

## Article

# Increasing Low-Temperature Toughness of 09Mn2Si Steel through Lamellar Structuring by Helical Rolling

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**Abstract:** The aim of the paper was to investigate the helical rolling parameters (a number of passes) for the microstructural modification and the low-temperature impact toughness improvement of the 09Mn2Si High Strength Low-Alloyed (HSLA) steel. In order to achieve this purpose, work spent to crack initiation and propagation was analyzed and compared with patterns of fracture surfaces. The microstructure and impact toughness values were presented in the temperature range from +20 to −70°C. Also, the fracture mechanisms in individual regions on the fracture surfaces were discussed. In addition, a methodology for computer simulation of the process was developed and implemented within the framework of the excitable cellular automata method and its integration with the kinetic theory of fracture. Finally, a theoretical analysis of the effect of grain shapes and orientations on the strain response patterns of a certain meso-volume simulating the material after the helical rolling was carried out.

**Keywords:** high strength low alloyed steels; helical rolling; low-temperature impact toughness; lamellar microstructure; computer simulation; excitable cellular automata method



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## 1. Introduction

The active development of the Arctic region justifies enhanced requirements for the mechanical properties of materials used for application in the cold climate. To solve this issue, in particular upon the construction of pipelines, steels of new generations (the X80 and X100 grades, or higher) are being designed, which are characterized by a continuously increasing ratio between ductility and strength [1,2]. Achieving high fracture resistance is possible due to the microstructure refinement by manufacturing processes that causes smaller bainitic grain sizes, reduced inclusion contents, etc. [3].

However, high strength low alloyed steels (HSLA) of the X52 grade, which have been used in the pipeline construction for a long time, continue to be of interest. Their properties, especially impact toughness at low climatic temperatures, have to be improved. It is known that the X65 and X70 grade steels exhibit a high ductility-to-strength ratio, whereas the X100 ones have its lower level and their fracture is, typically, rather brittle. The above metallurgical (technological) aspects of increasing impact toughness at low climatic temperatures have also been discussed in various papers [4–6]. A special role in improving this property is assigned to the bainite formation [7]. Moreover, upper bainite, formed due

to the presence of martensite–austenite (M–A) and an increased amount of low-angle lath boundaries, significantly decreases the ductility in these conditions.

To quantify crack resistance of a material, the ductile–brittle transition (DBT) temperature (DBTT) under impact bending test is applied [8]. This parameter actually characterizes the permissible operating temperature range. In [9], the microstructural aspects of embrittlement and toughening have been discussed in terms of the DBT. It has been shown that micro-voids hardly nucleate at fine M–A regions or carbides (they factually grow, but through bainite laths with high misorientation between them). Also, it is known from [10] that the DBTT is not intrinsic to a material, but to a significant extent depends on the specimen type and test conditions, which actually determine the plastic constraint.

The DBTT aspect is typically discussed in relation to materials used in nuclear power plants [11], since its development is due to aging through thermal and/or irradiation effects during long-term operation. Also, it is actively being studied for high-strength stainless pipeline steels. The key parameters that affect the DBT value are grain sizes, as well as the crack propagation direction relative to the boundaries of microstructural elements (determined by the rolling direction) [12]. If strength of specimens cut at different angles to the rolling axis has not changed significantly, the DBTT is extremely sensitive to the predominant grain orientation. Thus, the targeted formation of the boundaries of microstructural elements (especially, high-angle ones) can be an effective way to control the DBTT.

Finally, many authors have discussed the DBTT aspect for relatively inexpensive stainless steels with low nickel contents [13]. It has been shown that the DBT (as expected) is strongly affected by the crack propagation direction, depending on crystallographic texture and microstructural morphology. Moreover, the DBTT becomes extremely low in the case of fracture accompanying delamination.

Nowadays, “strain nano-technologies” [14] are being actively developed that based on the formation of micro-scale structural elements in metals and alloys, according to the ‘top-bottom’ principle. It has been shown that (a) the Hall–Petch law on the relationship between grain sizes and yield stresses ceases to be fulfilled in the dimension range below 10 nm; (b) it is not necessary to reduce the sizes of all microstructural elements below the “generally accepted” threshold of 100 nm to ensure high strength properties and maintain acceptable plasticity [15], but it is important to form a hierarchically organized microstructure in this case [16]. In addition, one of the critical factors is the strong interaction of the microstructural elements with each other. This causes an increase in strength and excludes rapid strain localization processes.

Within this formulation of the problem, a promising direction is associated with the formation of “heterogeneous materials”, which is widely discussed in recent papers in the field of materials science [17]. It should be noted that all structural materials are inherently heterogenous in one way or another. According to [18], the principles of the heterogenous microstructure formation are determined both by the available technologies for processing metals and alloys, and by the possibility of its hierarchical alignment. Enhancing strength of these materials is achieved by heterogeneity of their microstructures, crystal structures, or compositions. The domain sizes can be in the range from micro- to millimeters, and their shapes vary to form very diverse material systems. Additional factors to ensure extraordinary strength and ductility are “back-stress strengthening and back-stress work hardening”. From the standpoint of the mechanisms providing the relaxation-ability of such materials, “statistically stored dislocations and geometrically necessary dislocations” are typically discussed.

At a higher (meso-scale) level, grain shapes and sizes, as well as texture and grain size variations, are crucial. The “lamellar microstructure with elongated soft coarse-grained domains embedded in the ultrafine grain matrix” is considered to be the “optimal” grain one [19]. Thus, the principle of the hierarchical microstructure formation in the heterogeneous materials is determined by the back-stress produced due to dislocation pile-up at the domain boundaries, when their interface density is high. On the other hand, the

spacing between them should be sufficient to ensure the possibility of dislocation strains. The authors have called this principle “strain partitioning among heterogeneous domains”, when the strain gradient increase the back-stress work hardening.

Undoubtedly, one of the factors limiting the widespread industrial use of the heterogeneous materials is the lack of high-performance procedures for the formation of such a microstructure.

The above-described principle of the lamellar microstructure formation at various scale levels is clearly illustrated in [20]. It is shown that following the natural principles of the hierarchical microstructure alignment (like a seashell) enables to achieve high strength properties (with satisfactory ductility) for structural steels. This implies the formation of hierarchical sandwich microstructures, including gradient grain sizes, as well as both martensite and austenite phases. In this case, the lamellar microstructure has been observed both at the highest structural level (over the cross-sectional landscape of the material) and at the nano-scale (FCC austenite lamellae 5 nm wide alternate with 6 nm HCP martensite ones). It should be emphasized once again that the formation of such small microstructural elements does not suppress development of dislocation strain mechanisms even under conditions when grain sizes are about 20 nm.

Aspects related to the possibility of strength increasing with toughness retention (strength vs. toughness or hardness vs. ductility) for materials of various classes are discussed in a large number of papers (for example [21,22]) using the concepts of fracture mechanics. In [23], transcrystalline fracture patterns are shown from the perspective of various scale levels as applied to structural steels. At the micro-level, it is proposed to use local fracture stress values to assess the local “work-of-fracture” parameter. Furthermore, at the nano-scale level, ahead of a sharp crack is proposed “to increase the compliance by a cooperative movement of atoms (involving extra work) to allow the crack-tip bond to displace sufficiently”. These representations are consistent with the plastic distortion mechanism considered in V.E. Panin’s paper [24].

The formation of such a hierarchical microstructure, including the lamellar one, can be achieved using the radial–shear rolling procedure. The results of studies of the strain mechanisms for various materials subjected to this processing have been widely published recently [25–29]. In all these papers, the importance of the lamellar microstructure formation in the conditions of multilevel structuring is noted for increasing strength and maintaining toughness of rolled bars.

A fairly complete review devoted to the formation of multiscale hierarchical microstructures, including hardening and providing ductility and plasticity by forming a laminate structure, is presented in [30]. According to the authors, elastic modulus and yield strength obey the “rule of mean”. However, elongation and fracture toughness cannot be estimated by this principle. It is known that (for laminate composites) elongation is determined by a number of parameters: layer thickness, interface structure, gradient structure, strain hardening exponent, strain rate, etc. They, in turn, depend on crack deflection, blunting, and bridging, as well as stress redistribution, local plane stress deformation, and interfacial delamination cracking. Thus, the ideology of the formation of the laminate composites can be efficiently used to develop the heterogeneous materials with increased strength and toughness. In this case, the laminate microstructure can be formed at various structural levels.

In general, the implementation of the laminate microstructure as an effective way to improve fracture toughness can be traced to the example of the Egyptian pyramids [31]. The laminate microstructure has shown particular efficiency in terms of increasing strength and toughness in the manufacture of bladed weapons such as Damascus steel, Japanese combat swords, etc. [32]. In this case, the composites are formed from single base metals, while the layer-by-layer joining of stronger (for example, ultra-high carbon (UHC)) and more ductile (for example, Fe-3Si) steels is considered to be more efficient [33].

A fairly detailed review of research results related to laminate metal composites (LMC) is presented in [34]. It has been shown that one of the effective methods for

the formation of such composites is accumulated roll bonding (ARB), and the highest both strength and toughness values are achieved by forming packets from materials with high and low strength. Numerous examples show that damage critical properties (fracture toughness, fatigue, and impact behavior) are largely depend on local delamination at the layer interfaces. In addition, the key factors that determine the mechanisms of distortion and failure are the parameters of the laminate formation, its 'architecture' and the microstructure of the layers and their interfaces.

It has been shown in [35] that post-processing heat treatment of a twelve-layer "UHC steel/mild steel" composite significantly improves its impact toughness by enhancing strength at the layer interfaces during tests at a temperature of 79 °C.

One of the most effective methods of increasing the strength properties of conventional steels (for example, the X52 grade HSLA ones) is to form certain microstructures, phase compositions, and internal stresses in the bulk metal [36]. Varying these parameters at the processing stage enables to obtain the required hardness, ductility, and crack resistance. Most often, processing includes heat treatment to form the optimal phase composition and plastic strain for additional hardening. Low-alloyed low-carbon steels of the pearlite class are best suited to such procedures [37,38]. In this regard, development of complex thermo-mechanical processing technologies is relevant to obtain materials that combine both high strength and sufficient ductility. To control and determine the optimal processing modes, the high-temperature helical rolling (HR) is widely applied, which can strain and heat treat the studied materials to a predetermined degree [39].

The aim of this research has been to investigate the HR parameters (a number of passes) for the microstructural modification and the low-temperature impact toughness improvement of the 09Mn2Si HSLA steel. In order to achieve this purpose, work spent to crack initiation and propagation has been analyzed and compared with patterns of fracture surfaces. The second section describes the details of the experimental studies, including the HR modes. In the third one, the results of studies of the microstructure (optical and SEM micrographs) are shown and impact toughness values are presented in the temperature range from +20 to −70 °C. Also, the failure mechanisms in individual regions on the fracture surfaces are discussed. In the fourth section, the authors propose a methodology for computer simulation of the process within the framework of the excitable cellular automata method and its integration with the kinetic theory of fracture (S.N. Zhurkov). A theoretical analysis of the effect of grain shapes and orientations on the strain response patterns of a certain meso-volume simulating the material after the HR processing is carried out in the fifth section. The results presented in Section 5 are not a complete analogy of the full-scale experiment, but are aimed at identifying the role of the processes at the internal interfaces on the propagation of energy flows. Finally, the authors draw conclusions based on the obtained results. It should be noted that the effect of such processing on the resistance to fatigue fracture for this steel has been considered in the previous paper of the authors [40].

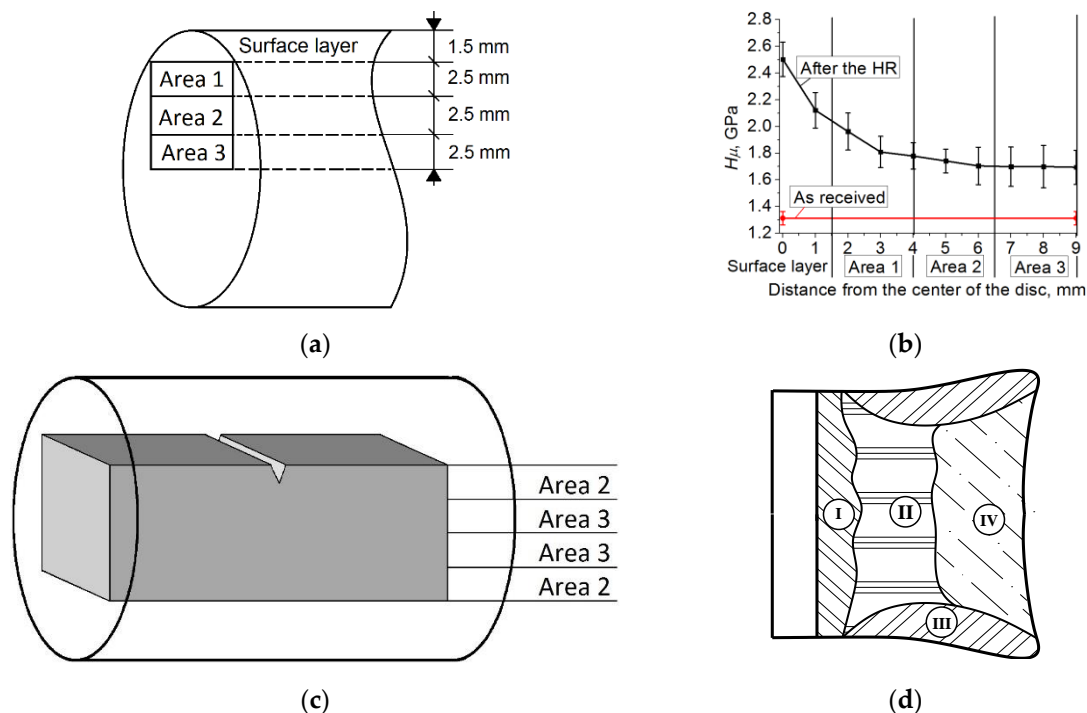
## 2. Materials and Methods

The HR processing of the 09Mn2Si HSLA steel (chemical composition:  $C \leq 0.12$ ;  $Si = 0.5\text{--}0.8$ ;  $Mn = 1.3\text{--}1.7$ ;  $Ni \leq 0.3$ ;  $S \leq 0.04$ ;  $p \leq 0.035$ ;  $Cr \leq 0.3$ ;  $N \leq 0.008$ ;  $Cu \leq 0.3$ ;  $As \leq 0.08$ ;  $Fe = 96\text{--}97$ ) was carried out using an RSP "14-40" three-roll mini-mill for screw rolling in the temperature range of 850–500 °C with a decrease in the rolling temperature at each subsequent pass (i.e., thermal schedule of the HR were are as follows: first pass: 850 °C; second pass: 750 °C; third pass: 650 °C; fourth pass: 580 °C; fifth pass: 500 °C.). After each rolling pass, water quenching was performed. Both three- and five-pass HR processing procedures were completed. The obtained results after various numbers of passes were compared to assess the dynamics of changes in the mechanical properties and determine the optimal mode of thermo-mechanical processing. The total degree of true logarithmic strain/accumulated strain of the steel after every pass of helical rolling was



equal to: (first) 0.09/8.8%; (second) 0.23/21.9%; (third) 0.45/41.5%; (fourth) 0.63/57.9%; (fifth) 0.80/73.4%.

Due to the metal inhomogeneity along the bar cross-section after the HR processing, the grain microstructure was evaluated within four regions (Figure 1a): the surface layer; area 1 (near the bar surface); area 2 (transition), and area 3 (the bar core). This differentiation is conditional, since the change in the microhardness values was smooth over the cross-section (Figure 1b). A 'PMT-3' microhardness tester (LOMO, St. Petersburg, Russia) was used to assess microhardness levels.



**Figure 1.** Scheme of the bar volume partitioning into areas (a), microhardness curve over the bar cross-section (b), the scheme of the specimen preparation for the Charpy impact tests (c), and the designation of the fracture areas in them (d): I—crack initiation area; II—crack propagation area; III—shear lips; IV—subcritical crack propagation area (rupture).

Specimens for the Charpy impact tests were cut from the bars by electrical discharge machining (Figure 1c). Impact toughness was determined using an 'Instron 450MPX' automated pendulum and the Charpy specimens of  $55 \times 10 \times 10 \text{ mm}^3$  (actually contained the metal from areas 2 and 3) with a V-notch 2 mm deep (according to ASTM E 23:2007). The tests were carried out at temperatures ( $T$ ) of +20, 0, −20, −40, and −70 °C. During 10 min immediately before testing, the specimens were cooled in a silicone oil environment using a 'Lauda rp 870' installation (chiller). Duration between the specimen removal from the cooling chamber and the test beginning was within 5 s. The lowest test temperature threshold was determined by the environmental operating temperatures in the Arctic region for the tested HSLA steel, as well as the possible temperature range of the 'Lauda rp 870' installation. The fracture surfaces of the specimens were observed using a 'LEO EVO 50' scanning electron microscope (Zeiss, Oberkochen, Germany) at the 'NANOTECH' Center for Collective Use of the ISPMS SB RAS.

### 3. Results

#### 3.1. Effect of the HR Processing on the Microstructure and Microhardness

Figure 1b shows the microhardness ( $H_\mu$ ) values on the bar cross-section after the three-pass HR. The test was performed to a depth of 9 mm, while the bar diameter was 18 mm after the HR processing. Compared to the as-received steel, strengthening had occurred throughout the entire volume of the bar with a gradual  $H_\mu$  decrease in the core.

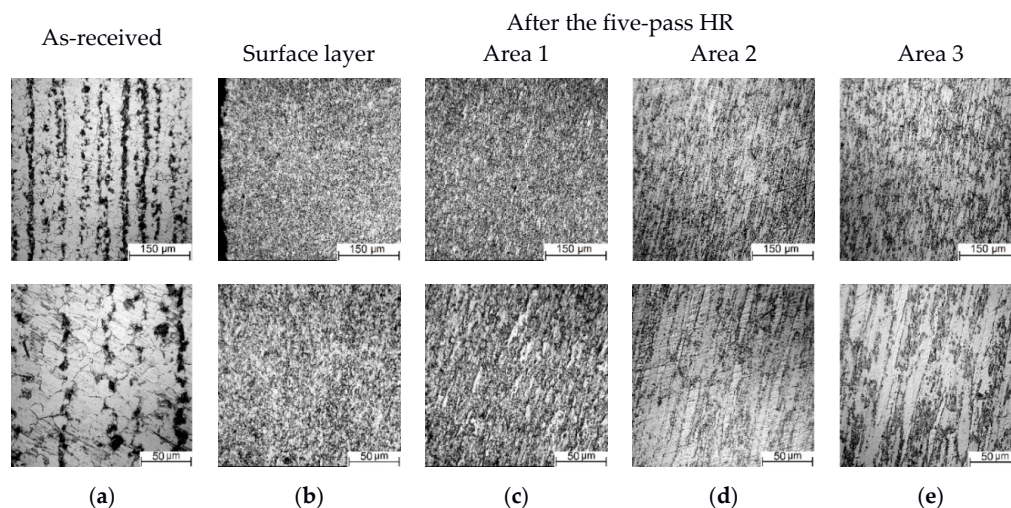
The greatest increase in microhardness was observed in the surface layer up to 3 mm thick. Table 1 presents the microhardness values in area 2 according to Figure 1a for the as-received steel specimen and after both three- and five-pass HR processing. The lowest  $H_{\mu}$  values were typical for the as-received steel, while the increase in the number of the HR passes resulted in a gradual  $H_{\mu}$  enhancing up to 20%.

**Table 1.** Microhardness values in area 2 according to Figure 1a.

Steel Conditions	Microhardness HV
As-received	$97 \pm 7$
After the three-pass HR	$116 \pm 3$
After the five-pass HR	$123 \pm 7$

The as-received steel was characterized by the ferrite-pearlite microstructure with a characteristic equiaxial grain size of  $21 \pm 2 \mu\text{m}$  (Figure 2a). Some striping was observed due to the use of rolling procedures for its manufacturing. In the surface layer after the five-pass HR, a fine microstructure was presumably composed of ferrite and fractured pearlite plates (Figure 2b). At a depth of more than 1 mm in the longitudinal cross-section, individual elongated ferrite grains were distinguished (Figure 2c; area 1), preserved after strains and slight deviated from the main axis towards the bar core at an angle of  $\sim 20^\circ$ . This orientation was due to a significant reduction in the bar diameter as a result of the HR processing, as well as the inclination of the rolls relative to the workpiece axis.

An increase in both the size and the number of elongated ferrite grains was observed towards the bar core. At a distance of more than 4 mm from the surface, the microstructure consisted entirely of elongated grains that formed the lamellar steel microstructure (Figure 2d, area 2). It was similar in the bar core (Figure 2e, area 3).



**Figure 2.** Optical images of the 09Mn2Si steel microstructure in the longitudinal cross-section (explanations for a–e are given in the text above).

The most intense elongation of grains had occurred along the bar axis in the longitudinal cross-section, while only a decrease in their sizes, comparing with the as-received steel, was observed in the transverse one (Table 2). In areas 2 and 3, the lamellar microstructure consisted of thin elongated ferritic grains oriented predominantly along the bar axis. In the surface layer at a depth of  $\sim 1$  mm, a finely dispersed microstructure without a pronounced texture had been formed, where the steel had been subjected to the maximum strains during the HR processing. Any clear boundaries between these layers were not observed. Since more pronounced microstructural changes were found after the five-pass HR, quantitative analysis of the grain microstructure was carried out only for these specimens.

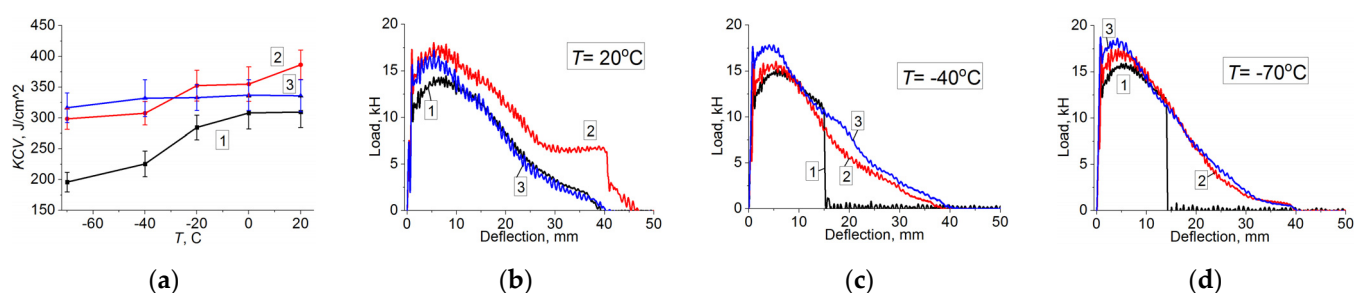
**Table 2.** Grain sizes in the 09Mn2Si steel after the HR processing.

Distance from the Bar Surface, mm	Grain Sizes, $\mu\text{m}$		
	As-Received	After the Five-Pass HR	
		Transverse Cross-Section	Longitudinal Cross-Section
Surface layer		$1.8 \pm 0.3$	$2.7 \pm 0.3$
1		$2 \pm 0.3$	$8.5 \pm 1.3$
2	$21 \pm 2$	$2.5 \pm 0.3$	$17.3 \pm 2.4$
3		$3.3 \pm 0.4$	$21.6 \pm 3.8$
4		$6.3 \pm 0.6$	$32.7 \pm 5.7$
5		$6.7 \pm 0.8$	$34.5 \pm 4.1$

### 3.2. Impact Toughness

The specimens for the Charpy impact test were cut from the bar core (actually from areas 2 and 3, according to Figure 1a,c). Due to the smooth nature of the decrease in the microhardness levels over the cross-section, the specimens possessed the hardened surface layers and the ‘softer’ core. The hardest layers (surface one and area 1), the presence of which during the Charpy impact tests could be accompanied by brittle fracture, were removed at the specimen fabrication stage (Figure 1c).

Both three- and five-pass HR processing of the 09Mn2Si steel caused its increased impact toughness over the entire test temperature range (Figure 3a, Table 3).



**Figure 3.** Impact toughness vs. test temperature curves (a); impact deformation curves for the Charpy specimens in the “load-deflection” coordinates (b–d); 1—as-received; 2—after the three-pass HR; 3—after the five-pass HR.

**Table 3.** Impact toughness (KCV) of the 09Mn2Si steel.

Test Temperature, °C	KCV, J/cm <sup>2</sup>		
	As-Received	After the Three-Pass HR	After the Five-Pass HR
−70	$196 \pm 16$	$299 \pm 17$	$317 \pm 24$
−40	$226 \pm 21$	$308 \pm 19$	$332 \pm 30$
−20	$285 \pm 20$	$353 \pm 25$	$333 \pm 21$
0	$308 \pm 26$	$355 \pm 28$	$337 \pm 25$
20	$310 \pm 25$	$387 \pm 24$	$336 \pm 26$

Unlike the as-received steel specimens, there was no critical decrease in toughness (KCV) of the specimens after the HR as the test temperature decreased. At the temperature of +20 °C, its higher values were observed for the ones after the three-pass HR processing. On the other hand, the difference in the KCV values decreased with reducing the test temperature for the steel after the HR processing using both modes. Thus, the three-pass HR proved to be a transitional one. In this case, both enhancing the impact toughness values (in comparison with the as-received steel specimens) and its reducing with the decrease in the test temperature were observed.

The most important obtained result was the stable retention of the impact toughness level throughout the studied temperature range for the steel specimens after the five-pass

HR processing. The absence of this effect for the ones after the three-pass HR could be justified by the following reasons: (i) the homogeneous ‘lamellar’ microstructure had not been formed due to the fewer number of passes; (ii) the softer metal was due to a lower level of accumulated strains; (iii) quenching in water after the third pass had been carried out from the higher temperature, which had caused the formation of the more non-equilibrium microstructure with lower toughness.

### 3.3. Analysis of the Loading Diagrams for the Charpy Impact Test Specimens

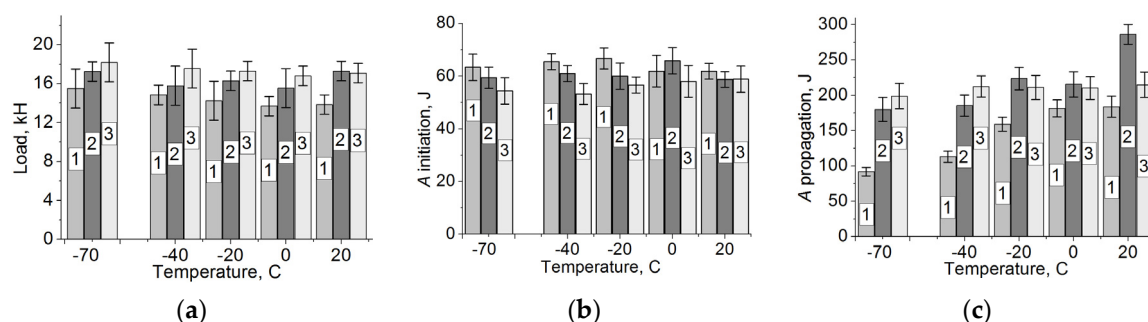
According to the data presented in the ‘load–deflection’ graphs (Figure 3b–d), the curves possessed noticeable variations in their shapes even at the temperature of +20 °C. In this case, all specimens fractured in the ductile manner, which confirmed fairly smooth changes in the diagrams without abrupt transitions and bifurcation points (Figure 3b; curve 1). This testified to the effective relaxation of stresses at both macro-crack initiation and propagation stages.

The shape of the diagram for the steel after the three-pass HR at the pre-fracture stage should be highlighted (Figure 3b, curve 2, constant  $p$  value in the deflection range of 30–40 mm). According to the authors, such curve indicated a significant margin of the material ductility even at the final fracture stage. The curvature of the diagram could be caused by a sharp turn (branching) of the fracture front. On the other hand, the steel specimens after the five-pass HR showed diagrams similar to those for the as-received ones, but with a higher energy intensity to fracture, as evidenced by the greater maximum load ( $P_{max}$ ) values (Figure 4a), as well as a larger area under the curves, reflected work spent to the specimen failure.

At the temperatures of −40 and −70 °C, impact fracture of the as-received steel specimens was accompanied by ductile deformation at both the crack initiation and (partially) propagation stages. However, when the critical crack length was reached, the character of its propagation changed from ductile to brittle. This fact was justified by load breakdown in the diagram (Figure 3c,d; curve 1).

On the other hand, the diagrams of the specimens after both three- and five-pass HR showed maintaining of the ductile nature of fracture, which indicated the preservation of relaxation processes (at the micro-scale level) and was accompanied by the noticeable increase in fracture toughness (Table 3).

One of the parameters reflected the effect of the HR processing on both strength and dynamic crack resistance of the steel was the maximum load ( $P_{max}$ ) that the specimens withstood upon the Charpy impact tests (Figure 4a). This parameter corresponded to the moment of transition from the macro-discontinuity (crack) initiation to its propagation. Its highest values for the specimens after the five-pass HR, primarily at the low test temperatures, showed the important role of the lamellar (fine-grained) microstructure formation in resistance to impact loading.

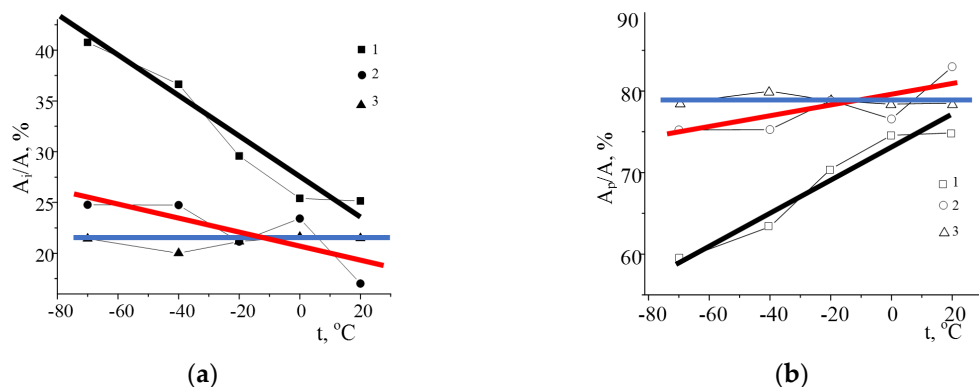


**Figure 4.** Maximum load to the striker during the Charpy impact tests (a); work spent to crack initiation (b) and propagation (c); as-received; 2—after the three-pass HR; 3—after the five-pass HR.

### 3.4. Calculation of Work Spent to Crack Initiation and Propagation under Impact Loading

Since impact toughness was determined by the energy levels for both crack initiation and propagation, their values were analyzed separately. It was found that the crack initiation energies  $A_i$  for the as-received steel specimens were slightly higher at almost all test temperatures than those for the steel after the HR processing. Its values changed in the range of 60–70 J (Figure 4b). The increase in the number of the HR passes was accompanied by a slight decrease in the  $A_i$  values, which, nevertheless, did not reduce below 55 J. On the other hand, the crack propagation energies  $A_p$  were always higher for the steel specimens the HR processing. After the three-pass HR,  $A_p$  decreased from ~270 down to ~180 J as the test temperature reduced, while it was about 200 J after the five-pass HR processing (Figure 4c). The greatest drop in the amount of work spent to crack propagation was observed for the as-received steel specimens, which was primarily associated with their brittle fracture at the subzero test temperatures (Figure 3c,d).

In order to estimate the ratio of energy spent to crack initiation and propagation to the total amount of energy spent to the specimen fracture, both  $A_i/A$  and  $A_p/A$  versus temperature graphs were plotted. Analysis of the data shown in Figure 5a enabled to conclude that the  $A_i/A$  ratio changed from 25.2% at the temperature of +20 °C up to 40.8% at −70 °C. Such a significant specific part of the energy intensity spent to crack initiation for the as-received steel specimen at the temperature of −70 °C indirectly indicated its embrittlement. For the ones after the three-pass HR, this was much less pronounced: there was an increase in the  $A_i/A$  ratio from 17.0% at +20 °C up to 24.8% at −70 °C. With increasing the number of passes to five, the  $A_i/A$  ratio values practically did not change, remaining at a level of ~21.5% both at +20 and −70 °C. Thus, the increase in the number of the HR passes resulted in the exclusion of the sensitivity of the 09Mn2Si steel to the embrittlement effect of temperature (from the point of view of crack initiation in the investigated temperature range).



**Figure 5.** Energy intensity to failure at the main crack initiation (a) and propagation (b) stages: 1—as-received; 2—after the three-pass HR; 3—after the five-pass HR.

A matching pattern, only in the ‘mirror’ reflection, was characteristic of the  $A_p/A$  parameter (Figure 5b). This enabled to assume the similarity of the fracture micro-mechanisms and morphology of the corresponding fractographic surfaces on the steel specimens after the HR processing, but their significant distinctions for the as-received ones. The HR processing had changed the grain microstructure hierarchy, their shapes and sizes, as well as affected the mechanical properties of the steel, including resistance to impact fracture. The decrease in the grain sizes caused the significant increase in the microstructure heterogeneity. This resulted in two main effects:

- reducing energy consumption to crack initiation, which was expected due to increasing its structural heterogeneity and enhancing the specific length of the grain boundaries (local heterogeneity zones) in micro-volumes of the steel specimens after the HR [41], in comparison with the as-received ones;



- increasing the energy intensity to crack propagation, which was associated with high energy consumption in such a “micro-composite material” with the lamellar mesostructure [42,43].

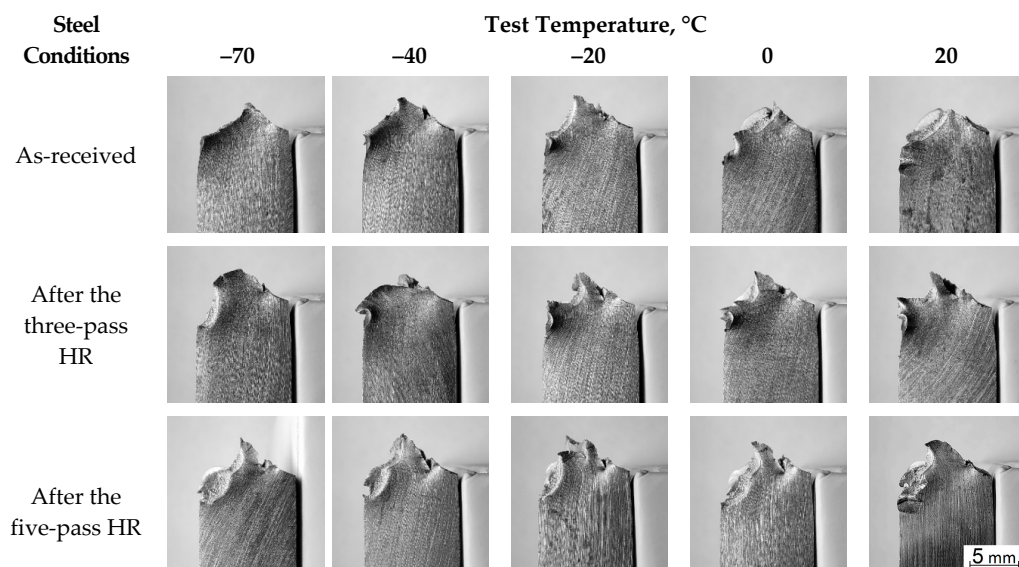
At the same time, the lamellar microstructure, formed during the HR processing, affected similar both  $A_i/A$  and  $A_p/A$  ratios. It should be noted the additional ductility depletion for the 09Mn2Si HSLA steel after the five-pass HR, which was especially noticeable from the microhardness test results (Table 1). These patterns indicated that this type of steels could possess increased sensitivity to aging, hydrogen absorption, and other degradation factors during long-term operation [44].

Thus, the main reason for the decrease in the steel toughness at the low test temperatures was a gradual decrease in resistance to crack propagation. In this regard, the microstructural modification by the HR processing, despite the increase in microhardness, contributed both to the fracture resistance maintenance almost at the initial level during the crack initiation stage and to its significant enhance upon the macro-discontinuity propagation. From the point of view of increasing the specific area of the grain boundaries along the crack propagation path at the second stage, this result was natural for the ‘lamellar’ microstructure of the steel after the HR processing. However, the fact of impact fracture toughness improving, even at the low test temperatures, required a deeper understanding.

### 3.5. Shapes of Shear Lips

The 09Mn2Si HSLA steel was characterized by a sufficient margin of ductility (toughness), which was noticeable by the angle of inclination of the shear lips relative to the specimen axis (Figure 1d), which was about 130° (Figure 6).

It is known from previous papers [45] that shear lips had formed near the macro-band of strain localization, both in the direction of the maximum shear stresses and perpendicular to the tension axis under the greatest normal ones. Shapes and dimensions of the shear lips characterized the effect of lowering the test temperature on the impact resistance of the material. They were bent at the steel specimens after the three-pass HR, which indicated an increase in the energy intensity to failure. At the same time, the rounded shape of the shear lips, even at the test temperatures of  $-40$  and  $-70$  °C, corresponded to a significant margin of fracture toughness. After the five-pass HR, discontinuities were at the shear lips due to the formation of strain localization zones (however, they were also observed on the steel specimens after the three-pass HR). This contributed to an additional increase in the energy intensity to failure.



**Figure 6.** Effect of the Charpy impact test temperature on the 09Mn2Si steel specimens (both as-received and after the HR processing).

### 3.6. Fractographic Analysis

SEM micrographs of the fracture surfaces are shown in Figures 7 and 8. At the temperature of +20 °C, all steel specimens after the HR processing fractured in a rather ductile manner. This fact was evidenced by a noticeable change in their shapes over the cross-sections (Figure 7a–c). The fracture surfaces were characterized by developed shapes with strain localization zones, as well as local delamination. Apparently, depth of these layers largely determined the energy intensity to failure, and as mentioned above, was reflected in the shape of the shear lips [46].

- At the micro-level, rounded pits surrounded by shear fracture spots had been formed on the as-received steel specimen in area 1 (crack initiation, Figure 1g)). In general, it had a ductile-brittle pattern (Figure 7a,g). The fracture surface was quite flat, formed according to the ductile-brittle mechanism. This was confirmed by the presence of multiple “chips”, which alternated with areas of ductile fracture of a dimple type. In area 2 (crack propagation, Figure 1m)), the surface morphology was very rough that indicated a clear heterogeneity in the strain process development at the main crack propagation (failure) stage (Figure 7s). It had been formed by large elongated dimples oriented across the crack propagation direction, and also included the large zone 4 (Figure 7s), corresponded to the specimen failure according to Figure 1g, which had a smoother surface with a gradual bend towards the point of impact loading.
- After the HR, the crack initiation mechanism was predominantly ductile-brittle, with a minor brittle component (Figure 7b,c,h,i). Dimples had deformed profiles oriented along the crack front direction that indicated an increase in the material ductility and its sufficiently high crack resistance. This was confirmed by data on the energy intensity to failure at this stage. Apparently, the formation of large dimples could not be realized due to local plastic strains. The proportion of areas with micro-dimples decreased on the fracture surfaces. At the same time, the fraction of the fracture surface characterized only by flattened cone pits and tunnel ones, as well as their dimensions, increased significantly. Crack propagation was accompanied by significant strains and delamination. At the same time, the fractographic relief of the areas between them looked rather brittle. This indicated the localization of the fracture processes within individual micro-zones (Figure 7k,l). Area 4 (Figure 7t,u) had a characteristic rough relief, without pronounced distinctions from a similar fracture of the as-received steel specimen.

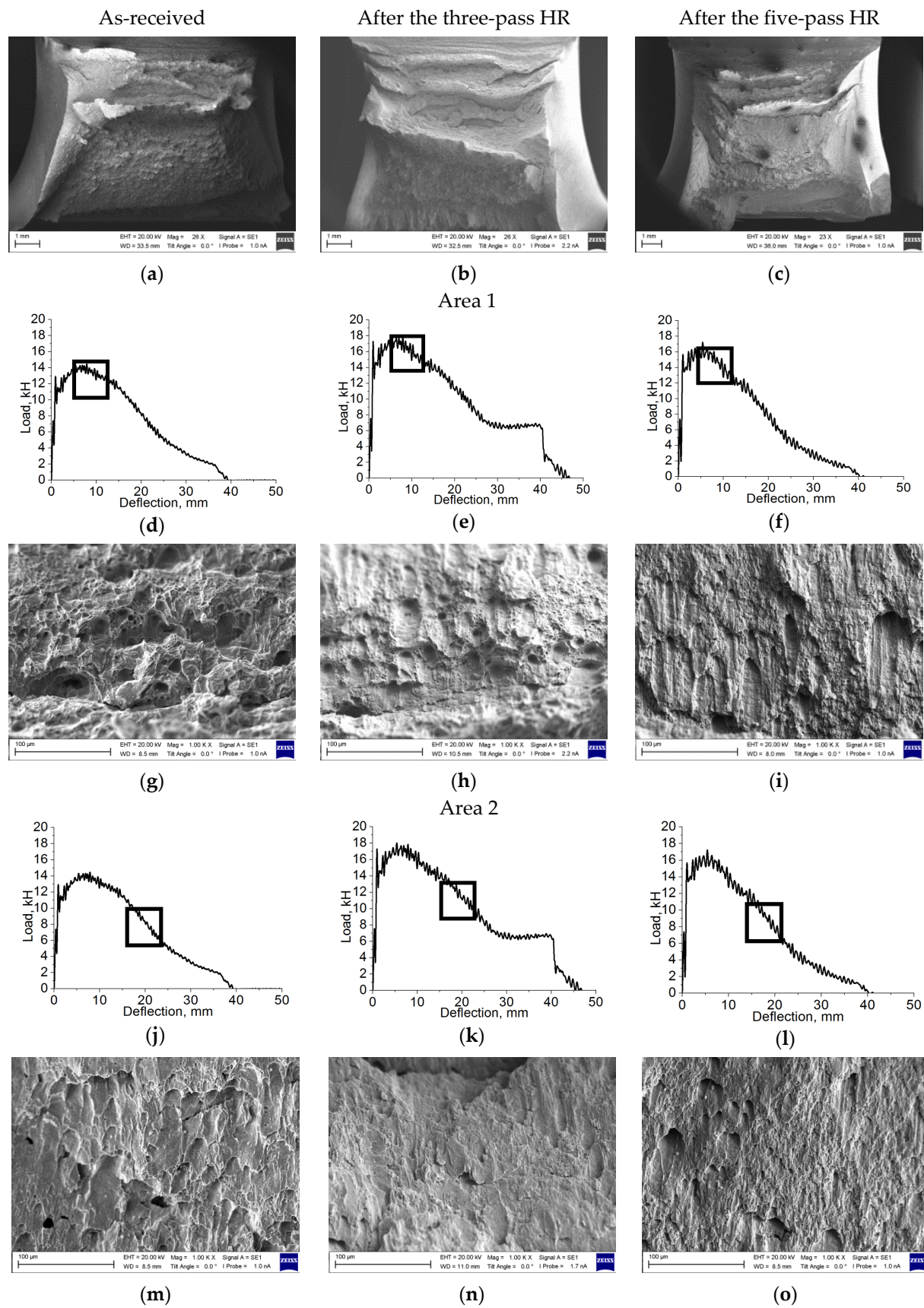
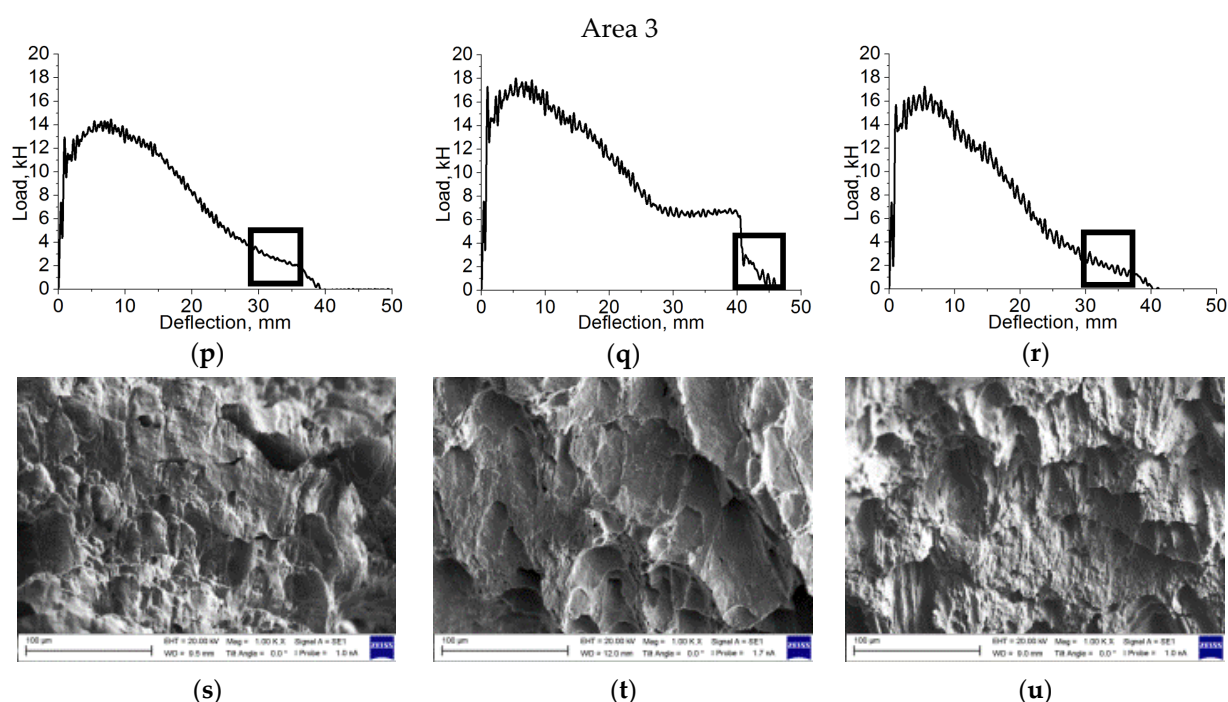


Figure 7. Cont.



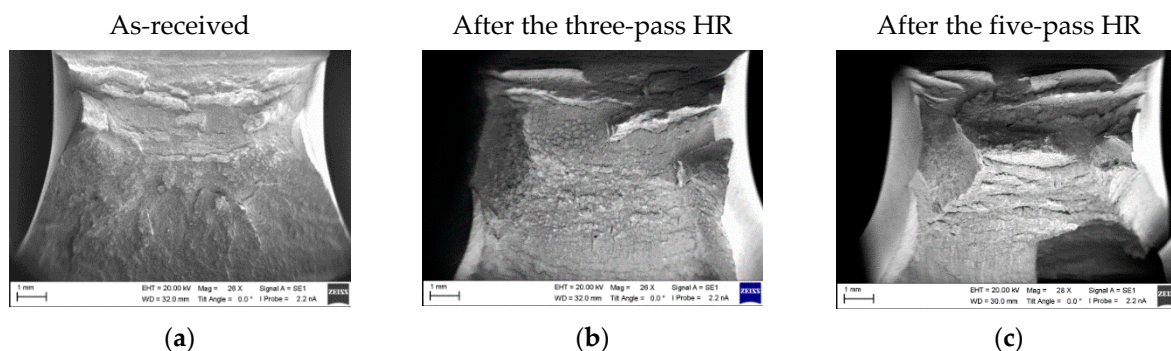


**Figure 7.** Fracture surfaces of the 09Mn2Si steel after the Charpy impact tests at the temperature of +20 °C. (explanations for **a–u** are given in the text above).

At the temperature of  $-70^{\circ}\text{C}$ , the fracture surface macro-irregularities were observed near the crack initiation area (Figure 8), which indicated the presence of local strains during specimen failure. The shapes of the specimen transverse cross-sections were deformed at the entire investigated temperature range, which reflected high energy intensity to failure.

At the micro-level, crack initiation and propagation areas had been formed due to shear strains for the as-received steel specimens (Figure 8g,m). The fracture surfaces had a smoothed 'structureless' appearance with several micro-chips. The fracture mechanism had been mixed with the predominant realization of micro-chips at the boundaries of deformed micro-zones. The fractographic relief reflected the ductile-brittle pattern of fracture. At the failure area (Figure 8s), where the sharp drop in load had been previously found in the load-deflection curves (Figure 3c,d), a brittle fracture surface was observed.

The specimens after the HR retained shear fracture pattern. Ductile tear dimples were absent (Figure 8b,c). The surfaces had an even more 'blurry' and 'structureless' character (Figure 8h,i,n,o). No significant variations were found in the fractographic surface pattern for the specimens after both three- and five-pass HR. This was also consistent with the data presented above on the study of the fracture macro-characteristics. The main difference from the as-received steel specimens was the ductile fracture type at all areas, including the failure one (Figure 8t,u).

Figure 8. *Cont.*

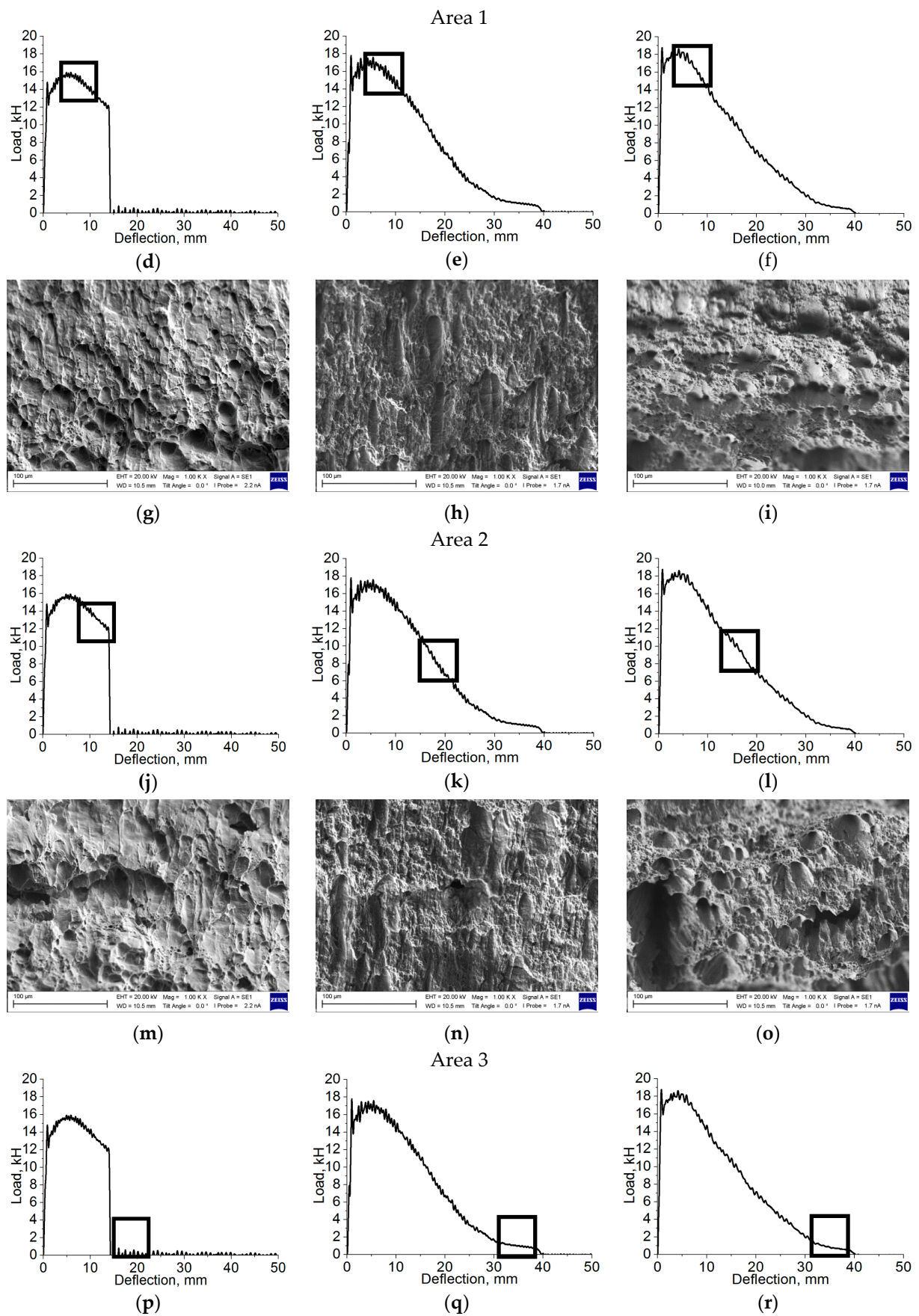
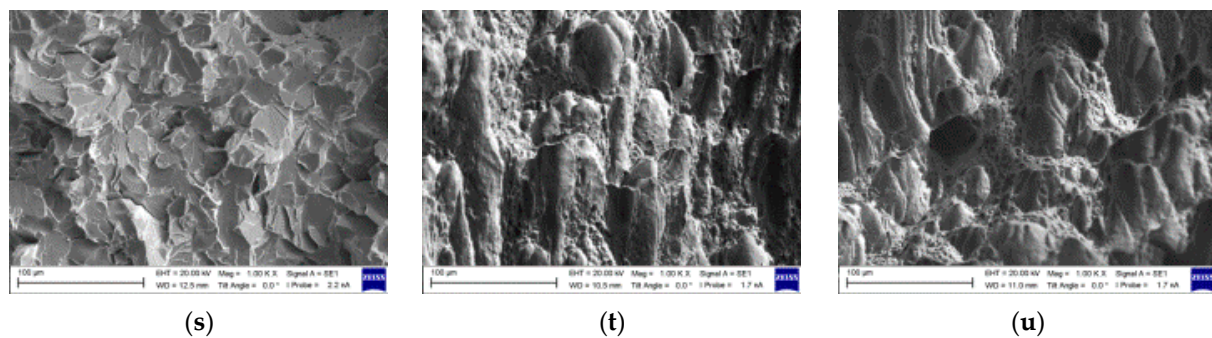


Figure 8. Cont.





**Figure 8.** Fracture surfaces of the 09Mn2Si steel after the Charpy impact tests at the temperature of  $-70\text{ }^{\circ}\text{C}$ . (explanations for a–u are given in the text above).

In this paper, it was shown that the increase in the number of the HR passes from three to five gave a small effect on increasing impact toughness of the 09Mn2Si HSLA steel, which was associated with the exhaustion of the material ductility margin. Analyzing the data in Figure 8, it was possible to note the similarity of the fracture mechanisms after the three- and five-pass HR processing. It could be assumed that a further increase in the number of passes would result in embrittlement of the steel and a decrease in its toughness [47].

#### 4. Justification of the Theoretical Approach to the Analysis of Deformation Processes under Impact Loading

The effects revealed in the experimental section required complex analysis, including the use of computer simulation methods. A multilevel approach for simulation of deformation and fracture processes seemed promising [48–50]. However, before solving the issue of combining mechanisms operating at different scale levels (solved, for example, within the framework of molecular dynamics, crystal plasticity, finite element modeling, etc.), it was necessary to analyze the contribution of the processes at each scale. In this regard, it was expedient to apply the excitable cellular automata method at the meso-level, which showed its effectiveness for analyzing the behavior of materials under intense loads, considering rotational strain modes and material fragmentation at the pre-fracture stage. In this paper, the authors presented how this method could be developed to consider this period in terms of micro-damage initiation.

In the previous paper [51], the authors—within the framework of the excitable cellular automata method [52–54]—implemented a model for the numerical description of the process of material fatigue under cyclic loading. This model used Mott’s theory [55], which postulated the dominance of kinetic (thermal fluctuation) processes when considering the dynamics of sliding of crystallite planes. According to this, the general expression for the shear activation energy  $E$ , taking into account the external stress  $\sigma$ , was written as

$$E = nF \pm \sigma n\omega a/2, \quad (1)$$

where  $a$  was an interatomic distance;  $n$  was the number of atoms in a cluster of the crystallographic plane, the area of which was  $n\omega$ ;  $F$  was the free energy required for shear due to cluster disordering. In this case, the expression for the free energy was written as

$$F = nL (1 - T/T_M), \quad (2)$$

where  $T_M$  was the melting point;  $L$  was latent heat of fusion per atom.

Mott’s theory described inter-crystalline shear as an oscillatory process directly related to thermal vibrations of the atomic lattice. Thus, for example, the frequencies of cluster shifts by distance  $a$  could be written as

$$\zeta_1 = \zeta_0 e^{-(nU \pm \sigma n\omega a/2)/kT} \quad (3)$$

where  $\zeta_0$  was the frequency of atomic vibrations.

Thus, internal stresses relaxed in a crystal at shear areas of the atomic planes. On the other hand, the mechanical energy of an external load enabled to dissipate due to thermal fluctuation processes associated with these shears.

In this paper, the authors, adhering to the line of dominance of the kinetic processes at the micro- and mesoscale levels, supplemented the excitable cellular automata model in the part of the thermo-fluctuation theory of strength, first proposed by S.N. Zhurkov in [56] and greatly developed later [57–60]. Considering the structure-dependent kinetic processes was also very important to establish the dependence of strength characteristics from temperature and loading rate upon material fracture. These tasks were relevant when solving, for example, such a challenge as low-temperature brittleness [61].

In the kinetic concept of S.N. Zhurkov, fracture was considered as a process that began immediately after application of an external load to a body and continuously developed in time. According to this theory, duration to failure  $\tau$  was determined by the following relation:

$$\tau = \tau_0 e^{(U_0 - \chi\sigma)/kT}, \quad (4)$$

where  $k$  was Boltzmann's constant;  $\sigma$  was tensile stress;  $T$  was temperature;  $\chi$ ,  $\tau_0$ , and  $U_0$  were material constants that had a clearly defined physical meaning. Thus,  $\tau_0$  was duration of natural vibrations of solid atoms (about  $10^{-13}$  s);  $U_0$  was the level of the energy barrier that determined the probability of rupture of the interatomic bond.

It could be concluded that  $\tau_0$  was the reciprocal of  $\zeta_0$  in expression (3) by S.N. Zhurkov, who described the frequency of an elementary act of shear strain in Mott's theory. Thus, frequency of natural vibrations of the crystal lattice was a fundamental quantity both in describing plastic strain and fracture of a solid.

The parameter  $\chi$  corresponded to a wide range of defect microstructures nucleated in a solid due to fracture under external loads and strongly depended on the material. It was found for the case of metallic single crystals that this coefficient was strictly inversely proportional to the density of dislocations generated under the external load. For pure polycrystalline metals, this value was determined by the density of dislocations at the grain boundaries.

Within the framework of the developed cellular automata approach, the material 'torque' (rotation) occurred under a local moment of force in an active element at a certain angle. The crystal lattice was distorted under this rotation, which was accompanied by the formation of single dislocations, and, then, multiple slip lines. Calculation of the material torsion angle enabled to quantify the induced change in density. For this purpose, the authors applied the approximation of the rod torsion model, considered in detail in [53,54]. Thus,  $\chi$ , being the reciprocal of the density of dislocations, could be calculated proceeding from the local moment of forces by the formula:

$$\chi = \frac{k_{lat}}{\Delta\rho_d} = \frac{k_{lat}G\pi r S_{CA}}{M}, \quad (5)$$

where  $k_{lat}$  was the proportionality coefficient associated with the type of crystal lattice;  $r$  was the radius of an active element;  $G$  was shear modulus;  $S_{CA}$  was the cross-sectional area of the cellular automaton active element.

The action of local moments of forces in the material was the reason for the transition of a cellular automaton active element from the initial 'elasticity' state to the 'pre-failure' one. The probability of this transition  $P_1$  was calculated for the element with index  $i$  as

$$P_1 = \sum_{j=1}^J \frac{|\vec{M}_i - \vec{M}_j|^2}{\left(|\vec{M}_i| + |\vec{M}_j|\right)^2}, \quad (6)$$

where  $j$  was the local index of a neighboring element on the first coordination sphere of the element with index  $i$ ;  $J$  was the number of such neighboring elements (in the case of FCC packing  $J = 12$ ).

Before the beginning of the rapid fracture stage, stationary accumulation of cracks occurred when the rate  $N$  was constant, and the number of cracks tended to a certain critical (threshold) value  $N^*$  (the number of cracks per unit volume). The critical value was determined from the concentration law [62]

$$N^{*-1/3}/L = K^*, \quad (7)$$

where  $L$  was a mean crack length;  $K^*$  was the dimensionless material parameter in the range from 2.1 to 6.0. It was easy to see that  $K^*$  had the meaning of the cubic root of the quantity, which was inversely proportional to the critical volume fraction of all cracks in the material.

Within the framework of the excitable cellular automata method, the last relation could be used to calculate the probability of the transition of an active element from the 'pre-fracture' state to the 'crack' one ( $P_2$ ). This was reflected in the laws of interaction with neighboring active elements. This probability was defined as the ratio of the density (amount per unit volume) of accumulated microcracks to the critical density:

$$P_2 = \frac{N}{N^*}. \quad (8)$$

Furthermore, in terms of a cellular-automaton evolutionary scheme, the probability was expressed as a sum over time steps of the algorithm as

$$P_2 = \dot{N}_0 K^3 \Delta t \sum_0^t e^{-(U_0 - \chi \sigma)/kT}, \quad (9)$$

The authors believed that the modification of the steel microstructure by the HR processing contributed to a longer stage of stationary damage accumulation, delaying the loss of the material's bearing capacity. The degree of structural 'imperfection' of the material was effectively considered in the theory of S.N. Zhurkov by means of the  $\chi$  parameter, which was determined by the density of dislocations at the crystallite level and by the density and energy state of the internal interfaces at the macro-scale one. Thus, within the framework of the described method, it was possible to consider both the temporal features of the damage accumulation process (according to the theory of S.N. Zhurkov) and the large-scale one (analyzed using the physical mesomechanic postulates), including the mesoscopic consideration of the deformation processes in Mott's theory. The method explicitly took into account the energy dissipation during the nucleation of defect structures in accordance with the relation [54]

$$\Delta E_d = \frac{k_{\text{diss}} G |\Delta \vec{\gamma}| 2\pi r_c^3}{2}, \quad (10)$$

where  $\Delta \vec{\gamma}$  was an increment of the vector torsion angle of the material in the active element;  $G$  was shear modulus;  $k_{\text{diss}}$  was the experimentally measured dissipation coefficient;  $r_c$  was the radius of the active element.

The authors suggested the development of the presented approach of cellular automata modeling in determining the critical tensile strength in order to use its values for simulation processes at higher scale levels.

## 5. Numerical Simulation Results

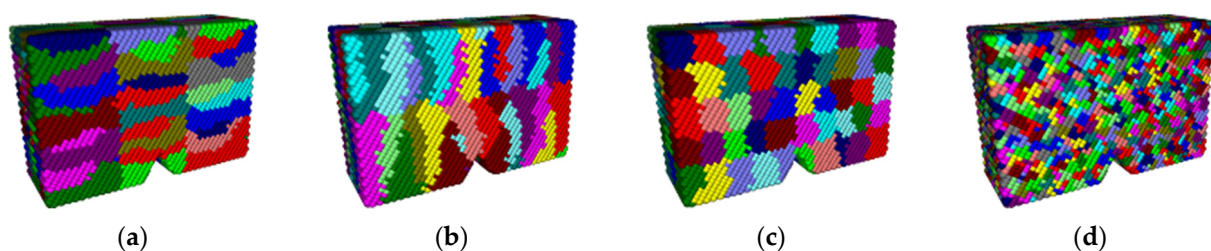
As noted in the experimental section, resistance to impact fracture was significantly affected by the processes at the macro-discontinuity propagation stage. However, the

microstructure formed upon the HR processing was decisive in terms of energy spent to crack initiation. The HR could have had a dual effect. On the one hand, an increase in yield stress restrained the transition of the material to plastic flow and, accordingly, the moment of crack initiation. A competing process was strain hardening (revealed, for example, by the microhardness tests) that reduced toughness. Together with the retention of strain dislocation mechanisms in the steel after the HR processing, it was possible to ensure the KCV values at the as-received steel level, while there was no significant decrease in impact toughness with the decrease in the test temperature compared to the as-received steel.

Thus, it was necessary to study how the modification of the microstructure by the HR processing affected propagation of energy fluxes in the steel specimens under impact loading. This section was devoted to solving the issue. The authors noted that they did not try to achieve a complete analogy with the experimental conditions (as was typically done for the finite element method). Nevertheless, a comparison of the obtained results with the above-described experimental data enabled to both develop the simulation method and outline ways to further improve the properties of structural materials by their microstructure modification via the HR processing.

Implementing the modified excitable cellular automata method, two computer simulation series were carried out to study in detail impact loading of the steel specimens with different microstructures characterized by grain sizes of  $30 \times 6 \times 6$  (grains were elongated along the  $x$ -axis),  $6 \times 6 \times 30$  (grains were elongated along the  $z$ -axis),  $9 \times 10 \times 12$  and  $3 \times 3 \times 3$   $\mu\text{m}$  (Figure 9). Each specimen had dimensions of  $90 \times 30 \times 60$   $\mu\text{m}$  and a symmetrical V-shaped notch 15  $\mu\text{m}$  wide and 13  $\mu\text{m}$  deep in the middle of the lower face. The size of a cellular automaton element was 1.5  $\mu\text{m}$ , the time step was 1 ns.

In the first simulation series, the temperature of the specimens was  $+20$   $^{\circ}\text{C}$ , but it was  $-70$   $^{\circ}\text{C}$  in the second ones. During each simulation with a duration of 100  $\mu\text{s}$ , an elliptical area with semi-axes of 30 and 10  $\mu\text{m}$  and a center in the middle of the upper face was under impact loading, and the smooth distribution of compression strain was set so that its maximum value ( $100 \text{ s}^{-1}$ ) was reached in the ellipse center, and the minimum (zero) was on its boundary. The smooth load distribution on the lower face simulated the support reaction: the value of the compression strain rate was maximal ( $10 \text{ s}^{-1}$ ) at the face edges, but it was minimal ( $2 \text{ s}^{-1}$ ) at the face center. The rest of the specimen faces were not loaded and did not exchange energy with environment.



**Figure 9.** Microstructures of the steel specimens with grain sizes:  $30 \times 6 \times 6$   $\mu\text{m}$  (a),  $6 \times 6 \times 30$   $\mu\text{m}$  (b),  $9 \times 10 \times 12$   $\mu\text{m}$  (c),  $3 \times 3 \times 3$   $\mu\text{m}$  (d).

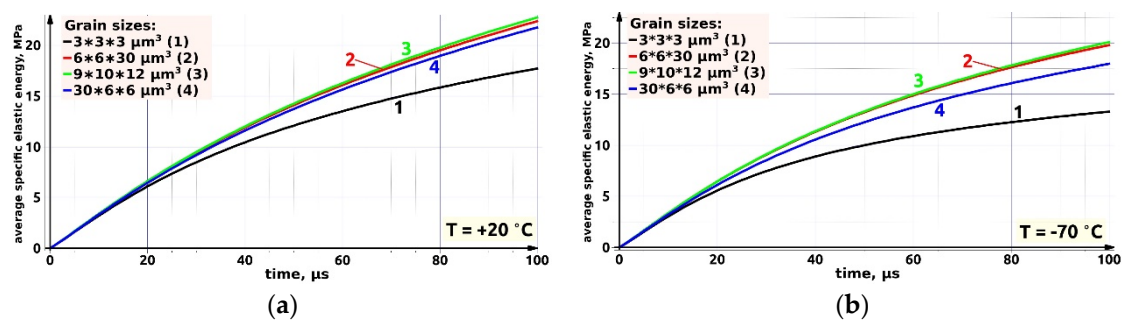
Figures 10–12 show plots of time dependences of specific elastic energy, specific dissipated torsion energy and their proportions in the total energy, as well as mean moduli of specific moment of force for each of the four considered specimens. The expressions, according to which the calculations of these parameters were carried out, had been given in [52].

The behavior of the graphs in Figure 10 enabled to conclude that the rate of accumulation of elastic energy gradually decreased in all specimens, and the minimum values of elastic energy were reached in the specimen with fine grains (3  $\mu\text{m}$  in size). Of the other three specimens, the highest value of the accumulated specific elastic energy was shown



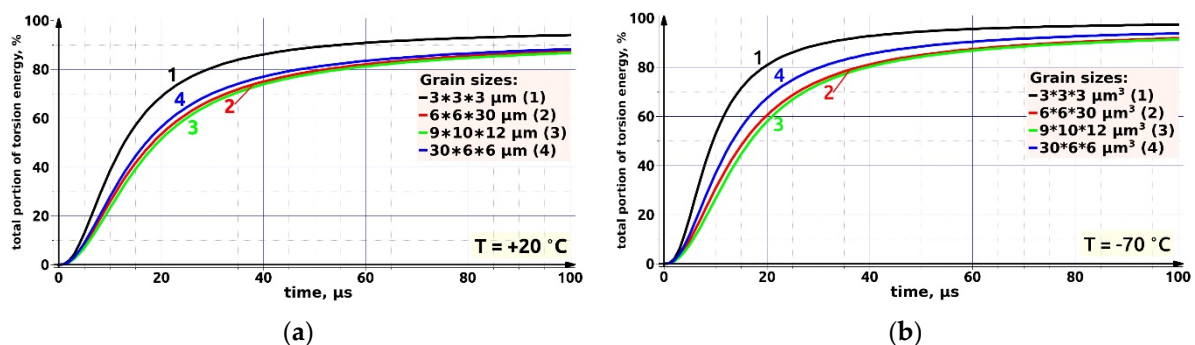
by the specimen with almost equiaxed grains of  $\sim 10\ \mu\text{m}$ , and the smallest one possessed the specimen with horizontally oriented (elongated) grains.

The curves shown in Figure 10b, had very significant distinctions in the specimen energy capacity with lowering temperature. For instance, the accumulated energy in the finest-grained specimen at the temperature of  $-70\ ^\circ\text{C}$  was 30% lower than that for the same microstructure at  $+20\ ^\circ\text{C}$ . It could also be noted that the specimen with the lamellar horizontal microstructure ( $30 \times 6 \times 6\ \mu\text{m}^3$ ) showed a noticeable sensitivity to the decrease in temperature compared with the ‘coarse-grained’ one. In this case, low temperature, which changed the energy state of the grain boundaries, caused a higher barrier to the elastic energy front propagation.



**Figure 10.** Graphs of time dependences of specific elastic energy for the specimens with grain sizes  $30 \times 6 \times 6\ \mu\text{m}$  (blue, 4),  $6 \times 6 \times 30\ \mu\text{m}$  (red, 2),  $9 \times 10 \times 12\ \mu\text{m}$  (green, 3) and  $3 \times 3 \times 3\ \mu\text{m}$  (black, 1) at the temperatures of  $+20\ ^\circ\text{C}$  (a) and  $-70\ ^\circ\text{C}$  (b).

The value of the specific torsional energy characterized the ability of the material to nucleate new defect structures according to expression (10). Figure 11 shows a gradual increase in the growth rate of the specific torsional energy for all specimens. At the same time, the maximum values were for the fine-grained material during the entire simulation. From three other specimens, they were for the material with grains oriented across the load direction. The minimum values of the torsional energy were accumulated in the specimen with large equiaxed grains (highly likely, due to the small extent of the boundaries themselves). At the end of the loading stage, the values of the torsional energy fractions stabilized in the range of 85–95%. The difference was noticeable only for the fine-grain specimens.



**Figure 11.** Graphs of time dependences of the proportion of the specific dissipated torsional energy in the total energy for the specimens with grain sizes  $30 \times 6 \times 6\ \mu\text{m}$  (blue, 4),  $6 \times 6 \times 30\ \mu\text{m}$  (red, 2),  $9 \times 10 \times 12\ \mu\text{m}$  (green, 3) and  $3 \times 3 \times 3\ \mu\text{m}$  (black, 1) at the temperatures of  $+20\ ^\circ\text{C}$  (a) and  $-70\ ^\circ\text{C}$  (b).

Analysis of the behavior of the plots in Figures 10 and 11 enabled to conclude that the maximum dissipation of elastic energy occurred in the fine-grained microstructure. Analysis of the data on the specific torsional energy indicated that plastic strain developed according to the ‘shear plus rotation’ mechanism. In this case, the relaxation of excess stresses was in the grain bodies due to interplanar slip (described within the framework



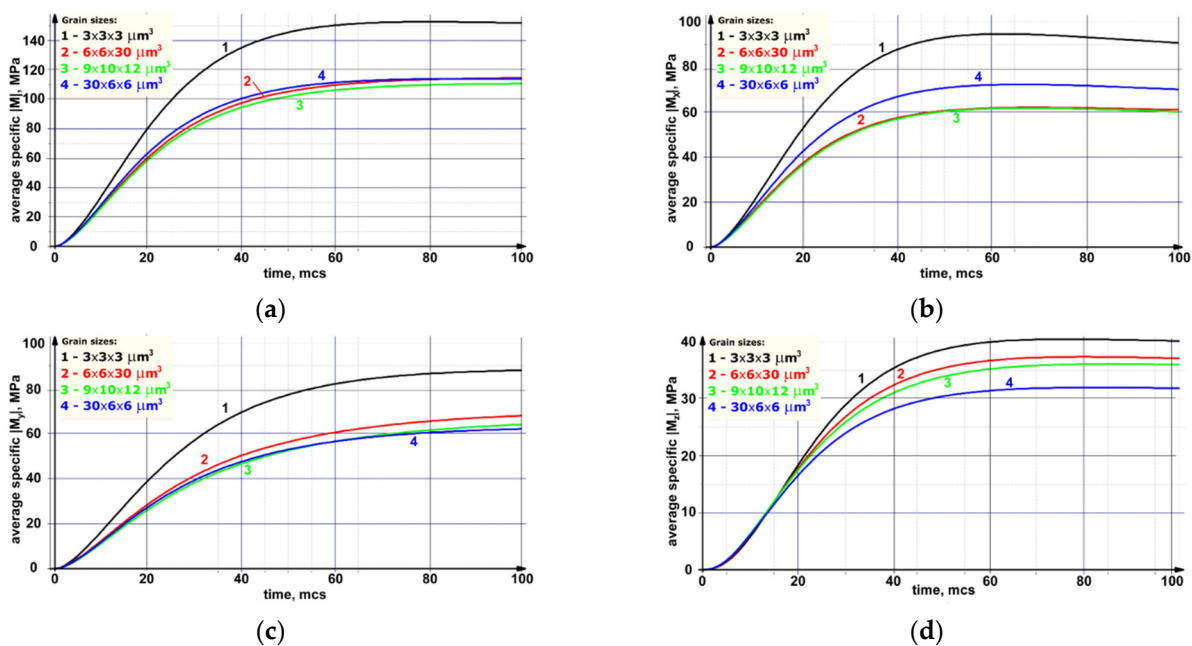
of Mott's theory), while the rotational mechanism had to prevail at the grain boundaries. It should be noted that grain boundaries were not only the barrier preventing the elastic energy transfer and absorbing part of it, but they also significantly changed the stress–strain state patterns in the grain bodies.

Of the three specimens with the same mean volume, but having different grain orientations, the highest dissipation of elastic energy was shown by the one with the preferred orientation of the boundaries across the load direction. This was due to the fact that the strain front was forced to overcome the area of the grain boundaries larger than in other specimens during the propagation process. On the other hand, the 'vertical' grain orientation also increased the degree of energy dissipation compared to the equiaxial structure. In this case, the reason was the elastic energy front propagation to the support areas on the lower face, which passed mainly through the grain boundaries oriented parallel to its direction. Thus, there were two competing mechanisms of elastic energy absorption: (i) the microstructure with horizontally oriented grains realized micro-plastic shears in the grain bodies (according to Mott's theory); (ii) the microstructure with vertically oriented grains was subjected by the intensive development of rotational wave of defect flows along the grain boundaries (according to the theory of V.E. Panin—V.E. Egorushkin).

The general tendency towards stabilization of the proportions of elastic and torsional energy for all types of the grain microstructures and temperatures was due to the fact that there was a gradual decrease in the energy gradients and moments of forces both inside the grains and at their boundaries upon loading because of the establishment of an equilibrium regime for the specimen macro-volume. This effect was not only due to the balancing of elastic macro-stresses, but also because of the achievement of a certain threshold value of the torsional energy, when the material was saturated with dislocation microstructures and exhausted the ductility margin [63]. In this sense, the transition point on the curves, shown in Figure 11, to the horizontal section ('saturation') indicated the beginning of the pre-fracture stage, the features of which are discussed below.

The magnitude of the modulus of the specific moment of force was analogous to the potential energy of torsion, which later (in the process of increasing the density of dislocations and the formation of slip lines) was transformed into the work spent to plastic strain. The time dependences presented in Figure 12, a show a gradually decaying increase in the mean absolute values of the moment of force. For a more detailed study of the processes of the curvature formation at the grain boundaries, time dependences were also built for the three components of the moment of force for all simulation specimens (Figure 12b–d). Since the relationship of the separate components of the moment of force was characterized exclusively by structural features, these graphs were plotted only for the temperature of +20 °C.

The maximum mean values of the moduli of the moments of forces and its components were shown by the specimen with the fine-grained microstructure. Again, this fact was due to the largest area of the grain boundaries in comparison with other investigated microstructures. The plots of the mean moduli of the moments of forces were almost identical for all other specimens with the same mean grain volume. Nevertheless, the dependences of the mean values of the moduli of various components of the moments were significantly different from each other. Accordingly, the grain shapes affected the spatial distribution pattern of the torsion zones of the material around one or another coordinate axis. For example, torsion of the material under simulation around the (horizontal)  $x$  axis was more intense during deformation if the grains were stretched in the horizontal direction. At the same time, 'swirls' around the vertical  $z$  axis was most pronounced in the microstructure with vertically oriented grains. This confirmed the above thesis about the competing mechanisms of the elastic energy dissipation. In the case of predominantly horizontally oriented grains (across the load direction), the dissipation was due to the formation of defect microstructures in the grain bodies. In the case of the grain microstructure oriented coaxially to impact load, the predominant development of the rotational wave of defect flows was along the grain boundaries.



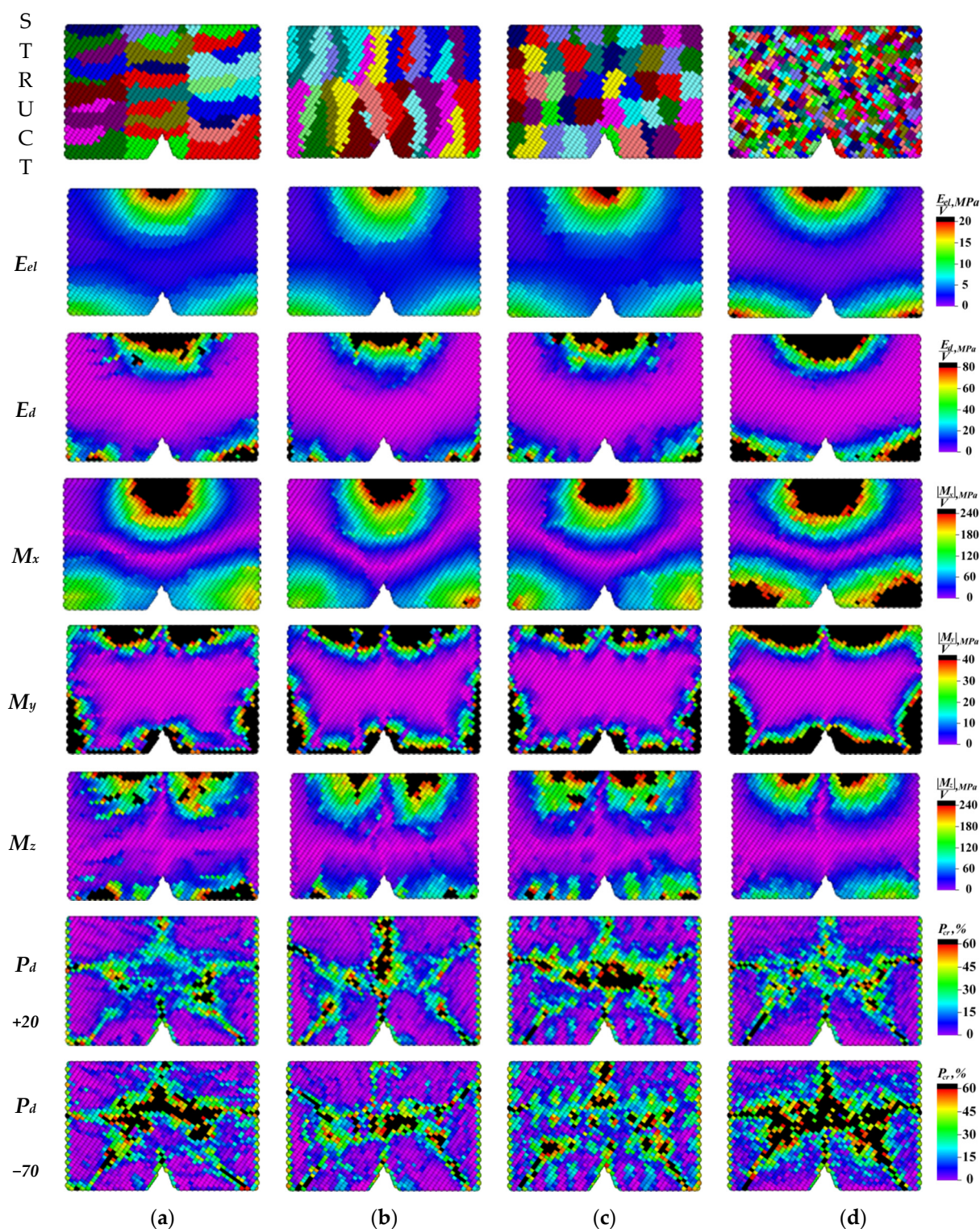
**Figure 12.** Graphs of time dependences of the mean moduli of the specific moment of force (a) and its components X (b), Y (c), Z (d) for the specimens with grain sizes  $30 \times 6 \times 6$  (blue, 4),  $6 \times 6 \times 30$  (red, 2),  $9 \times 10 \times 12$  (green, 3) and  $3 \times 3 \times 3$   $\mu\text{m}^3$  (black, 1) at the temperature of  $+20$  °C.

Figure 13 shows maps of the grain microstructure on the front faces of the simulation specimens and the distribution patterns of the specific elastic energy, the specific dissipated torsion energy, the mean moduli of the components of the specific moment of force, and the probabilities of micro-crack initiation on these faces after  $20 \mu\text{s}$ .

The elastic energy propagation dynamics, presented in Figure 13, demonstrated a gradual inflow of energy from the load application area at the upper face (as well as from the support side located at the lower face) into the specimens. The grain microstructure had a significant effect on the elastic energy front propagation rate (the grain boundaries slowed it down). For this reason, the ‘saturation of the material with elastic energy’ was slowest in the fine-grained microstructure compared to the other ones. The elastic energy dissipation was most intense both at the point of the impact load application and near the supports at the lower face.

It should be noted that the greatest increase in the torsional energy was in the regions of the grain boundaries, behind which this energy propagated deep into the specimen, mainly along the grain boundaries. Such a character of the torsional energy front propagation and its dissipation at the nucleation of the defect microstructures was caused by the action of the moments of forces arose as a result of the elastic energy gradients at the grain boundaries. In this case, the action of the local moments of forces also spread into the grain bodies, which subsequently caused their fragmentation and transcrystalline failure.





**Figure 13.** Specimen grain microstructure maps and the patterns of spatial distributions of specific elastic energy ( $E_{el}$ ), specific torsional energy ( $E_d$ ), components of the moment of force and the probability of micro-crack initiation at the temperatures of +20 °C and -70 °C after 20  $\mu$ s. (explanations for a–d are given in the text above).

In Figure 13, the distribution patterns of the X component absolute values of the specific moment of force  $M$  were in many respects similar to the same distributions of the specific torsional energy  $E$ . A distinctive feature of these patterns was a more significant

difference between the fine-grained and other microstructures in the degree of accumulation of the X moment component values near the support areas.

It was noted above that the Y component values of the moment of force reached their maximum values on the lateral specimen faces at the beginning stage of impact loading. Then, they were gradually increased within the specimen. As a result, quasi-symmetric patterns of these component distributions were found at the front faces of the specimens. Their type was determined by the shapes and the sizes of the grains. In the case of the Z component of the moment of force, the distribution patterns were also quasi-symmetric about the central vertical axis (the axis of the applied impact load). In this case, they gradually ‘stabilized’ as they were loaded. The maximum modulus values were achieved in the grain bodies located near the loaded upper and lower faces. This was due to the fact that the grain bodies accumulated the torsional energy as a result of the action of the moments of forces generated along the grain boundaries and the specimen surface. On the other hand, the grain boundaries were channels for the propagation of rotational wave flows deep into the material.

The probability distributions of micro-crack initiation, calculated by formula (6), showed that the shapes and the sizes of the grains also had a significant effect on the nature of micro-cracking (Figure 13;  $P_d$ ). Thus, there was an extremely effective damping effect of the microstructure with the preferred grain orientation perpendicular to the loading axis. Despite the “potentially incipient macro-cracks” in the top notch and punch regions, the material resisted to accumulation of micro-ones in the specimen core to a critical level regulated by the concentration ratio (9). In the case of vertically oriented grains, the rapid formation of the ‘main crack’ through the entire specimen was very likely (Figure 13b;  $P_d$ ). In the case of the equiaxial coarse-grained microstructure, an even more critical result was found (Figure 13c;  $P_d$ ): the total specimen area affected by micro-cracking was the largest of all the studied microstructures. Moreover, the macro-crack propagation direction was difficult to predict in this case. It could be assumed that further fracture dynamics was of a bifurcation nature. The refinement of the microstructure enabled to reduce the mean level of the micro-crack initiation probability (Figure 13d;  $P_d$ ). Nevertheless, no apparent damping effect was observed in the specimen volume due to the low efficiency of the microstructure refinement without its anisotropy.

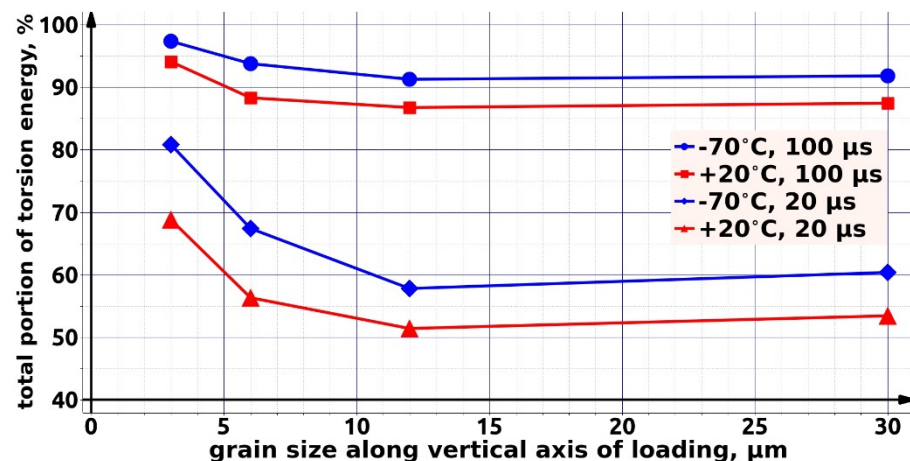
The final stage of the loading process simulation was discussed above, where the torsion energy curve was transformed into the horizontal section corresponding to saturation with defects that correlated with the pre-fracture stage beginning. This section of the material saturation with the dislocation microstructures and broken interatomic bonds was in reality the beginning stage of microcrack accumulation into larger ones with their subsequent rapid propagation. The start of this step was regulated by the concentration law, expressed by Equation (7). Upon completion of the stage of accumulation of the defect microstructures, the propagation rate of the number of microcracks ( $N_0$ ) in formula (9) ceased to be a constant and began to rapidly increase. It should be noted that development of a special algorithm described accumulation of the microcracks into the macro-ones is a separate issue, which is planned for implementation by the authors in future works.

Similar calculations of the probability distribution of microcrack initiation were carried out for identical simulation specimens at the temperature of  $-70\text{ }^{\circ}\text{C}$ . The simulation results are also shown in Figure 13 (bottom line). They indicated a greater ability to generate damage with decreasing temperature. In this case, the fine-grained specimen showed the greatest ability to low-temperature embrittlement (Figure 13d). The area of supposed microcracks was the largest of all specimens at the temperature of  $-70\text{ }^{\circ}\text{C}$ . A very unexpected result was observed for the lamellar microstructure with grains oriented along the load axis. While at the temperature of  $+20\text{ }^{\circ}\text{C}$  there was a clear growth of the longitudinal crack across the entire specimen, micro-cracks dispersed without agglomerating into macro-ones at  $-70\text{ }^{\circ}\text{C}$ . In this case, an intense accumulation of the micro-damages was along the grain boundaries in the lamellar microstructure oriented across the loading axis, which was more resistant to impact loads at temperatures above zero. This fact highlighted requirements



to form multilayer composite materials with alternating ‘parallel’ and ‘perpendicular’ lamellar microstructures.

Since the relative torsional energy was directly related to the activation volume of fracture and, ultimately, to the micro-crack initiation probability, the authors considered it appropriate to provide a summary graph of the relative torsion energy for all simulation specimens and both temperatures (Figure 14).



**Figure 14.** Summary plots of the relative torsional energy for all simulation specimens at two loading stages.

According to Figure 14, the greatest difference in the behavior of various specimens, as well as the effect of temperature, was at the first stages of impact loading. With an increase in the load duration, the distinction between the microstructures decreased due to the tendency of the system to equilibrium. This was also consistent with the experimental fact that the ductile-brittle transition phenomenon was not observed under the quasi-static test conditions. This transition was not only a microstructure-dependent phenomenon (from a type of the crystal lattice) but was also an extremely non-equilibrium process in which strains of various natures competed with each other in grain bodies.

Thus, the performed theoretical studies enabled to trace explicitly the influence of the sizes and the orientations of the grain microstructure on the evolution of elastic and dissipated strains, including the crack initiation probability. It should be emphasized once again that the authors did not strive to obtain complete agreement between theoretical and experimental data, if only because of the specificity of the model approach they developed. However, the obtained results indicated the correctness of the chosen direction, within which research on fracture process simulation will be continued for structural materials after the HR processing.

By way of the discussion, authors would like to refer the interpretation of helical rolling effect given in [64]. The effect of a lattice curvature on impact toughness at low temperatures has been studied for the 09Mn2Si and 10Mn2VNbAl high-strength low-alloy (HSLA) steels. It was shown that the formation of the lattice curvature of several degrees/micrometer significantly increases the low-temperature impact toughness of the HSLA steels, regardless of the alloying element compositions. It is suggested that the perlite–bainite phase transformation develops under helical rolling at a temperature of 850 °C. As a result, bainite is formed in the lattice curvature zones characterized by the presence of interstitial mesoscopic structural states where random rotational deformation modes could develop. This was responsible for drastic improvement of the temperature impact toughness of the HSLA steels.

For the above described reasons, the revealed increasing of impact toughness should be attributed to the structural transformations at various scale levels, including electron spectra modification (nanoscale), ensuring dislocation ductility and inverse perlite–bainite



transformation (microscale level), lamellar grain structure formation (mesoscale), gradient pattern of structure variation (macroscale).

## 6. Conclusions

A study of the microstructure and the mechanical properties of the 09Mn2Si HSLA steel was carried out after the three- and five-pass HR processing. After three passes, toughness increased by 10% at the temperature of +20 °C, while it enhanced by 25% after five passes. When the temperature dropped below −20 °C, the impact toughness of the as-received steel specimens sharply decreased, and the difference (in comparison with the specimens after the HR) reached 60%.

It was found that the contribution of the crack initiation energy  $A_i$  to the total energy consumption decreased for all specimens with reducing test temperature. Higher values were typical for the as-received steel, while lower levels were observed for ones after the five-pass HR. Thus, the main reason for the increase in impact toughness of the steel after the HR, according to the test results of work spent to crack initiation and propagation, was the more effective fracture resistance at both stages. The lowest values of the spent energy and its sharper decrease with reducing temperature were for the as-received steel specimens.

It was shown that the HR processing caused the change in the hierarchy of the microstructure of grains, their shapes and sizes, which affected the mechanical properties of the steel, including impact toughness. A decrease in the grain sizes resulted in a significant increase in the structural heterogeneity of the steel at the micro-level. This resulted in two main effects. First, the crack initiation energy intensity decreased. Second, the crack propagation energy intensity increased, which was associated with high energy spent to crack propagation in the micro-composite with the lamellar mesostructure.

It was shown that the greatest difference in impact toughness, which manifested itself at the temperature of −70 °C, was reflected in the corresponding fracture surfaces. At the micro-level, the surface of the crack initiation and propagation areas was formed due to shear strains for the as-received steel specimen. The fractographic surface had a smoothed, 'structureless' appearance with micro-fractures mainly at the boundaries of the deformed micro-zones. The fractographic relief reflected the ductile-brittle pattern of fracture. The specimens after the HR retained the shear fracture (there were no ductile tear dimples). The surface had the more 'smeared' and 'structureless' character. No noticeable distinctions were found in the fractographic pattern of the specimens after the three- and five-pass HR.

Within the framework of the theoretical section, the advanced algorithm was proposed to calculate the probabilities of the initiation of microcracks under impact loading, which was based on the kinetic theory of fracture. Analysis of the graphs of the changes in various energy components enabled to identify the points of saturation of the specimen with the defect microstructures and, thereby, to determine the macro-crack initiation beginning. The nature of the distribution of the probabilities of microcrack nucleation at the temperature of +20 °C showed the highest energy intensity and damping ability of the anisotropic microstructure with grains oriented across the impact load axis. On the other hand, the lamellar microstructure with grains oriented along the load axis demonstrated the ability to disperse microcracks at low temperatures. For this reason, the authors concluded that it is advisable to form layered 'transverse-longitudinal' lamellar mesostructures in order to achieve optimal parameters of strength and fracture toughness, including at low temperatures.

A comparison of the simulation results immediately after the load application and at the end of loading showed the importance of considering the non-equilibrium process of energy distribution in the analysis of the behavior of various microstructures, especially at under-zero temperatures.

To develop a theoretical method, the authors plan to assess the energy capacity of various types of (experimentally detected) grain microstructures on the basis of the

concentration criterion, primarily from the standpoint of the ability to accumulate micro-defects without loss of bearing capacity (failure).

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## References

1. Graef, M.K.; Hillenbrand, H.G.; Heckmann, C.J. High-strength large-diameter pipe for long-distance high-pressure gas pipelines. *Int. J. Offshore Polar Eng.* **2004**, *14*, 69–73.
2. Pontremoli, M. A new generation of ultra-high strength X100/X120 pipelines: A breakthrough for economic long-distance gas transportation. In Proceedings of the 15th International Offshore and Polar Engineering Conference, Seoul, Korea, 19–24 June 2005; pp. 31–39.
3. Enrico, L.; Chris, N.M.; Ray, L.S. Impact characterization of line pipe steels by means of standard, sub-size and miniaturized charpy specimens. In *NIST Technical Note 1865*; NIST: Boulder, CO, USA, 2015. [\[CrossRef\]](#)
4. Gorynin, I.V.; Khlusova, E.I. Nanostructured Steels for Developing the Shelf of the Arctic Ocean. *Her. Russ. Acad. Sci.* **2010**, *80*, 507–513. [\[CrossRef\]](#)
5. Yan, J.-B.; Liew, J.Y.R.; Zhang, M.-H.; Wang, J.-Y. Mechanical properties of normal strength mild steel and high strength steel S690 in low temperature relevant to Arctic environment. *Mater. Des.* **2014**, *61*, 150–159. [\[CrossRef\]](#)
6. Avila, J.A.; Lucon, E.; Sowards, J.; Mei, P.R.; Ramirez, A.J. Assessment of Ductile-to-Brittle Transition Behavior of Localized Microstructural Regions in a Friction-Stir Welded X80 Pipeline Steel with Miniaturized Charpy V-Notch Testing. *Metall. Mater. Trans. A* **2016**, *47A*, 2855–2865. [\[CrossRef\]](#)
7. Yu, C.; Yang, T.C.; Huang, C.Y.; Shiue, R.K. Low-Temperature Toughness of the Austempered Offshore Steel. *Metall. Mater. Trans. A* **2016**, *47A*, 4777. [\[CrossRef\]](#)
8. Zia-Ebrahimi, F. *Ductile-to-Brittle Transition in Steel Weldments for Arctic Microstructures*; Fracture and Deformation Division, Center for Materials Science National Bureau of Standards: Boulder, CO, USA, 1985; NBSIR 85-3020.
9. Chen, J.; Tang, S.; Liu, Z.-Y.; Wang, G.-D. Microstructural characteristics with various cooling paths and the mechanism of embrittlement and toughening in low-carbon high performance bridge steel. *Mater. Sci. Eng. A* **2013**, *559*, 241–249. [\[CrossRef\]](#)
10. Pluinage, G.; Amara, B.M.; Capelle, J.; Azari, Z. Role of constraint on ductile brittle transition temperature of pipe steel X65. *Procedia Mater. Sci.* **2014**, *3*, 1560–1565. [\[CrossRef\]](#)
11. Bruchhausen, M.; Holmstrom, S.; Lapetite, J.-M.; Ripplinger, S. On the determination of the ductile to brittle transition temperature from small punch tests on Grade 91 ferritic-martensitic steel. *Int. J. Press. Vessel. Pip.* **2017**, *155*, 27–34. [\[CrossRef\]](#)
12. Niu, Y.; Jia, S.; Liu, Q.; Tong, S.; Li, B.; Ren, Y.; Wang, B. Influence of Effective Grain Size on Low Temperature Toughness of High-Strength Pipeline Steel. *Materials* **2019**, *12*, 3672. [\[CrossRef\]](#)
13. Takahashi, O.; Shibui, Y.; Xu, P.G.; Harjo, S.; Suzuki, T.; Tomota, Y. Microstructural Features and Ductile-Brittle Transition Behavior in Hot-Rolled Lean Duplex Stainless Steels. *Quantum Beam Sci.* **2020**, *4*, 16. [\[CrossRef\]](#)
14. Valiev, R.Z.; Langdon, T. Principles of equal-channel angular pressing as a processing tool for grain refinement. *Prog. Mater. Sci.* **2006**, *51*, 881–981. [\[CrossRef\]](#)
15. Naik, S.N.; Walley, S.M. The Hall–Petch and inverse Hall–Petch relations and the hardness of nanocrystalline metals. *J. Mater. Sci.* **2020**, *55*, 2661–2681. [\[CrossRef\]](#)
16. Cao, S.; Liu, J.; Zhu, L.; Li, L.; Dao, M.; Lu, J.; Ritchie, R.O. Nature-Inspired Hierarchical Steels. *Sci. Rep.* **2018**, *8*, 5088. [\[CrossRef\]](#)
17. Wu, X.; Zhu, Y. Heterogeneous materials: A new class of materials with unprecedented mechanical properties. *Mater. Res. Lett.* **2017**, *5*, 527–532. [\[CrossRef\]](#)
18. Wu, X.; Jiang, P.; Chen, L.; Zhang, J.; Yuan, F.; Zhu, Y. Synergetic strengthening by gradient structure. *Mater Res Lett.* **2014**, *2*, 185–191. [\[CrossRef\]](#)

19. Yang, M.; Pan, Y.; Yuan, F.; Zhu, Y.; Wu, X. Back stress strengthening and strain hardening in gradient microstructure. *Mater. Res. Lett.* **2016**, *4*, 145–151. [\[CrossRef\]](#)
20. Chen, A.; Shi, S.; Tian, H.; Ruan, H.; Li, X.; Pan, D.; Lu, J. Effect of warm deformation on microstructure and mechanical properties of a layered and nanostructured 304 stainless steel. *Mater. Sci. Eng. A* **2014**, *595*, 34–42. [\[CrossRef\]](#)
21. Launey, M.E.; Ritchie, R.O. On the Fracture Toughness of Advanced Materials. *Adv. Mater.* **2009**, *21*, 2103–2110. [\[CrossRef\]](#)
22. Ritchie, R.O. The conflicts between strength and toughness. *Nat. Mater.* **2011**, *10*, 817–822. [\[CrossRef\]](#)
23. Knott, J. Brittle fracture in structural steels: Perspectives at different size-scales. *Phil. Trans. R. Soc. A* **2015**, *373*. [\[CrossRef\]](#)
24. Panin, V.E.; Shulepov, I.A.; Derevyagina, L.S.; Panin, S.V.; Gordienko, A.I.; Vlasov, I.V. Nanoscale Mesoscopic Structural States in Low-Alloy Steels for Martensitic Phase Formation and Low-Temperature Toughness Enhancement. *Phys. Mesomech.* **2020**, *23*, 376–383. [\[CrossRef\]](#)
25. Tsay, K.; Arbuz, A.; Gusseyinov, N.; Nemkaeva, R.; Ospanov, N.; Krupen'kin, I. Refinement of the microstructure of steel by cross rolling. *J. Chem. Technol. Metall.* **2016**, *51*, 385–392.
26. Wang, Y.L.; Molotnikov, A.; Diez, M.; Lapovok, R.; Kim, H.-E.; Wang, J.T.; Estrin, Y. Gradient microstructure produced by three roll planetary milling: Numerical simulation and microstructural observations. *Mater. Sci. Eng. A* **2015**, *639*, 165–172. [\[CrossRef\]](#)
27. Skripalenko, M.M.; Galkin, S.P.; Karpov, B.V.; Romantsev, B.A.; Kaputkina, L.M.; Danilin, A.V.; Patrin, P.V. Forming Features and Properties of Titanium Alloy Billets After Radial-Shear Rolling. *Materials* **2019**, *12*, 3179. [\[CrossRef\]](#)
28. Bogatov, A.A.; Panov, E.I. Effect of Stress-strain State during Helical Rolling on Metal and Alloy microstructure and Ductility. *Metallurgist* **2013**, *57*, 434–441. [\[CrossRef\]](#)
29. Surikova, N.S.; Panin, V.E.; Narkevich, N.A.; Mishin, I.P.; Gordienko, A.I. Formation of a multilevel hierarchical mesosubstructure by cross rolling and its influence on the mechanical behavior of austenitic steel. *Phys. Mesomech.* **2018**, *21*, 36–47. [\[CrossRef\]](#)
30. Liu, B.; Huang, L.; Geng, L.; Yin, F. Multiscale Hierarchical microstructure and Laminated Strengthening and Toughening Mechanisms. In *Lamination—Theory and Application*; Intechopen: London, UK, 2018; pp. 19–49. [\[CrossRef\]](#)
31. Wadsworth, J.; Lesuer, D.R. Ancient and modern laminated composites: From the great Pyramid of Gizeh to Y2K. *Mater. Charact.* **2000**, *45*, 289–313. [\[CrossRef\]](#)
32. Wadsworth, J.; Sherby, O.D. On the bulat (Damascus steels). *Annu. Bull. Met. Mus.* **1979**, *4*, 7–23.
33. Wadsworth, J.; Kum, D.W.; Sherby, O.D. Welded Damascus steels and a new breed of laminated composites. *Met. Prog.* **1986**, *129*, 61–67.
34. Lesuer, D.R.; Syn, C.K.; Sherby, O.D.; Wadsworth, J.; Lewandowski, J.J.; Hunt, W.H., Jr. Mechanical behavior of laminated metal composites. *Int. Mater. Rev.* **1996**. [\[CrossRef\]](#)
35. Kum, D.W.; Oyama, T.; Wadsworth, J.; Sherby, O.D. The impact properties of laminated composites containing ultrahigh carbon (UHC) steels. *J. Mech. Phys. Solids* **1983**, *31*, 173–186. [\[CrossRef\]](#)
36. Derevyagina, L.S.; Surikova, N.S.; Gordienko, A.I.; Korznikov, A.V. Structural Changes in Steel 12GBA During Warm Rolling and Resistance to Brittle Fracture at Negative Temperatures. *Met. Sci. Heat Treat.* **2018**, *59*, 721–728. [\[CrossRef\]](#)
37. Panin, V.E.; Derevyagina, L.S.; Panin, S.V.; Shugurov, A.R.; Gordienko, A.I. The Role of Nanoscale Stain-Induced Defects in the Sharp Increase of Low-Temperature Toughness in Low-Carbon and Low-Alloy Steels. *Mater. Sci. Eng. A* **2019**, *768*, 138491. [\[CrossRef\]](#)
38. Panin, V.E.; Derevyagina, L.S.; Lemeshev, N.M.; Panin, A.V.; Kazachenok, M.S.; Korznikov, A.V. On the Nature of Low-Temperature Brittleness of BCC Steels. *Phys. Mesomech.* **2014**, *17*, 89–96. [\[CrossRef\]](#)
39. Derevyagina, L.S.; Gordienko, A.I.; Kashiro, P.O. The Influence of Regimes of Thermomechanical Treatment on the Structural-Phase State, Mechanical Properties, and Fracture Toughness of Low-Alloy Steel API 5L Grade X52. *Russ. Phys. J.* **2019**, *61*, 1971–1977. [\[CrossRef\]](#)
40. Panin, S.; Vlasov, I.; Maksimov, P.; Moiseenko, P.; Maruschak, P.; Yakovlev, A.; Schmauder, S.; Berto, F. Increasing Fatigue Life of 09Mn2Si Steel by Helical Rolling. Theoretical-Experimental Study on Governing Role of Grain Boundaries. *Materials* **2020**, *13*, 4531. [\[CrossRef\]](#)
41. Chi, G.; Liu, H.; Yi, D. Drastic improvement in elongation and impact toughness of Ti–Al–V–Mo–Zr alloy tube via three-step heat treatment. *Mater. Lett.* **2021**, *284*, 128925. [\[CrossRef\]](#)
42. Lin, Y.; Yu, Q.; Pan, J.; Duan, F.; Ritchie, R.O.; Li, Y. On the impact toughness of gradient-structured metals. *Acta Mater.* **2020**. [\[CrossRef\]](#)
43. Luo, H.; Wang, X.; Liu, Z.; Yang, Z. Influence of refined hierarchical martensitic microstructures on yield strength and impact toughness of ultra-high strength stainless steel. *J. Mater. Sci. Technol.* **2020**, *51*, 130–136. [\[CrossRef\]](#)
44. Maruschak, P.O.; Danyliuk, I.M.; Bishchak, R.T.; Vuherer, T. Low temperature impact toughness of the main gas pipeline steel after long-term degradation. *Cent. Eur. J. Eng.* **2014**, *4*, 408–415. [\[CrossRef\]](#)
45. Panin, S.V.; Maruschak, P.O.; Vlasov, I.V.; Ovechkin, B.B. Impact toughness of 12Cr1MoV steel. Part 1—Influence of temperature on energy and deformation parameters of fracture. *Theor. Appl. Fract. Mech.* **2016**, *83*, 105–113. [\[CrossRef\]](#)
46. Kryzhanivskyy, Y.; Poberezhny, L.; Maruschak, P.; Lyakh, M.; Slobodyan, V.; Zapukhliak, V. Influence of test temperature on impact toughness of X70 pipe steel welds. *Procedia Struct. Integr.* **2019**, *16*, 237–244. [\[CrossRef\]](#)
47. Maruschak, P.; Sorochak, A.; Brezinova, J.; Guzanova, A.; Yasnii, O. Micromechanisms of ferrite-pearlite steels fracture under cyclic and impact loading. *Acta Metall. Slovaca* **2017**, *23*, 345–355. [\[CrossRef\]](#)
48. McDowell, D.L. A perspective on trends in multiscale plasticity. *Int. J. Plast.* **2010**, *26*, 1280–1309. [\[CrossRef\]](#)

- 
49. Roters, F.; Eisenlohr, P.; Hantcherli, L.; Tjahjanto, D.D.; Bieler, T.R.; Raabe, D. Overview of constitutive laws, kinematics, homogenization and multiscale methods in crystal plasticity finite-element modeling: Theory, experiments, applications. *Acta Mater.* **2010**, *58*, 1152–1211. [[CrossRef](#)]
  50. Trusov, P.V.; Sharifullina, E.R.; Shveykin, A.I. Multilevel Model for the Description of Plastic and Superplastic Deformation of Polycrystalline Materials. *Phys. Mesomech.* **2019**, *22*, 402–419. [[CrossRef](#)]
  51. Panin, S.V.; Maruschak, P.O.; Vlasov, I.V.; Moiseenko, D.D.; Maksimov, P.V.; Berto, F.; Bishchak, R.T.; Vinogradov, A. The role of notch tip shape and radius on deformation mechanisms of 12Cr1MoV steel under impact loading. Part 2. Influence of strain localization on fracture and numeric simulations. *Fatigue Fract. Eng. Mater. Struct.* **2017**, *40*, 1838–1853. [[CrossRef](#)]
  52. Panin, S.V.; Moiseenko, D.D.; Maksimov, P.V.; Vlasov, I.V.; Byakov, A.V.; Maruschak, P.O.; Berto, F.; Schmauder, S.; Vinogradov, A. Influence of Energy Dissipation at the Interphase Boundaries on Impact Fracture Behaviour of a Plain Carbon Steel. *Theor. Appl. Fract. Mech.* **2018**, *97*, 478–499. [[CrossRef](#)]
  53. Moiseenko, D.D.; Panin, V.E.; Elsukova, T.F. Role of local curvature in grain boundary sliding in a deformed polycrystal. *Phys. Mesomech.* **2013**, *16*, 335–347. [[CrossRef](#)]
  54. Moiseenko, D.D.; Maruschak, P.O.; Panin, S.V.; Maksimov, P.V.; Vlasov, I.V.; Berto, F.; Schmauder, S.; Vinogradov, A. Effect of structural heterogeneity of 17Mn1Si steel on the temperature dependence of impact deformation and fracture. *Metals* **2017**, *7*, 280. [[CrossRef](#)]
  55. Mott, N.F. Slip at Grain Boundaries and Grain Growth in Metals. *Proc. Phys. Soc.* **1948**, *60*, 391–394. [[CrossRef](#)]
  56. Zhurkov, S.N. Kinetic Concept of the Strength of Solids. *Int. J. Fract. Mech.* **1965**, *1*, 311–323. [[CrossRef](#)]
  57. Mishnaevsky, L., Jr.; Brondsted, P. Modeling of fatigue damage evolution on the basis of the kinetic concept of strength. *Int. J. Fract.* **2007**, *144*, 149–158. [[CrossRef](#)]
  58. Hilarov, V.L. Kinetic theory of strength and a self-organized critical state in the process of fracture of materials. *Phys. Solid State* **2005**, *47*, 832–835. [[CrossRef](#)]
  59. Mishnaevsky, L., Jr. *Damage and Fracture of Heterogeneous Materials*; CRC Press: Boca Raton, FL, USA, 1998; 200 p.
  60. Suresh, S. *Fatigue of Materials*; Cambridge University Press: Cambridge, UK, 1998; 704p.
  61. Chernov, V.M.; Kardashev, B.K.; Moroz, K.A. Cold brittleness and fracture of metals with various crystal lattices: Dislocation mechanisms. *Tech. Phys.* **2016**, *61*, 1015–1022. [[CrossRef](#)]
  62. Ivanov, V.V.; Egorov, P.V.; Klimov, V.I.; Kolpakova, L.A.; Malyshin, A.A.; Tugolokova, L.F.; Yunnikov, V.V.; Berveno, V.P. Determination of the kinetic strength constants and of the critical size of destruction of composite materials on the basis of recording of pulse electromagnetic radiation during their destruction. *J. Appl. Mech. Tech. Phys.* **1994**, *35*, 628–633. [[CrossRef](#)]
  63. Sangid, M.D. The physics of fatigue crack initiation. *Int. J. Fatigue* **2013**, *57*, 58–72. [[CrossRef](#)]
  64. Panin, V.E.; Egorushkin, V.E.; Panin, S.S. Mechanism for improving low temperature impact toughness and fatigue durability of high-strength low-alloy steels for applications in Arctic region. *Acta Mech.* **2020**. [[CrossRef](#)]