Ő	A_t
$\omega t = 90^{\circ} + 360^{\circ} n$	R	ωA_f	0	0	$-A_t$
$\omega t = 180^\circ + 360^\circ$	L	0	$-\omega A_t$	$-A_f$	0
п	R	0	ωA_t	$-\dot{A_f}$	0
$\omega t = 270^\circ + 360^\circ$	L	$-\omega A_f$	0	0	$-A_t$
n	R	$-\omega A_f$	0	0	A_t

Accordingly, the sequence of the motion is given as:

- (1) At $\omega t = 0$, the torsional vibration velocity reaches the peak value. Meanwhile, the flexural vibration displacement is maximal in the +*x* axis. Since the vibrator operates in the 1st torsional mode, the torsional vibration velocities at left (L) and right (R) points are in inverse directions.
- (2) At $\omega t = 90^{\circ}$, the torsional vibration velocity and the flexural vibration displacement are zero.
- (3) At $\omega t = 180^\circ$, the torsional vibration velocity is maximal but the direction is inverse to those in step (1). On the other hand, the flexural vibration displacement is maximal in the-*x* axis.
- (4) At $\omega t = 270^{\circ}$, the torsional vibration velocity and the flexural vibration displacement are zero.

When the voltages applied to the flexural and torsional PZT disks are $-U_0 \sin(\omega t)$ and $U_0 \cos(\omega t)$, the sequence becomes (1)–(4)–(3)–(2); in this case, the rotation direction was inverse to that corresponding to the voltages of $U_0 \sin(\omega t)$ and $U_0 \cos(\omega t)$.

3. Vibration Properties

3.1. Impedance Characteristics

First, the impedance characteristics of the vibrator were explored with the testbed shown in Figure 4a. The electrical ports corresponding to the flexural (or torsional) PZT disks and the ground were connected to the impedance analyzer (4294 A, Agilent, Santa Clara, CA, USA.). Figure 4b,c illustrate the impedance curves of the flexural and torsional vibrations, respectively. To further evaluate the impedance characteristics, some equivalent circuit parameters were estimated as follows [1,2,14]: initially, the resonance frequency (f_r), the anti-resonance frequency (f_a), mechanical quality factor (Q_m), and the motional resistance (R_m) could be directly obtained from the impedance curve. Subsequently, the electromechanical coupling factor (k), motional inductance (L_m), and motional capacitance (C_m) were calculated with the following equations [22–25]:

$$k = \sqrt{1 - \left(\frac{f_r}{f_a}\right)^2},\tag{2}$$

$$L_m = \frac{Q_m R_m}{2\pi f_r},\tag{3}$$

$$C_m = \frac{1}{2\pi f_r Q_m R_m}.$$
(4)

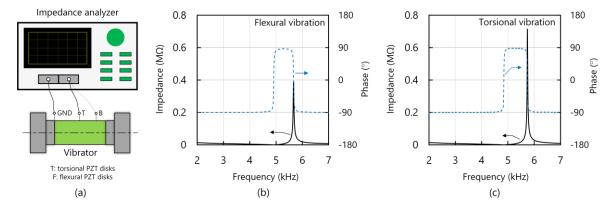


Figure 4. (a) Schematic of the testbed for measuring the impedance characteristics. The impedance characteristics corresponding to (b) the flexural vibration and (c) the torsional vibration.

Table 2 lists the results. The f_r discrepancy was <0.1 kHz between the torsional and flexural vibrations, satisfactorily small for driving the piezoelectric motor. The *k* was 1.2 times higher for the torsional than for the flexural vibration. Because of the high *Q* factor, R_m corresponding to the torsional

vibration was relatively high compared to that of the flexural vibration; this feature implies that the mechanical loss induced by torsional vibration was relatively low [1,18]. Here, since we aimed to prove the potential of using torsional and flexural vibrations to achieve high performance, we adopted a simple structure to achieve modal degeneration; this caused f_r to be in the audible range. In the future, we will shorten the vibrator to change f_r into the inaudible range [18,26–30].

	Flexural Vibration	Torsional Vibration
Resonance frequency, f_r (kHz)	4.911	4.838
Anti-resonance frequency, f_a (kHz)	5.070	5.058
Electromechanical coupling factor, k	24.8%	29.2%
Mechanical quality factor, Q_m	215	302
Motional admittance, Y_{m0} (mS)	9.3	14.1
Motional resistance, R_m (Ω)	107.7	70.9
Motional inductance, L_m (H)	0.750	0.705
Motional capacitance, C_m (nF)	1.40	1.53

Table 2. Equivalent circuit parameters of the T/F vibrator.

3.2. Vibration Velocity Distribution

Then, the vibration velocity distributions were measured to confirm whether the designed vibrations were excited. The vibrator was fixed to a frame with a screw. As shown in Figure 5a, the torsional and flexural vibration velocities were measured with an out-of-plane laser Doppler vibrometer (LDV; type LV-S01, Sunny Optical Corporation, Yuyao, China) [31]. In the measurements, the LDV was arranged horizontally and the light beam's direction was changed with a mirror. Different from the flexural vibration velocity, the torsional vibration velocity was in plane, which is theoretically difficult to directly measure with an out-of-plane LDV. As shown in Figure 5b,c, we set the illumination point on a plastic plate bonded to the vibrator. Since the plastic material has far lower density than the stainless-steel vibrating body, the plate should show negligibly small effect on the vibration properties. It is also worth mentioning that, as shown in Figure 6, the in-plane vibration velocity measured by the aforementioned method showed a linear velocity (in the *y* axis), equal to the product of the angular torsional vibration velocity (in the *z* axis) and the radii. Here, the torsional vibration velocities shown in Figure 6b were in the angular form.

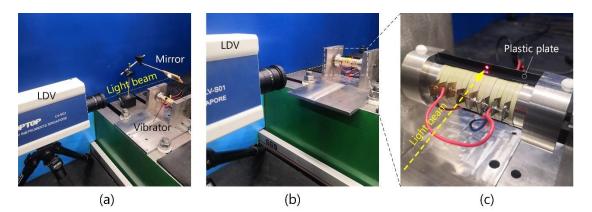


Figure 5. Experimental setup for measuring (**a**) flexural and (**b**) torsional vibration velocity distributions. (**c**) shows the plastic plate bonded onto the vibrator.

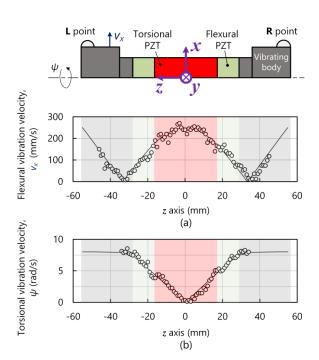


Figure 6. (a) Flexural and (b) torsional vibration velocity distributions in the *z* axis. The dots and curves indicate the measured and simulated results, respectively.

Figure 6a,b illustrate torsional and flexural vibration velocity distributions, respectively. The voltage amplitude and the working frequency was, respectively, 50 V and 4.837 kHz. Here, the driving frequency did not markedly deviate from the resonance frequency measured with the impedance analyzer; this indicates little effect of the fixing mechanism on the vibration properties [29–31]. Additionally, the experimental results well agree with the simulated ones; this implies the successful excitation of the first-order torsional and first-order flexural vibrations.

4. Load Characteristics

Figure 7 illustrates the testbed for measuring the T/F motor's load characteristics. The preload was orthogonally applied to the vibrator by tightening the coil springs. The torque τ was measured by pulling up the weight and the rotation speed θ was estimated with a rotation meter (HG200, Aero-top Hitech Corporation, Beijing, China). The input power P_{in} was measured with power meters (3332, Hioki Corporation, Nagano, Japan). The efficiency is that

$$\eta = \frac{\tau \cdot \theta}{P_{in}} \times 100\%. \tag{5}$$

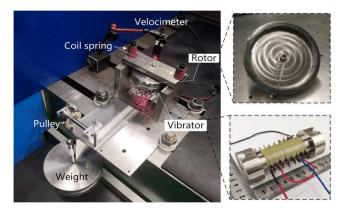


Figure 7. Experimental setup for measuring the load characteristics.

Figure 8 shows the variation in no-load rotation speed versus the phase shift between the voltages applied to the torsional and flexural vibration. The no-load rotation speed reached the peak values along the counter-clockwise and clockwise directions when the phases were -110° and 70° , respectively.

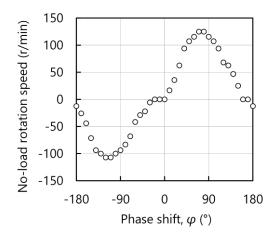


Figure 8. No-load rotation speed as a function of the phase shift of the voltages applied to torsional and flexural vibrations. The voltages applied to the flexural and torsional PZT disks have the phase shift φ .

Figure 9a–c plot the variations in rotation speed, output power, and efficiency versus torque, respectively. The working frequency and the preload were set to 4.678 kHz and 600 N, respectively. The working frequency was slightly smaller than the resonance frequency dominantly because the vibration loss and the PZT's elastic compliance become larger with increasing voltage [32,33]. When the voltage had the zero-to-peak amplitude of 250 V, the motor yielded the maximal rotation speed of 125 r/min and the torque of 4.3 Nm. Meanwhile, the maximal output power and the maximal efficiency were, respectively, 16.9 W and 10.5%. At 200 V, the no-load rotation speed decreased to 97 r/min. The efficiency was <5% probably because the frictional loss became larger.

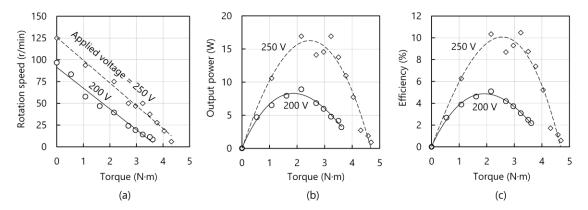


Figure 9. Load characteristics: (a) rotation speed, (b) output power, and (c) efficiency vs. torque.

Table 3 compares the T/F motor's performance with several other motors based on d₁₅ modes. Compared to Yan's micro motor working in the T/T modes [2], the T/F motor exhibited high torque density and high power density despite its large profile. In addition, Satonobu's T/L motor provided larger torque because the longitudinal vibration produces large normal force, which may enhance the frictional force to drive the rotor [5]. However, the T/L motor provided lower output power than the T/F one because undesirable vibrations exist on the vibrator especially when the diameter-to-length ratio approaches one. Additionally, in the fabrication of torsional PZT, several wedge-shaped PZT pieces should be independently prepared and bonded with each other; this procedure causes higher expense of making torsional PZT than other types particularly in the small batch of producing torsional PZT.

However, the production expenses of a large amount of the T/F motors was assumed to have a small difference compared to other types of piezoelectric motors as the workloads for making a large batch of PZT disks in d_{33} , d_{31} , and d_{15} modes did not have significant difference. Thus, these comparisons allow us to infer that superposition of torsional and flexural vibrations is suitable for piezoelectric motors.

	This Motor	[2]	[5]
Vibration modes	T/F	T/T	T/L
Structure	Cylindrical	Square bar	Cylindrical
Dimension (mm)	$\Phi 45 \times 110$	$1.55 \times 0.75 \times 0.75$	$\Phi 60 \times 72$
Weight (kg)	0.89	1.9×10^{-5}	1.1
Maximal torque (Nm)	4.3	2.5×10^{-6}	8
No-load rotation speed (r/min)	125	700	-
Maximal output power (W)	16.9	8.3×10^{-5}	3
Maximal efficiency	10.5%	_	<15%
Torque density (Nm/kg)	4.8	0.1	7.3
Power density (W/kg)	19.0	4.3	2.7

Table 3. Performance comparison between our motor and those driven with torsional modes.

5. Conclusions

In this study, we produced a dumbbell-shaped vibrator with the first-order torsional and first-order flexural modes, and used it to form a piezoelectric motor, where the maximal torque, the maximal rotation speed, and maximal output power reached 4.3 Nm, 125 r/min, and 16.9 W, respectively. These results imply the suitability of torsional/flexural vibrations for piezoelectric motors. In further studies, we would like to, as mentioned above, increase the working frequency into the inaudible range by shortening the vibrator's length [34,35]. Additionally, we will model the T/F motor and conduct structural optimization in the future [34,36].

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