

Article

Orientation Dependence of High Cycle Fatigue Behavior of a $\langle 111 \rangle$ Oriented Single-Crystal Nickel-Based Superalloy

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Abstract: High cycle fatigue failure has been recognized as one of the major forms of failure of aero-engine blades. This paper presents the high cycle fatigue testing of a Ni-based superalloy near $\langle 111 \rangle$ orientation at 800 °C. The fracture morphology and dislocation configuration were analyzed in detail by scanning electron microscopy (SEM) and transmission electron microscopy (TEM) to indicate the influence of orientation deviation degree on the high cycle fatigue properties. The results show that the orientation deviation significantly affects the initiation of the slip systems, which is closely related to fatigue performance. The best fatigue life appears on the precise $\langle 111 \rangle$ orientation, and the deformation behavior is controlled by multiple sets of equivalent $\langle 110 \rangle$ $\{111\}$ slip systems. With the increase in orientation deviation, the fatigue properties of the alloy degenerate significantly. On the boundary of $\langle 111 \rangle$ - $\langle 001 \rangle$, two groups of $\langle 110 \rangle$ $\{111\}$ slip systems with the maximum Schmid shear stress dominate the deformation behavior. On the other hand, on the $\langle 111 \rangle$ - $\langle 011 \rangle$ boundary, the formation of stacking faults and rapid cutting of γ' precipitates results in a negative effect on the fatigue life.

Keywords: superalloys; high cycle fatigue; crystal orientation



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1. Introduction

Nickel-based single-crystal superalloys have been widely used as components in the hot end of advanced aircraft engines and gas turbines due to their excellent high-temperature strength, oxidation resistance and creep properties [1,2]. In the actual service process, turbine blades are often subjected to fatigue damage with cyclic loading, except for creep damage caused by constant centrifugal stress. Among the many factors that cause engine failure, high cycle fatigue (HCF) has proven to be the most common reason [3].

In the past few decades, high cycle fatigue has attracted much attention because of numerous influential factors. Temperature [4,5], frequency [6], load ratio [7], casting conditions [8] and trace elements [9] all have effects on the initiation and propagation of fatigue cracks. However, most of these studies focus on a specific [001] crystal orientation [10–13]. Vibration stresses resulting in fatigue damage scarcely follow a single axial direction when applied to turbine blades with complex structures. Therefore, crystal orientation, as a significant factor affecting mechanical properties, is noteworthy in the study of high cycle fatigue.

Previous research on the anisotropy of nickel-based single-crystal superalloys mainly focused on creep properties [14–17]. In Leverant and Kear's [18] paper, the [111] orientation of Mar M200 exhibits superior creep performance compared to [001] at 760 °C/689 MPa because of the lower Schmid factor of $\langle 112 \rangle$ $\{111\}$ slip system. MacKay and Meier's early study on MarM247 SX superalloys indicated that orientation deviation significantly affected the creep properties near [001] orientation [19]. Creep lives of crystals near the $\langle 001 \rangle$ - $\langle 011 \rangle$

boundary of the stereographic triangle are significantly longer than that of the $\langle 001 \rangle$ - $\langle 111 \rangle$ boundary. Further studies of CMSX-4 by Sass [20] and Matan [21] showed that anisotropy of the primary creep rate is closely related to the number of dominant slip systems and the magnitude of shear stress. In general, the difference in the activation process of the slip systems during plastic deformation caused by different crystal orientation is considered to be the main reason for the creep anisotropy at immediate temperatures. However, the anisotropy of high cycle fatigue properties has only received limited attention.

In this paper, the high cycle fatigue behaviors of a Ni-based single-crystal superalloy with different orientations near $\langle 111 \rangle$ orientation were studied. Both coplanar double-slip orientation and non-coplanar double-slip orientation were considered to explain the orientation influencing factors of high cycle fatigue behaviors. Characteristics and mechanisms of fatigue fracture caused by orientation deviation angle and direction are revealed from macroscopic fracture mode and microscopic dislocation movement perspectives.

2. Materials and Methods

A model nickel-based single-crystal superalloy with nominal composition (wt.%) of (7.4–7.8) Al, (3.0–4.5) Ta, (8.9–9.3) Mo, (1.3–2.7) Cr, (1.4–1.6) Re and Ni as the balance, is used as the research object. The master alloy is prepared by vacuum induction melting, and the single-crystal (SX) rods with different deviations from $[\bar{1}11]$ were prepared by high-rate solidification (HRS) directional solidification with the seed crystal method. Twenty-millimeter-long seeds with the same desired crystal orientation as the axial direction of test bars were cut directionally from the single-crystal specimens and placed at the bottom of the ceramic mold. Then, the master alloy was melted at 1600 °C and poured on the seeds. Finally, the ceramic mold was withdrawn from the furnace at a withdrawal rate of 3.5 mm/min. The crystal orientations were determined by Laue X-ray diffraction, as shown in Figure 1a. In this paper, test rods with four types of crystal orientations are chosen for this work, which are precise $[\bar{1}11]$ orientation, labeled A, 14° deviations from $[\bar{1}11]$ orientation along the $\langle 111 \rangle$ - $\langle 001 \rangle$ boundary, labeled B, and 10° and 15° deviations from $[\bar{1}11]$ orientation along the $\langle 111 \rangle$ - $\langle 011 \rangle$ boundary marked as C and D, respectively. After a full heat treatment with 1300 °C/2 h + 1310 °C/2 h + 1320 °C/2 h + 1330 °C/6 h + 1340 °C/10 h, air cooling (AC) + 1080 °C/2 h and AC + 870 °C/24 h, these rods are processed into test specimens with a gauge diameter of 4 mm and a length of 52 mm, as shown in Figure 1c. Two fatigue samples were machined for each orientation from the same bar.

All tests were performed in 800 °C/ $\sigma_{\max} = 445$ MPa under load control at a stress ratio $R = \sigma_{\min}/\sigma_{\max} = -1$ and a frequency of 83.3 Hz. The specimens were loaded with sinusoidal wave by the rotating bending method. The fracture surfaces of specimens were observed by EVO-10 scanning electron microscope (SEM, Zeiss, Jena, Germany). Foils for the transmission electron microscope test were cut by a wire cutting machine along the direction perpendicular to the principal stress axis near the fatigue fracture surface, then mechanically ground and electro-chemically thinned at 20 °C in a solution (10 vol% perchloric acid in methanol). Whereafter, the dislocation configurations were analyzed by JEM-2100F transmission electron microscope (TEM, Tokyo, Japan). Observation regions of TEM were selected near the crack source to accurately represent the plastic deformation behavior.

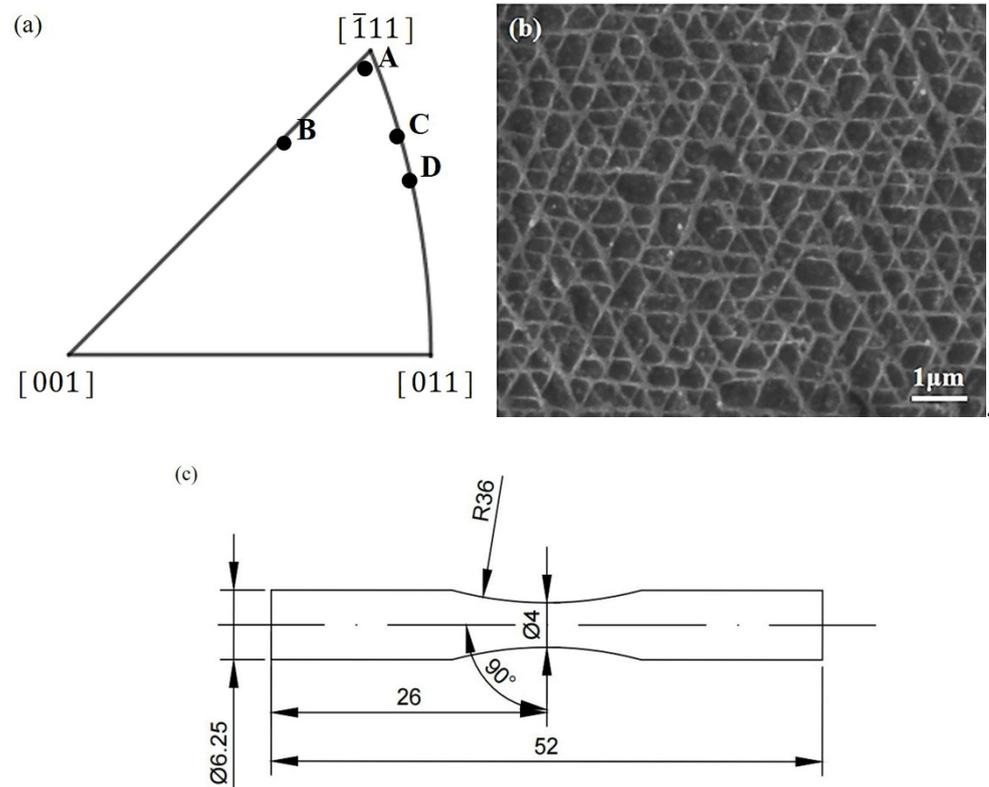


Figure 1. (a) Orientations of specimens A–D in the standard stereographic triangle. (b) Secondary electron SEM micrograph of initial γ/γ' microstructure with the crystallographic orientation $\langle 111 \rangle$ after the standard heat treatment. (c) Schematic of specimen for high cycle fatigue test (unit: mm).

3. Results

3.1. HCF Properties

The results from HCF tests at $800\text{ }^{\circ}\text{C}/\sigma_{\max} = 445\text{ MPa}$ for specimens with different crystal orientations are presented in both Table 1 and Figure 2. It is obvious that the HCF properties of this alloy exhibited significant small-angle deviation sensitivity. The precise $[111]$ orientation showed the best HCF performance over 107 cycles. With the increased orientation deviation, the HCF properties degraded gradually, while the degree of decline varied according to the deviation direction. Fatigue life of specimens deviated by 15° from $[111]$ towards $[001]$ and decreased to an average of 2.42×10^6 cycles, whereas misorientations of 10° and 14° towards $[011]$ directions resulted in a rapid decline in fatigue performance with 1.28×10^6 and 6.76×10^5 cycles, respectively.

Table 1. $850\text{ }^{\circ}\text{C}/\sigma_{\max} = 445\text{ MPa}$ high cycle fatigue (HCF) properties for different orientations to the model alloy.

Orientation	Sample 1	Sample 2	Average
A	9.99×10^6	$>1 \times 10^7$	1×10^7
B	2.48×10^6	2.36×10^6	2.42×10^6
C	1.08×10^6	1.48×10^6	1.28×10^6
D	6.55×10^5	6.96×10^5	6.76×10^5

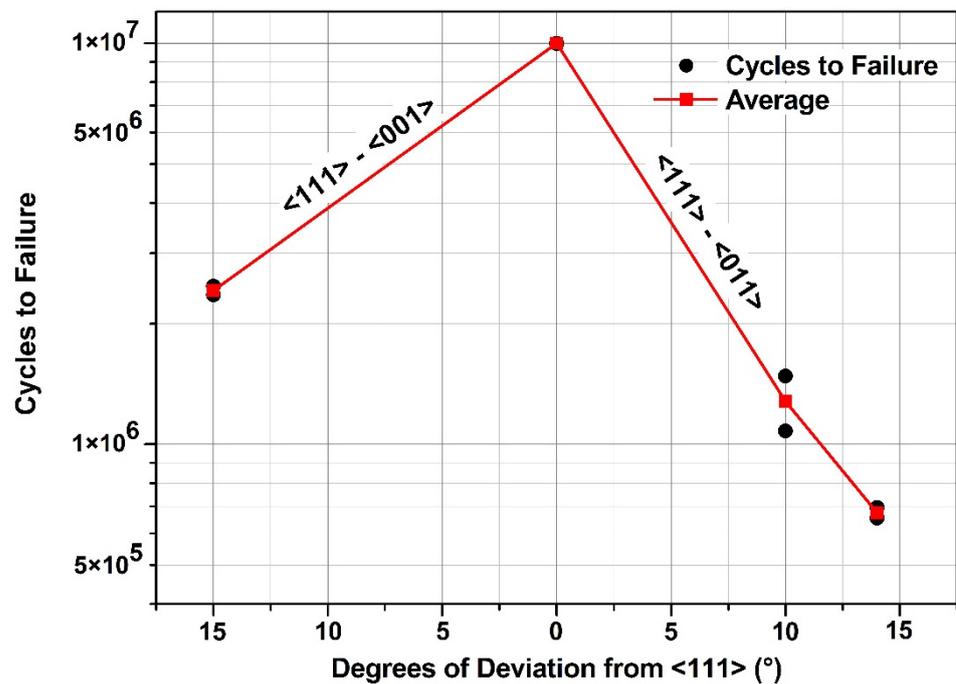


Figure 2. Effect of orientation deviation from $\langle 111 \rangle$ on the $850\text{ }^{\circ}\text{C} / \sigma_{\max} = 445\text{ MPa}$ high cycle fatigue life.

3.2. SEM Fractography

SEM observations showed that mixed cleavage and tearing fracture modes exist in the vicinity of $\langle 111 \rangle$ orientation. Due to the rotary bending loading mode, the surface of the specimens bears the maximum cyclic stress, which results in the initial cracks mostly at or near the surface. Fractures of specimens with different orientations consist of crack initiation stage, crack propagation stage and final rupture stage, similar to numerous previous studies. However, with the change of deviation direction and angle, the fracture morphology exhibits obvious differences as shown in Figure 3a–d.

The observations from Figure 3a show typical multi-source fracture characteristics in precise $\langle 111 \rangle$ orientation. Fatigue cracks generate on the surface and propagate uniformly to the interior along multiple cleavage facets with distinct river pattern and obvious crystallographic characteristics. Many researchers have reported the stage 1 crack growth behavior in nickel-based SX superalloys. These smooth cleavage fracture facets formed during the stage of crack growth and are significantly dependent on the slip planes. Due to face-centered cubic structures, the dominating slip system of nickel-based SX superalloys operated at elevated temperatures is $\langle 110 \rangle \{111\}$ and $\langle 112 \rangle \{111\}$, hence crack propagation will follow the habit plane $\{111\}$ [22–24]. Fatigue steps are observed at the intersection of different cleavage surfaces as a result of the cross slip of dislocations under cyclic stress.

On the boundary of $\langle 111 \rangle$ - $\langle 001 \rangle$, with the orientation deviating from $\langle 111 \rangle$, although multi-source fracture patterns are still observed, crack propagation at stage 1 is not strictly uniform along different $\langle 111 \rangle$ planes, but mainly on two types of cleavage planes connected with ridges. Figure 4 shows the schematic diagram of $\{111\}$ planes on which the different cleavage planes are located. In order to further determine the type of $\{111\}$ plane, the method of calculating the angle relationship is adopted, as shown in Figure 4c. By measuring the angular relationship between the stress axis and the normal direction of the cleavage plane after cutting the sample longitudinally, one can observe that both fracture surfaces are 63° off the stress axis, which is close to the (111) and $(1\bar{1}\bar{1})$ planes (64°). Cracks propagate inward along these cleavage planes, resulting in a gradual decrease in the bearing capacity area and eventual failure under tensile overload. The occurrence of river

patterns provides evidence for cyclic plastic deformation, and the variation of distance between slip trace indicates that plastic deformation accelerates with crack growth.

A distinct fracture mode occurs when the orientation deviates towards [011] along the $\langle 111 \rangle$ - $\langle 011 \rangle$ boundary. Both Figure 3c,d show a unitary crack source regardless of whether misorientation is 10° or 14° . Crack propagation at a single cleavage plane with river patterns and stepwise features can be observed simultaneously. Angular relation calculations indicate that the fracture surface is (111) plane. It was also verified by Laue X-ray diffraction after polishing the cleavage plane of the sample, and the result was added in Appendix B. The fatigue fracture consists of a dominant stage 1 and a transient stage 2. Further comparison of Figure 3a–d indicates that with the increase in orientation deviation, the fracture pattern with a single cleavage plane becomes more evident. Moreover, this transformation of deformation behavior results in rapid degradation of fatigue performance.

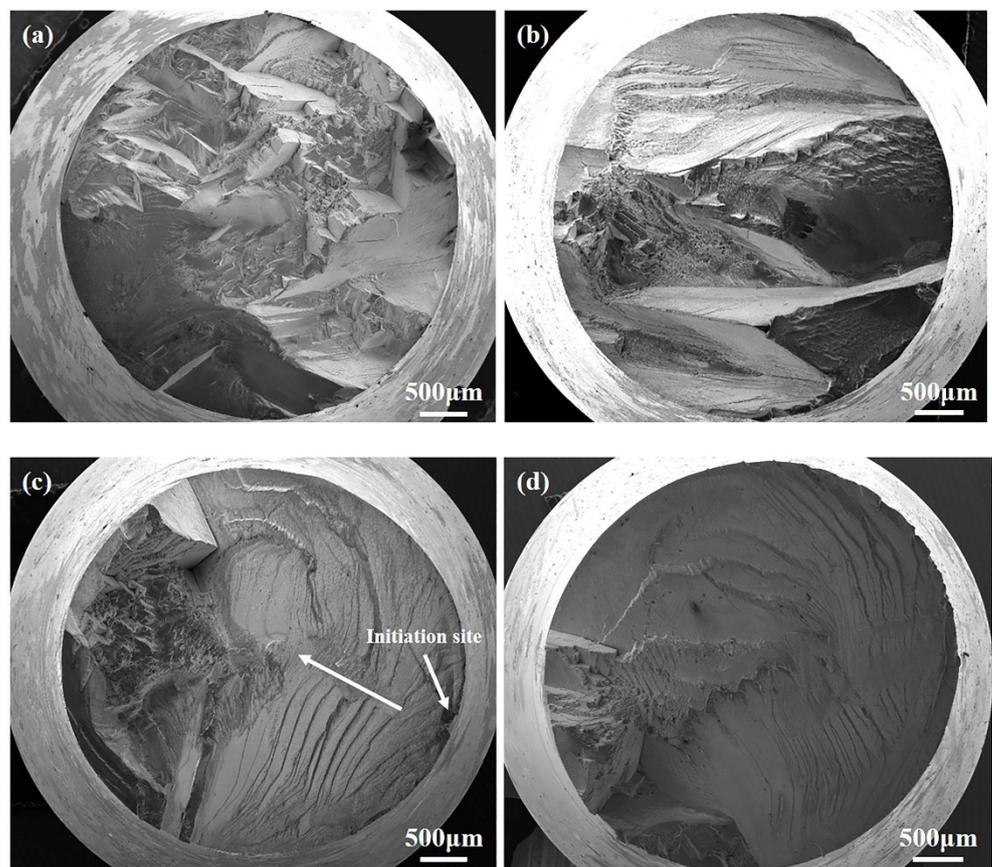


Figure 3. SEM images of fracture surface of specimens with different orientations tested at 800°C . (a) Precise $\langle 111 \rangle$ orientation A. (b) Orientation B on the $\langle 111 \rangle$ - $\langle 001 \rangle$ boundary. (c,d) Orientation C and D on the $\langle 111 \rangle$ - $\langle 011 \rangle$ boundary.

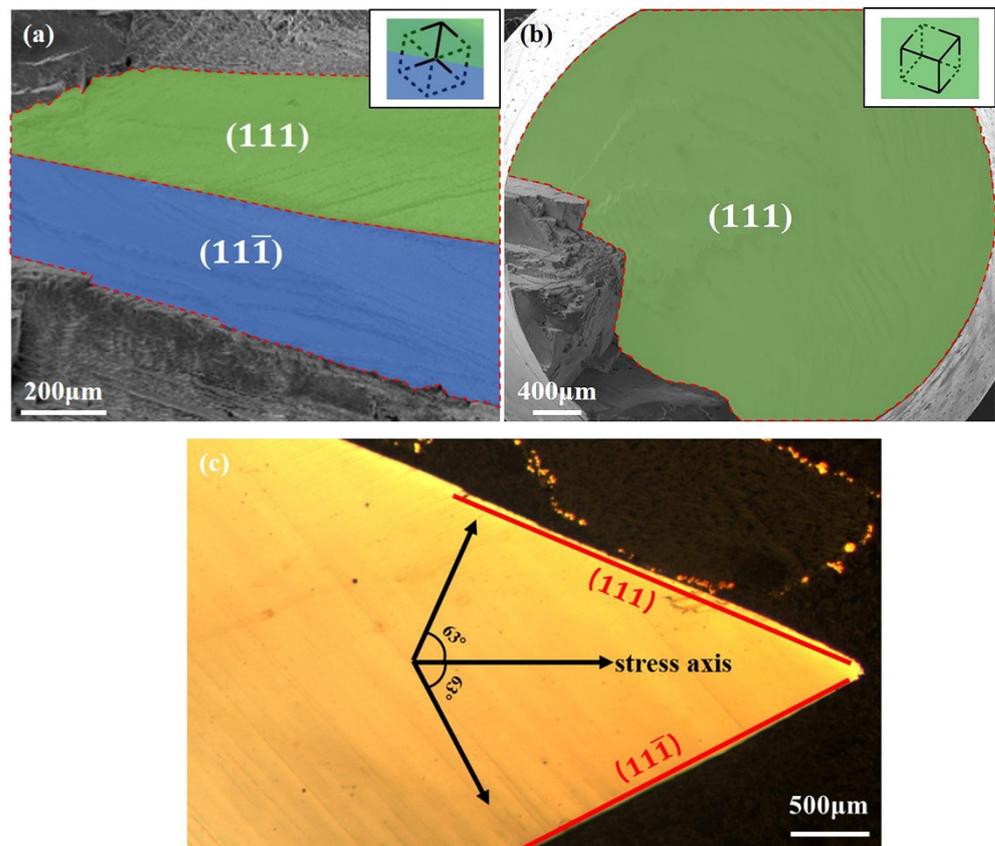


Figure 4. Diagram of slip planes at fracture surface. (a) Orientation B. (b) Orientation D. (c) Angle calculation diagram of B-oriented sample.

3.3. Dislocation Configurations

TEM micrographs in Figure 5a show typical deformation substructures from the samples with orientation A after HCF failure. Under this test condition, the γ/γ' microstructures after fracture were roughly the same as that of the heat treated. The distribution of dislocations in the two phases is not uniform, but mainly concentrated in the γ matrix. Several groups of dislocations with different Burgers vectors in the matrix deposit on the γ/γ' interface, forming dislocation networks. A few dislocations cut into the γ' precipitate in the form of $a/2 \langle 110 \rangle$, and no stacking faults cutting was observed. The basic characteristic of a nickel-based superalloy structure is that the cubic γ' phase is embedded in the γ matrix, resulting in the formation of three pairs of equivalent γ channels around the γ' cubic precipitate. When the applied stress is in the precise $\langle 111 \rangle$ orientation, three types of γ channels have the same stress state and similar dislocation distribution.

In the B oriented specimen, the dislocation configuration is similar to that of the sample with A orientation as shown in Figure 5b. Two types of $a/2 \langle 110 \rangle$ dislocations were observed in the γ matrix and at the γ/γ' interface, while there were also very few present in the γ' precipitates. The nature of the dislocations cutting into the γ' precipitate was analyzed in detail under two-beam condition as shown in Figure 6. Typical dislocations **b1** and **b2** as indicated by the arrows are presented in Figure 6a–f. Several sets of \mathbf{g} diffraction conditions have been used to form images. Both of these dislocations are visible under $\mathbf{g} = 220$. The difference is that the **b1** dislocation is visible under $\mathbf{g} = \bar{2}02$ operation reflection but out of contrast in the $\mathbf{g} = 02\bar{2}$ condition. Since $\mathbf{g} \cdot \mathbf{b} = 0$, the Burgers vector of **b1** is inferred to be $a/2[011]$. Similarly, the comparative analysis of dislocation **b2** indicates that it is visible under $\mathbf{g} = 02\bar{2}$, while invisible under $\mathbf{g} = \bar{2}02$ diffraction condition. The Burgers vector of **b2** is $a/2[101]$. According to the characterization of different operation reflections,

the actuation of the slip system is obviously changed with the crystal orientation deviating from $\langle 111 \rangle$ direction.

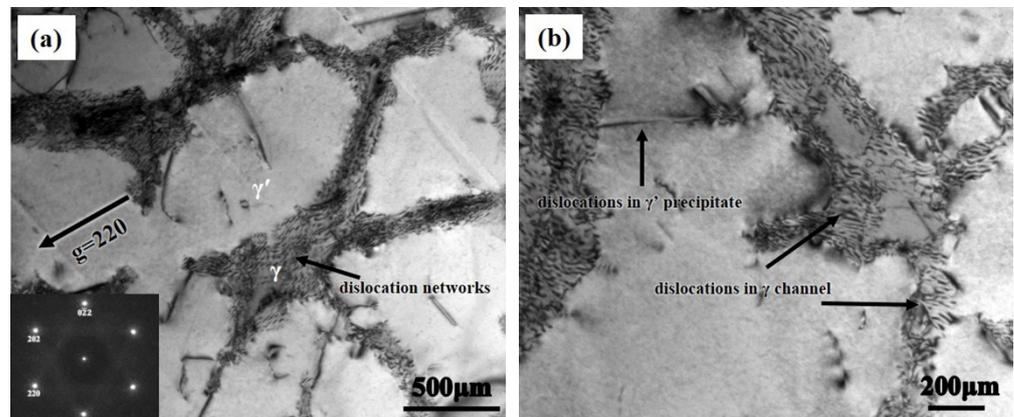


Figure 5. TEM images near fatigue source showing the dislocation configuration in the γ matrix and γ' precipitates of specimens ruptured at $800\text{ }^\circ\text{C}/\sigma_{\text{max}} = 445\text{ MPa}$. (a) Dislocation networks with no stacking fault in orientation A. (b) Irregular dislocation networks and a small amount of $a/2\langle 011 \rangle$ dislocations cutting into γ' precipitates in orientation B.

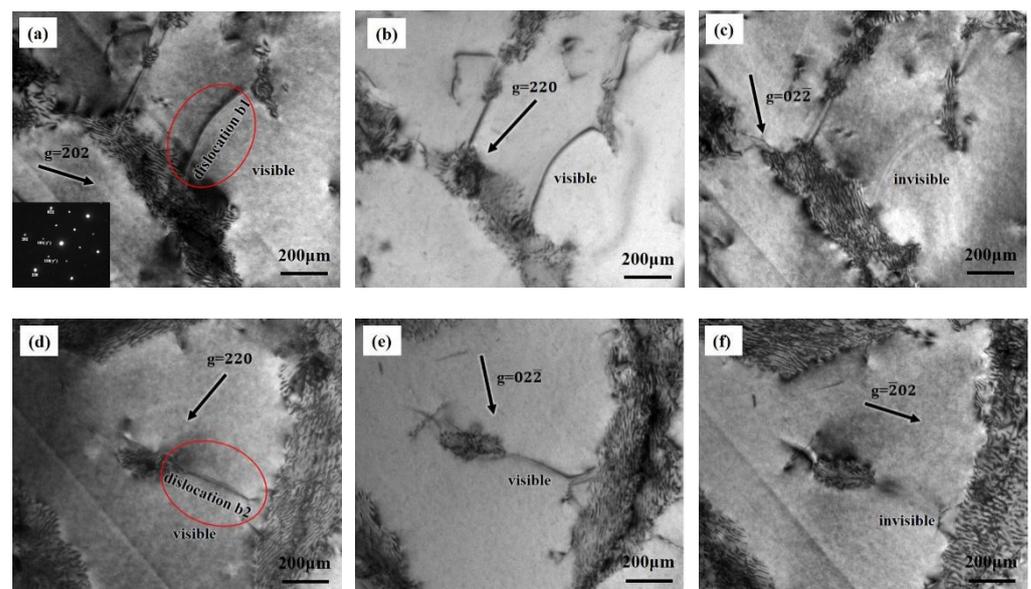


Figure 6. Two-beam TEM bright field images of typical dislocations labeled b1 and b2 in γ' precipitates of orientation B after rupture. (a) Dislocation b1 with $g = \bar{2}02$ diffraction condition. (b) $g = 220$. (c) $g = 02\bar{2}$. (d) Dislocation b2 with $g = 220$ diffraction condition. (e) $g = 02\bar{2}$. (f) $g = \bar{2}02$.

Compared with orientation A and B, the differences of samples on the $\langle 111 \rangle$ - $\langle 011 \rangle$ boundary are striking. Orientation deviation causes a variation of the stress state in the three different types of matrix channels, result in significant differences in the number of dislocations in different channels as shown in Figure 7. Meanwhile, in addition to high-density dislocations in γ channels, there are many unidirectional stacking faults associated with $\langle 112 \rangle$ cutting both the γ matrix and γ' precipitates as presented in Figure 6b, which indicates that $\langle 112 \rangle \{111\}$ slip systems participate in the deformation instead of being dominated only by $\langle 110 \rangle \{111\}$ slip systems.

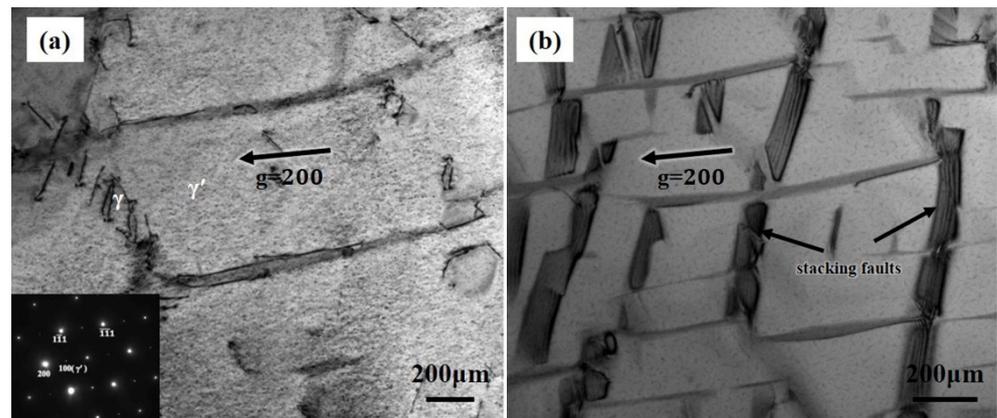


Figure 7. TEM micrographs of orientation C after rupture. (a) Dislocation density differences in different types of matrix channels. (b) Dense stacking faults along the same direction.

4. Discussion

Although HCF differs from creep in the manner of stress loading, both of them exhibit plastic behavior below yield strength. In this paper, the analysis of plastic behavior is mainly focused on the crack initiation on the sample surface to avoid the complex stress state in the late deformation. The initiation and propagation of fatigue cracks are the result of periodic slip of the slip system, that is, cyclic plastic deformation caused by dislocation movement. At elevated temperatures, the anisotropic behavior of nickel-based SX superalloys is closely related to the critical shearing stress of $\langle 110 \rangle \{111\}$ and $\langle 112 \rangle \{111\}$ slip systems [25]. Schmid factors of different specimens are shown in the Table 2. The value of the Schmid factor represents the ease with which the slip system is operated. It should be noted that the Schmid factors for all slip systems are not listed here. Only $\langle 110 \rangle \{111\}$ slip systems with the largest values and the $\langle 112 \rangle \{111\}$ slip systems that can be generated by the reaction are represented. Complete data for all slip systems are added in Appendix A. For the precise $\langle 111 \rangle$ orientation, TEM photographs in Figure 4a show that dislocation movement was mainly concentrated in the γ channel with few γ' precipitates cutting. It was indicated that plastic deformation was mainly controlled by the $\langle 110 \rangle \{111\}$ octahedral sliding system, since the shear stress of the $\langle 112 \rangle \{111\}$ sliding system on $\langle 111 \rangle$ orientation was relatively low. Six equivalent slip systems operated uniformly on three $\langle 111 \rangle$ planes, which was consistent with the macroscopic fracture. During the deformation process operated by multiple $\langle 110 \rangle \{111\}$ slip systems, dense dislocation networks were formed at the γ/γ' interface through a continuous cross-slip and reaction of $a/2[110]$ dislocations with different Burgers vectors. Previous studies [26] on creep behavior have shown that dense interfacial dislocation networks prevent dislocations in the γ matrix from cutting γ' precipitates, resulting in an obvious work hardening, which is an important reason for the transition from the first creep stage to the second. In the $[111]$ oriented sample, since the Schmid factors of initial slip and cross slip planes are almost equivalent, the cross slip of screw dislocations in matrix channels occurs easily under the cyclic stress [27], which promotes the generation of γ/γ' interface dislocation networks. Furthermore, $\langle 110 \rangle \{111\}$ slip systems have the lowest shear stress in $\langle 111 \rangle$ direction. Both the low resolved shear stress and interfacial dislocation networks can contribute to resistance to the movement of dislocations, resulting in the best high cycle fatigue performance of $\langle 111 \rangle$ specimens.

Table 2. Calculated Schmid factors for specimens A–D on part of the most potentially activated slip systems.

Orientation	Slip Systems	Schmid Factors
A	($\bar{1}11$) [110]	0.272
	($\bar{1}11$) [101]	0.272
	(11 $\bar{1}$) [112]	0.314
B	(111) [$\bar{1}01$]	0.380
	(11 $\bar{1}$) [011]	0.380
	($\bar{1}\bar{1}1$) [$\bar{1}12$]	0.088
C	(111) [$\bar{1}\bar{1}0$]	0.373
	(111) [$\bar{1}01$]	0.373
	(111) [$\bar{2}11$]	0.431
D	(111) [$\bar{1}\bar{1}0$]	0.405
	(111) [$\bar{1}01$]	0.405
	(111) [$\bar{2}11$]	0.467

As can be seen from Table 2, misorientation from $\langle 111 \rangle$ results in an obvious difference in the actuation of slip systems. Furthermore, (111) [$\bar{1}01$] and (11 $\bar{1}$) [011] slip systems dominate when the orientation deviates towards $\langle 001 \rangle$ along the $\langle 111 \rangle$ - $\langle 001 \rangle$ boundary. No stacking faults were observed according to the results of TEM, which reveals that the dominant slip system was also $\langle 110 \rangle$ {111} despite a slight increase in Schmid factor of $\langle 112 \rangle$ {111} in this direction. Previous research on creep anisotropy of nickel-based SX superalloys at elevated temperatures has suggested that stacking fault shear is frequently associated to a reaction of $a/2 \langle 110 \rangle$ {111} dislocations in the γ matrix [28]. A typical reaction might be written as:

$$a/2[011] + a/2[101] = a/3[112] + a/6[112]$$

For sufficiently applied stress, the $a/3\langle 112 \rangle$ dislocation may cut into γ' accompanied by rapid plastic deformation, while $a/6\langle 112 \rangle$ dislocations remain at the γ/γ' interface. Note that the reactions require appropriate dislocations rather than two arbitrary dislocations [29,30]. Considering the orientations along the $\langle 111 \rangle$ - $\langle 001 \rangle$ boundary, the most highly stressed dislocations $a/2[\bar{1}01]$ (111) and $a/2[011]$ (11 $\bar{1}$) with the Schmid factor of 0.38 can combine to form $a/3[\bar{1}12]$ (1 $\bar{1}\bar{1}$) and $a/6[\bar{1}12]$ (1 $\bar{1}\bar{1}$). However, a finite Schmid factor of 0.088 is inadequate to obtain $a/3[\bar{1}12]$ (1 $\bar{1}\bar{1}$) dislocations shear into γ' precipitates. In addition, $\langle 110 \rangle$ dislocations on different {111} planes are not easily combined to generate $\langle 112 \rangle$ dislocations. Macroscopic fractures are characterized as stage 1 fatigue along two major $\langle 111 \rangle$ planes, which provide evidence for the operation of (111) [$\bar{1}01$] and (11 $\bar{1}$) [011] from another side. Meanwhile, compared with the precise $\langle 111 \rangle$ orientation, the dominant $\langle 110 \rangle$ {111} slip systems subjected to greater shear stress result in a decline in fatigue performance.

Instead of uniform deformation on multiple slip planes, a coplanar double slip occurs when the orientation deviates from precise $\langle 111 \rangle$ by 10° towards $\langle 011 \rangle$ along the $\langle 111 \rangle$ - $\langle 011 \rangle$ boundary. In other words, two dominant $\langle 110 \rangle$ {111} slip systems with maximum shear stress are located on the same plane. TEM observations in Figure 6 indicated that unidirectional stacking faults shear γ' is the dominative mechanism of plastic deformation. According to the analysis above, coplanar (111) [$\bar{1}01$] and (111) [$\bar{1}\bar{1}0$] slip systems can participate in the generation of (111) [$\bar{2}11$] dislocation distribution in the same (111) plane. Furthermore, the generated [$\bar{2}11$] dislocation was dominant in the dodecahedral slip system, which can continue to form stacking fault bands cutting into γ' precipitates. The resulting rapid plastic deformation promotes crack initiation and propagation and leads to a rapid decline in fatigue life.

5. Conclusions

The present study is focused on the effect of orientation deviation near $\langle 111 \rangle$ orientation on high cycle fatigue properties of a nickel-based single-crystal superalloy. High cycle fatigue tests were carried out at 800 °C ($\sigma_{\max} = 445$ MPa), and the fatigue deformation mechanism was studied from both macroscopic and microscopic levels, respectively. The conclusions are summarized as follows:

- (1) The precise $\langle 111 \rangle$ orientation showed the best fatigue performance. With the increase in orientation deviation, the fatigue properties degenerate significantly.
- (2) Fractographic investigation revealed that, under testing conditions, the precise $\langle 111 \rangle$ orientation shows a multi-source fracture mode, and its deformation is mainly controlled by multiple sets of equivalent $\langle 110 \rangle \{111\}$ slip systems in three γ channels with the same stress state.
- (3) On the boundary of $[111]$ - $[001]$, although the alloy still maintains the multi-source fracture mode, the dominant slip systems are changed into two $\langle 110 \rangle \{111\}$ slip systems with maximum Schmid shear stress.
- (4) As the orientation deviates towards $[011]$ along the $[111]$ - $[011]$ boundary, the fracture mode of the alloy changes from multi-source to single-source. Stacking faults formed by two $\langle 110 \rangle \{111\}$ dislocations cutting into γ' precipitates result in rapid plastic deformation and has a negative effect on the fatigue life.

Author Contributions: S.L. and S.G. conceived the project and designed experiments. B.H. and Y.P. carried out most of the experiments. B.H. and S.L. wrote the paper. All authors contributed to the interpretation and presentation of the results. All authors have read and agreed to the published version of the manuscript.

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Conflicts of Interest: There are no conflict to declare.

Appendix A

Table A1. Schmid factors for $\langle 110 \rangle \{111\}$ slip systems of specimens A–D.

Slip System	A	B	C	D
$(111) [\bar{1}\bar{1}0]$	0.157	0.285	0.373	0.405
$(111) [\bar{1}0\bar{1}]$	0.157	0.380	0.373	0.405
$(111) [0\bar{1}\bar{1}]$	0.314	0.095	0.000	0.000
$(\bar{1}\bar{1}1) [110]$	0.000	0.000	0.152	0.218
$(\bar{1}\bar{1}1) [101]$	0.000	0.209	0.152	0.218
$(\bar{1}\bar{1}1) [0\bar{1}\bar{1}]$	0.000	0.209	0.000	0.000
$(\bar{1}\bar{1}\bar{1}) [011]$	0.314	0.076	0.221	0.187
$(\bar{1}\bar{1}\bar{1}) [110]$	0.157	0.000	0.038	0.045
$(\bar{1}\bar{1}\bar{1}) [\bar{1}0\bar{1}]$	0.157	0.076	0.184	0.142
$(11\bar{1}) [\bar{1}\bar{1}0]$	0.157	0.285	0.184	0.142
$(11\bar{1}) [101]$	0.314	0.095	0.038	0.045
$(11\bar{1}) [011]$	0.157	0.380	0.221	0.187

Table A2. Schmid factors for $\langle 112 \rangle$ $\{111\}$ slip systems of specimens A–D.

Slip System	A	B	C	D
(111) $[\bar{1}\bar{1}2]$	0.000	0.274	0.215	0.234
(111) $[\bar{1}2\bar{1}]$	0.000	0.110	0.215	0.234
(111) $[2\bar{1}\bar{1}]$	0.000	0.384	0.431	0.467
$(\bar{1}\bar{1}1)$ $[1\bar{1}2]$	0.157	0.241	0.088	0.126
$(\bar{1}\bar{1}1)$ $[12\bar{1}]$	0.157	0.121	0.088	0.126
$(\bar{1}\bar{1}1)$ $[211]$	0.314	0.121	0.175	0.251
$(1\bar{1}\bar{1})$ $[\bar{1}\bar{1}2]$	0.157	0.088	0.234	0.190
$(1\bar{1}\bar{1})$ $[121]$	0.314	0.044	0.149	0.134
$(1\bar{1}\bar{1})$ $[21\bar{1}]$	0.157	0.044	0.084	0.056
$(11\bar{1})$ $[112]$	0.314	0.274	0.149	0.134
$(11\bar{1})$ $[\bar{1}21]$	0.157	0.384	0.234	0.190
$(11\bar{1})$ $[2\bar{1}1]$	0.157	0.110	0.084	0.056

Appendix B

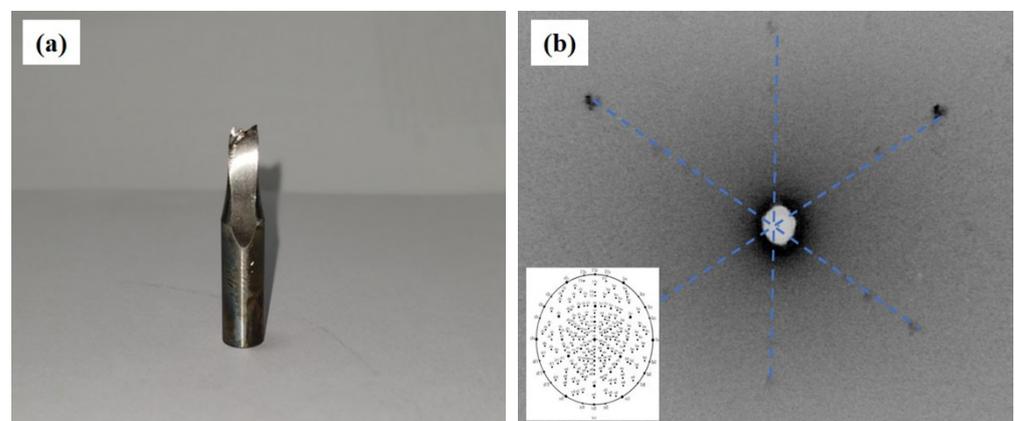


Figure A1. Laue diffraction spot pattern of a C-oriented sample in a single slip system. (a) The polished surface of the sample. (b) Laue's diffraction spot pattern.

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