

Review

Flux-Grown Piezoelectric Materials: Application to α -Quartz Analogues

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Received: 22 April 2014; in revised form: 11 June 2014 / Accepted: 12 June 2014 /

Published: 23 June 2014

Abstract: Using the slow-cooling method in selected MoO₃-based fluxes, single-crystals of GeO₂ and GaPO₄ materials with an α-quartz-like structure were grown at high temperatures ($T \ge 950$ °C). These piezoelectric materials were obtained in millimeter-size as well-faceted, visually colorless and transparent crystals. Compared to crystals grown by hydrothermal methods, infrared and Raman measurements revealed flux-grown samples without significant hydroxyl group contamination and thermal analyses demonstrated a total reversibility of the α -quartz $\leftrightarrow \beta$ -cristobalite phase transition for GaPO₄ and an absence of phase transition before melting for α -GeO₂. The elastic constants C_{IJ} (with I, Jindices from 1 to 6) of these flux-grown piezoelectric crystals were experimentally determined at room and high temperatures. The ambient results for as-grown α-GaPO₄ were in good agreement with those obtained from hydrothermally-grown samples and the two longitudinal elastic constants measured versus temperature up to 850 °C showed a monotonous evolution. The extraction of the ambient piezoelectric stress contribution e₁₁ from the C_{11}^{D} to C_{11}^{E} difference gave for the piezoelectric strain coefficient d_{11} of flux-grown α -GeO₂ crystal a value of 5.7(2) pC/N, which is more than twice that of α -quartz. As the α-quartz structure of GeO₂ remained stable up to melting, a piezoelectric activity was observed up to 1000 $\,^{\circ}$ C.

Keywords: single crystal; GaPO₄; GeO₂; SiO₂; raman; infrared; Brillouin; growth from high temperature solution; differential scanning calorimetry (DSC); X-ray diffraction

1. Introduction

Piezoelectric materials that can operate under very high temperature without degradation are sought for the control of structure materials and control system in turbines, engines, nuclear reactors, *etc.* [1,2]. However, the use of a piezoelectric material at elevated temperature presents many challenges such as possible phase transition, chemical degradation or structural defect propagation which can cancel or lead to instability of the piezoelectric properties.

Non-pyroelectric single crystals with the α -quartz-like structure exhibiting both higher piezoelectric constants and a higher thermal stability as compared to α -quartz would be promising materials to build miniaturized high temperature piezoelectric-operated devices without cooling. In the XO₂ (X = Si, Ge) and MPO₄ (M = Fe, Al, Ga, B) family, the α -quartz-like structure is composed of either only XO₄ corner-shared tetrahedra or of both MO₄ and PO₄ tetrahedra forming a trigonal system [3–16]. The α -phase ($P3_121$ or $P3_221$, respectively left-handed or right-handed, Z=3) is derived from the β -phase ($P6_222$ or $P6_422$) by a symmetry loss induced by a tilt of the tetrahedra around the b-axis. Two structural distortion parameters exist to quantify this phenomenon: the average intertetrahedral bridging angle θ (X–O–X or M–O–P) and the tetrahedral tilt angle δ [11,14,17–20]. Based on experimental characterizations and theoretical studies on α -quartz analogues, it was demonstrated that the higher piezoelectric properties of α -GaPO₄ and α -GeO₂ materials compared to α -SiO₂, α -FePO₄ and α -AlPO₄ compounds were directly related to their structural distortion with respect to the β -quartz structure type (for β -SiO₂ θ = 153.3° and δ = 0° at 575 °C) [11,14,17,21–27].

In addition, the well-known α - β phase transition, which appears around 573 °C in α -SiO₂ ($\theta = 144.2$ °) [19,26] does not occur when the tilt angle δ is over 22° (leading to θ under 136 °) [9,13,14,22,25]. In other words, the transition from an α phase to a β phase is absent for α -GeO₂ and α -GaPO₄ crystallized materials ($\theta = 130.04$ ° and $\theta = 134.60$ °, respectively) [13,14,17,22,28,29].

Under ambient conditions, GeO_2 exhibits two persistent forms of differing anion coordination around the central cation: the α -quartz-type modification (trigonal) with 4:2 coordination [16] and the rutile-type modification (tetragonal $P4_2$ /mnm, Z=2) with 6:3 coordination [4,11,30–34]. Naturally occurring, GeO_2 is known to be more stable in the rutile structure than in the α -quartz structure. The transformation from the rutile-like form to the α -quartz-type form has been reported to occur under normal atmospheric pressure in the 1024–1045 $^{\circ}$ C temperature range [11,35]. The kinetics of the allotropic transformation from a trigonal to a tetragonal structure is extremely low, and non-appreciable without the presence of any catalyst such as traces of GeO_2 rutile-like form, chlorides or water [11,36–39].

A transition from the thermodynamically stable α -quartz-like GaPO₄ phase to the β -cristobalite modification is known to occur above 960 °C [13,40–46]. This high temperature allotropic transition does not permit the growth of α -GaPO₄ crystals by conventional melt techniques. Single crystals of α -GaPO₄ were grown using epitaxial hydrothermal-based methods in a retrograde-solubility range at T < 250 °C and in the range of direct temperature dependence of solubility at T > 300 °C from highly corrosive acid solutions [40,43,47–62]. However, for gallium orthophosphate material, it seems quite difficult to obtain very high quality crystals with these solution-based growth methods. The presence of twins, dislocations and/or a quite high level of hydroxyl group incorporated via the growth medium tend to deteriorate the piezoelectric properties especially at high temperatures as physical properties are very sensitive to material perfection [43,49,50,57,63–73]. The OH impurities in α -GaPO₄ single crystals are

responsible for the "milky" appearance of the samples when exposed to high temperature (600 $^{\circ}$ C). This typical behavior is due to water trace precipitation effects and corresponds to the upper temperature limit of physical property measurements [65,69,74,75].

For GeO_2 , the temperature region of existence of the stable trigonal phase goes from 1033 °C up to the melting point of 1116 °C [25,30,76]. When epitaxially-grown in the temperature region where it is metastable, the hydrothermally-grown alpha-quartz form of GeO_2 contains high OH impurities which rather easily catalyze its return to the thermodynamically stable rutile-like structure when heated as low as 180 °C [77–81].

In this context, it appeared that another suitable growth technique for oxides could be applied for the crystallization of α -GaPO₄ and α -GeO₂ materials to get single crystals with a high degree of structural and chemical perfection; the high temperature solution growth technique also known as the fluxed melt growth [82–85]. Using molten inorganic salts, or flux, playing the role of solvent, crystals of a compound can be obtained below its melting or transformation point. The principle of flux growth is based on the spontaneous nucleation that occurs when a supersaturation is obtained either upon cooling of a high temperature solution or by boiling away a volatile solvent at a fixed temperature [86]. The fluxed melt method presents some advantages compared to hydrothermal-based growth methods: the material can be crystallized at atmospheric pressure and the inorganic solvents used are water-free at high temperature.

This paper presents an overview of the main results obtained from several structural and physical characterizations undertaken on flux-grown α -GaPO₄ and α -GeO₂ piezoelectric crystals. When possible, the results are compared and discussed in the view of hydrothermally-grown α -quartz-like SiO₂, GeO₂ and GaPO₄.

2. Flux-Growth and Morphology

2.1. Flux-Grown α-GaPO₄

In high temperature solution growths, GaPO₄ presents direct solubility. Non-toxic sodium chloride with its melting point at 800 °C, is a suitable solvent to grow α -GaPO₄ single crystals under the temperature of the allotropic α -quartz to β -cristobalite phase transition [75]. A graphite crucible, filled with an appropriate mixture of GaPO₄ and NaCl was sealed in a silica ampoule under Argon. By cooling from 900 to 600 °C and pulling down the ampoule (12 mm day⁻¹) within a temperature gradient (2–5 °C cm⁻¹) or by using the gradient freeze method, α -GaPO₄ crystals having NaCl "foreign" phase as shown by X-ray powder diffraction diagrams were obtained. Thus, the accelerated crucible rotation technique was applied and relatively clear and unshaped α -GaPO₄ crystals containing flux inclusions were grown up to 5 × 5 × 2 mm³ in size with a rough surface.

From spontaneous crystallization with the slow cooling method $(1.2-5 \text{ °C h}^{-1})$ between 950 and 600 °C, colorless, transparent and flux inclusion-free as-grown α -GaPO₄ crystals of millimeter-size have been obtained in MoO₃-based solvents $X_2O:3MoO_3$ (X=Li,K) [44,87–89]. The melting point of these molybdenum-based solvents are between 500 and 600 °C which allows α -GaPO₄ growth under the allotropic phase transition temperature [44,89]. The unseeded growth experiments were done under atmospheric conditions in Pt crucibles covered (but not sealed) with a lid containing a starting mixture

of α -GaPO₄ and X₂O:3MoO₃ (X = Li, K) powder in different weight ratios. Plate-like single crystals with very smooth surface roughness were grown in K₂O:3MoO₃ having a volume up to 6 × 4 × 1 mm³. With Li₂O:3MoO₃ flux, the as-grown α -GaPO₄ crystals presented unshaped bulk morphology with a quite rough surface and volume up to 5 × 2.5 × 2 mm³ [44,87,89].

With a slow cooling rate of 0.1 °C h⁻¹ from 950 to 750 °C, followed by a cooling rate of 2 °C h⁻¹ from 750 to 600 °C, in Li₂O:3MoO₃ flux, a visually high quality crystal of 8 mm long to 3 mm large and 2.5 mm thick ($V = 60 \text{ mm}^3$) presenting two smooth major faces was grown [89,90].

All these $X_2O:3MoO_3$ flux-grown α -GaPO₄ materials crystallized in the trigonal system without any secondary phase detectable and with lattice parameters in perfect agreement with those published on hydrothermally-grown α -GaPO₄ material [7,9].

2.2. Flux-Grown α-GeO₂

The first attempt at growing an α -GeO₂ crystal by the unseeded fluxed melt method was undertaken by Finch *et al.* [91] in the late sixties. Among many compounds explored as potential solvents, Li₂O:2MoO₃ and Li₂O:2WO₃ were selected. With the slow cooling growth technique from 1100 to 950 °C at 1 °C h⁻¹ or less, colorless and well developed α -GeO₂ crystals were grown up to 3 mm on the edge. Crystals from Li₂O:2MoO₃ contained impurities or defects. X-ray powder diffraction diagrams were consistent with the hexagonal modification of GeO₂. In the majority, crystals were predominantly bonded by {1011} rhombohedral faces with {1010} prism faces incompletely developed.

In the seventies using the top seeded solution growth (TSSG) method with $Li_2O:WO_3$ as flux, a colorless α -GeO $_2$ single-crystal with a maximum final diameter of 5 mm was obtained [92]. The as-grown crystal, which did not present the hexagonal-like morphology, was not sharply facetted and homogeneously transparent (white inclusions). The experimental procedure was quite unclear and some important parameters were missing. Concerning the chemical and structural quality of this TSSG-crystal, little information was given.

Recently, unseeded slow cooling growth experiments were performed from 970 to 600 °C at a cooling rate of 1 °C h⁻¹ with a solute to flux ratio of 10/90 by weight in a Pt covered crucible (not sealed) [93,94]. Using $K_2W_2O_7$ flux, the majority of the as-grown α -GeO₂ crystals contained yellow flux inclusions. Thus, other inorganic solvents were investigated in the MoO₃-based systems since they present low melting temperature and low evaporation at high temperatures (loss of mass <1% after two weeks at T > 700 °C for $K_2Mo_4O_{13}$ flux) [93,94]. Well-faceted and visually colorless and transparent α -GeO₂ single-crystals were obtained in fluxes such as $K_2Mo_4O_{13}$, $Rb_2Mo_4O_{13}$ and $Rb_2Mo_2O_7$. The as-grown crystals had no visible flux inclusions, bubbles or cracks and presented very smooth surface roughness [16,93,94]. The α -GeO₂ millimeter-size crystals, up to 3 mm, presented either an unshaped morphology or a pseudo-cubic shape as already found for natural α -quartz [95]. In the case of pseudo-cubic crystals, the positive r {10 $\overline{1}$ 1} rhombohedral faces were prevalent with the restricted presence of the negative r {01 $\overline{1}$ 1} rhombohedral faces while m {10 $\overline{1}$ 0} faces were absent [93,94].

The crystal structure and quality of these flux-grown α -GeO₂ materials were studied by both powder and single-crystal X-ray diffraction at room temperature [16,93,94]. The refinements confirmed the crystallization of the α -quartz-type structure and the lattice parameters and volumes were in good agreement with the literature data [4,10,11,32,80]. The excellent final reliability factors of the

refinements indicated that the as-grown α -GeO₂ crystals were of high quality precluding the presence of any considerable amount of twinning [16]. The value of the Flack coefficient which is an indicator of the presence of growth portions containing mirror-image was found close to zero. The uniform coloration of the natural faces under polarized light indicated the lack of optical domains in the flux-grown α -GeO₂ single crystals [16,93].

A colorless, well facetted, highly-transparent and large-size single crystal, up to $0.5~\text{cm}^3$, of the piezoelectric phase of GeO₂ was grown by TSSG from a high temperature solution using K₂Mo₄O₁₃ as solvent [27,93]. The obtained volume made this isometric flux-grown GeO₂ single crystal, Figure 1, with the metastable α -quartz like structure, the largest reported in the literature. The macroscopic hexagonal morphology, similar to the well-known morphology of α -SiO₂, was in accordance with an α -quartz-like structure and facilitated the identification of the different natural crystallographic faces [96–98].

Figure 1. Picture of an α -GeO₂ single crystal grown by the top seeded solution growth (TSSG) technique in $K_2Mo_4O_{13}$ flux.



3. Impurities Contamination

3.1. Flux-Grown α-GaPO₄

The OH contamination of the crystalline lattice during crystallization via growth medium, which decreases the Q-factor of the resonators, has been largely reported in the literature concerning α -GaPO₄ grown by the hydrothermal method [43,52,53,55,57,68,71,99,100].

The visual estimation of the OH contamination of the α -GaPO₄ single crystals flux-grown from NaCl solvent was studied by annealing experiments [75]. The "milky" clouding was a direct estimation of the OH $^-$ concentration; the higher the OH content, the more "milky" is the sample. These NaCl flux-grown α -GaPO₄ crystals presented slight clouding and thus significant OH contamination.

In a typical room temperature non-polarized infrared transmission spectrum of an α -GaPO₄ material containing significant OH-groups, a broad and intense band between 2500 and 3600 cm⁻¹ is observed (O–H infrared region) associated with H-bonded molecular water. This broad band is superimposed upon three well-separated peaks at 3167, 3290 and 3400 cm⁻¹ and sometimes upon a sharp absorption band at 3508 cm⁻¹ related to an isolated OH-group stretching band [43,55,65,68]. The amplitude of the peak centered around 3400 cm⁻¹ is attributed in one part to the absorption by the GaPO₄ lattice and in another part to O–H stretching vibrations [65].

For $X_2O:3MoO_3$ (X = Li, K) flux-grown α -GaPO₄ materials, the collection of the non-polarized infrared data was done in transmission mode at room temperature on as-grown samples, *i.e.*, not

polished samples. Only the three bands at 3167, 3290 and 3400 cm⁻¹ were registered. The OH content was materials [44,46,88–90,101]. Compared to the spectra of hydrothermally-grown GaPO₄, they did not present the characteristic broad and intense absorption band from 3600 to 2500 cm⁻¹ due to quite strong OH contamination estimated with the extinction coefficient α calculated at 3400 cm⁻¹ from the expression $\alpha = 1/d[\log(T_{3800}/T_{3400})] - \alpha_{3400}^*$ where d represents the sample thickness in cm, T represents the % IR transmission at respectively 3800 and 3400 cm⁻¹, and α_{3400}^* represents the absorption coefficient due to intrinsic lattice vibrations of α -GaPO₄ at 3400 cm⁻¹ [65]. The α_{3400}^* was estimated to be 0.078 cm⁻¹ based on hydrothermally-grown material [65,68,102].

The calculated value of the extinction coefficient α at 3400 cm⁻¹ was close to 0.03 cm⁻¹ for $X_2O:3MoO_3$ (X = Li, K) flux-grown GaPO₄ and close to 0.15 cm⁻¹ for hydrothermally-grown material [88]. Thus, the infrared spectra correspond to $X_2O:3MoO_3$ (X = Li, K) flux-grown α -GaPO₄ samples without significant OH-content in perfect accordance with the absence of "milky" clouding of the samples even after several times heating up to 850 °C [103].

Chemical analyses of $X_2O:3MoO_3$ (X = Li, K) flux-grown α -GaPO₄ crystals were done using glow discharge mass spectrometry (GDMS) and induced coupled plasma-atomic emission spectroscopy (ICP-AES) [89]. The results demonstrated the presence of foreign chemical elements coming from the solvent with a Mo-content as high as 0.3% for GaPO₄ crystals grown in $K_2O:3MoO_3$ flux. Na, Ca, Fe, Al and Si elements were also found with a major contamination from Al element (320 ppm to 0.13%).

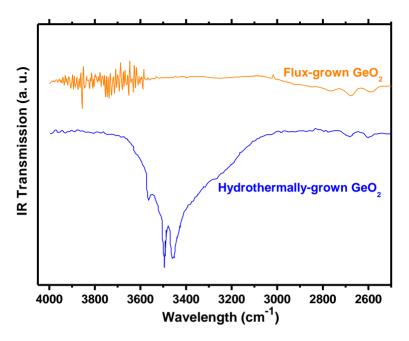
3.2. Flux-Grown α -GeO₂

As for hydrothermally-grown α -GaPO₄ single crystals, OH contamination of the lattice was reported for hydrothermally-grown α -GeO₂ [43,79,80,94]. In a non-polarized infrared spectrum collected in the transmission mode of a hydrothermally-grown α -GeO₂ material containing significant OH-groups, Figure 2, a well-pronounced broad band between 2500 and 3800 cm⁻¹ is observed with maxima at about 3455, 3500 and 3562 cm⁻¹ attributed to Ge–OH vibration groups [43,79,81,94]. For both spontaneously nucleated MoO₃-based flux-grown α -GeO₂ samples and TSSG-grown α -GeO₂ oriented plates of simple crystallographic orientations, infrared spectra characterize OH-free flux-grown α -GeO₂ crystals [27,93] as evidenced by the absence of both a broad band and sharp peaks in the 2800–3500 cm⁻¹ range, Figure 2.

Room temperature non-polarized Raman spectra of commercial α -GeO₂ powder or α -GeO₂ stain-etch sample exhibited some features around 760–780 cm⁻¹ which were found to disappear once the sample was annealed in air above 400 °C [39,93,104–106]. These features have been assigned either to oxygen vacancy complexes or to water bound to Ge–O entities or to a Ge–O stretching vibration of a water-distorted GeO₄ entity [39,104,105].

The non-polarized Raman signal of a MoO₃-based flux-grown α -quartz GeO₂ single crystal was collected in the 50–4000 cm⁻¹ range and no other modes than the four non-degenerate A1 ones and the eight doubly-degenerate E modes predicted by group theory for α -quartz-type GeO₂ (D₃ point group) were observed [93,106–108]. In agreement with infrared measurements [27,93], neither hydroxyl groups nor water inclusions were detected by Raman spectroscopy. Furthermore, these MoO₃-based flux-grown α -GeO₂ samples were heated up to 1000 °C several times without presenting the well-known milky hue attributed to the release of water from hydroxyl impurities with increasing temperature [109].

Figure 2. Infrared transmission spectra of flux- and hydrothermally-grown α -GeO₂ single crystals.



Ambient polarized Raman measurements performed on flux-grown α -quartz GeO₂ single crystals along with state of the art density functional theory (DFT) based calculations were reported [93,106,107]. An excellent agreement was obtained between experimental and theoretical Raman lines for both wavenumbers and relative intensities which permitted to unambiguously assign the symmetry and the nature of α -quartz GeO₂ modes. This Raman study pointed out that high temperature flux-grown GeO₂ single crystals of α -quartz-like structure were of high structural quality (impurities would have resulted in shifts of the Raman lines or in additional bands), and that vibrations in the α -quartz GeO₂ structure were relatively quasi-harmonic as the calculated frequencies at -273 °C were almost the same as the experimental values at 25 °C.

Optical transmittance was measured in the UV-VIS-NIR region for a plate with a surface perpendicular to the c axis prepared from a MoO₃-based TSSG-grown α -GeO₂ single crystal [27]. The transmission spectrum was taken on an as-grown crystal plate of 330 μ m in thickness. The transmission fraction, found to be over 97% in the VIS-NIR region, showed that the TSSG-grown α -GeO₂ crystal was of very good optical quality. The high chemical purity and the homogeneity of the TSSG-GeO₂ crystal were also revealed by the absence of absorption bands. The α -GeO₂ crystal exhibited transparency down to the UV region. Below 300 nm, a strong absorption was visible due to the fundamental absorption within the band gap with a cutoff at 205 nm.

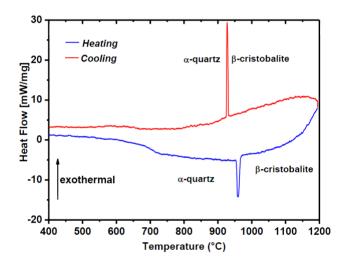
Optical transmission was reported for hydrothermally-grown α -GeO₂ crystal [79]. The spectrum showed a considerable absorption in the 0.20–0.210 μ m range associated with the presence of several bands [79].

4. Thermal Characterizations

4.1. Flux-Grown α-GaPO₄

Differential scanning calorimetric (DSC) experiments were done on flux-grown α -GaPO₄ single crystals crystallized via spontaneous nucleation in MoO₃-based fluxes as Li₂O:3MoO₃ and K₂O:3MoO₃ [44,46,89,90,101]. On heating runs up to 1200 °C (2–10 °C min⁻¹) of α -GaPO₄ as-grown single crystals or powdered samples, only a sharp endothermic peak appeared in the 950–964 °C temperature range, Figure 3, caused by the structural transition from the α -quartz GaPO₄ phase to the β -cristobalite modification stable above 960–980 °C [40–42].

Figure 3. Differential scanning calorimetric (DSC) curves of powdered flux-grown α-GaPO₄ crystals registered with a thermal cycle of 2 $^{\circ}$ C/min.



The successive cooling (-5 °C/min) data registered from 1200 °C back to room temperature of a sample containing some as-grown single crystals (different sizes and faces) presented a unique exothermic feature with double maxima at onset temperature of 937 °C for the first thermal cycle, and at onset temperature of 925 °C in the second cycle [101].

When the as-grown α -GaPO₄ crystals were powdered and sieved to a regular grain size (20 μ m), the successive cooling (-2 or -10 $^{\circ}$ C/min) curves, from 1200 to 20 $^{\circ}$ C, showed only one exothermic peak in the 942–908 $^{\circ}$ C temperature range depending on the flux composition, Figure 3.

This exothermic feature corresponded to a total transformation of the β -cristobalite GaPO₄ phase into the α -quartz phase as confirmed by the X-ray powder pattern of the end product of the DSC analysis [44,46,89,90]. This was the first time that GaPO₄ had been found, after a DSC cycle, exclusively in the α -quartz modification after cooling from the β -cristobalite phase without annealing periods.

Effectively, concerning the hydrothermally-grown α -GaPO₄ material grown at 300 °C (direct solubility), the successive cooling curve (-10 °C/min) from 1200 °C back to 20 °C, showed two main exothermic peaks [43,90]: one with an onset temperature of 910 °C due to a partial transformation of the β -cristobalite GaPO₄ phase in α -quartz phase and a second, close to 578 °C, attributed to the β -cristobalite/ α -cristobalite transition [40–42].

For hydrothermally-grown α -GaPO₄ material grown at 230 °C (indirect solubility), only one strong exothermic peak was visible on the cooling curve close to 578 °C corresponding to a total transformation of the β -cristobalite to α -cristobalite [43]. The crystallization temperature of the α -GaPO₄ phase would be an important parameter to favor the reversible α -quartz/ β -cristobalite transition.

The thermal evolution of flux-grown α -GaPO₄ single crystals grown from 950 to 600 °C in MoO₃-based solvents was also followed by Raman spectroscopy [45]. A direct α -quartz/ β -cristobalite transition was observed at a temperature close to 980 °C upon heating. Back transformation to the α -quartz-type form was found to occur readily at 920 °C in perfect agreement with the DSC experiments [44,46,89,90,101].

The temperature dependence of the cell parameters of a α -GaPO₄ sample grown in a Li₂O-3MoO₃ flux over the temperature range 950–600 °C was determined from powder X-ray diffraction data collected from 30 to 900 °C [103]. The X-ray patterns remained the same from room temperature to 900 °C. Both the lattice parameters and unit cell volumes were found to increase markedly and nonlinearly (third degree polynomial) as a function of temperature. Therefore, no indication of a structural transition was pointed out in the studied temperature range (30–900 °C) in perfect agreement with previous reported results obtained from DSC [44,46,89,90,101] or Raman studies [45,89] concerning MoO₃-based flux-grown α -GaPO₄ crystals.

4.2. Flux-Grown α -GeO₂

In the temperature range from room temperature to 1200 °C, the DSC heating-curve of GeO_2 crystals with the α -quartz structure grown by the spontaneous nucleation method in selected fluxes (MoO₃-based compounds) showed no other peak than an endothermic feature attributed to the melting of the studied material at maximum of 1116 °C [93,94].

Powder X-ray patterns of flux-grown α -GeO₂ were registered at several temperatures from room temperature up to 1050 °C [93,109]. For each studied temperature, the whole diffraction pattern was assigned to the α -quartz phase of GeO₂. The very interesting results brought to light by these thermal analyses were that this high temperature flux-grown oxide material did not present a phase transition before melting (unlike SiO₂ with the well-known α -quartz/ β -quartz transformation close to 573 °C [19,28]) and that no secondary phases such as the GeO₂ rutile-like phase or flux-derived phases were detected.

Variable-temperature Raman spectroscopy measurements performed on high temperature flux-grown α -quartz GeO₂ single crystals in MoO₃-based solvents were reported [93,106]. Vibrations in α -GeO₂ were shown to be very slightly anharmonic as evidenced by the very low wavenumber shifts and the weak damping of the modes between room temperature and 1100 °C. In contrast with what has been observed for other α -quartz homeotypes like SiO₂ or AlPO₄, which undergo an α -quartz to β -quartz phase transition [19,28], neither phase transitions nor a second phase were detected by this Raman study from room temperature to 1100 °C. First-principle calculations with the ABINIT code [93,106,107] revealed the absence of the tetrahedral libration mode in the α -quartz-like structure of GeO₂ which explained the very low degree of thermally-induced dynamic disorder registered in the 30–1100 °C range and further confirmed that the piezoelectric properties of flux-grown α -GeO₂ should not be degraded significantly up to its melting point (1116 °C).

To confirm the thermal stability (aging), an as-grown α -GeO₂ single crystal obtained by spontaneous nucleation in MoO₃-based flux was annealed in air at high temperature (800–900 °C) over several months [27]. The impact of this thermal cycle on the α -quartz-like structure of GeO₂ was followed by Raman measurement. This long annealing process validated the excellent aging behavior under very high thermal stress of the α -quartz GeO₂ obtained from the flux method since no phase transition and no evolution of the visual transparency were detected. This important result could be directly related to the high crystalline quality of the α -GeO₂ single crystals accessible with the flux growth techniques (water-free, chemical inclusion-free and rutile phase-free) as illustrated by the optical transmission curve [27].

5. Elastic Constants

5.1. Ambient Conditions

Using Voigt's notation and taking into account the crystal symmetry, the α -quartz structure (point group 32) presents six independent elastic stiffness moduli C_{IJ} (I, J indices from 1 to 6): C_{11} , C_{33} , C_{44} , C_{66} , C_{12} , C_{13} and C_{14} ($2C_{66} = C_{11} - C_{12}$). To identify the crystalline orientation, α -quartz analogues use a standard Cartesian coordinate system where the Z-axis of α -quartz analogues coincides with the crystallographic c-axis; the X-axis matches the crystallographic a-axis and the Y-axis is normal to the X and Z axes (b-axis being in the XY plane at 120 ° from a-axis).

5.1.1. Flux-Grown α-GaPO₄

By the ultrasonic method, a determination of four out of six independent single-crystal elastic stiffness constants C_{IJ}^{E} at constant electric field was undertaken on millimeter suitable shaped plates obtained from as-grown α -GaPO₄ single crystals spontaneously crystallized by slowly cooling a Li₂O:3MoO₃ flux saturated with GaPO₄ [46,89,110]. The average size of the MoO₃-based flux-grown GaPO₄ single crystals was too small [44,46,89,90,101] to get all the orientations useful for the measurements of the whole elastic constant set (ex: C_{13}^{E}).

Single-crystal high-resolution Brillouin spectroscopy experiments were carried out to measure five out of six (C^{E}_{13} is missing) elastic constants C^{E}_{IJ} of flux-grown α -GaPO₄ material [90,103]. Optical quality single-crystals of α -GaPO₄ with millimeter size were flux-grown from $X_2O:3MoO_3$ (X = Li, K) solvent in unseeded experiments over the 950–600 $^{\circ}$ C temperature range. The Brillouin measurements were done, on one hand, on plates of simple X-, Y- and Z-orientation [90] and on other hand, on an as-grown α -GaPO₄ single crystal polished as a cube with 2 mm side length showing X (100), Y (010) and Z (001) faces [103].

Table 1 gives the resulting single crystal elastic stiffness constants C^{E}_{IJ} obtained on high temperature flux-grown α -GaPO₄ materials in MoO₃-based solvents. They are compared with a reported experimental set of data concerning the hydrothermally-grown GaPO₄ single crystals as well as with reported computed values [111].

The values of the flux-grown single crystal elastic stiffness constants C^{E}_{IJ} compared with those obtained by Brillouin or pulse-echo methods on hydrothermally-grown samples, Table 1, are in good accordance while large discrepancies exist with the computed values concerning C^{E}_{11} and C^{E}_{66} .

| Table 1. Computed and experimental elastic stiffness constants C_{IJ}^{E} [GPa] of α -GaPC | Table 1. Computed an | nd experimental elastic stiffness constant | is C^{E}_{II} [GPa] of α -GaPC |
|--|----------------------|--|---|
|--|----------------------|--|---|

| Elastic constant | Hydrothermal-Growth | Computed values (-273 °C) | Flux-Growth |
|--|--|---------------------------|---|
| C^{E}_{11} | 66.58 [112] 66.35 [67] 66.60 [17] 66.58 [113] | 79.80 [111] | 64.01 [46,110] 66.37 [90] 66.52 [103] |
| $C_{12}^{E} = (C_{11}^{E} - 2C_{66}^{E})]$ | 21.81 [112] 21.65 [67] 21.80 [17] 17.38 [113] | 16.60 [111] | 13.51 [46,110] 21.45 [90] 21.04 [103] |
| C ^E ₁₄ | 3.91 [112] 4.20 [67] 3.90 [17] 5.14 [113] | 3.20 [111] | 5.52 [46,110] 4.93 [90] 5.53 [103] |
| C ^E ₃₃ | 102.13 [112] 101.31 [67] 102.10 [17] 102.13 [113] | 106.30 [111] | 103.29 [90] 103.88 [103] |
| C^{E}_{44} | 37.66 [112] 37.80 [67] 37.70 [17] 39.68 [113] | 39.90 [111] | 39.39 [46,110] 37.85 [90] 38.01 [103] |
| C ^E ₆₆ | 22.38 [112] 22.35 [67] 22.40 [17] 24.60 [113] | 31.60 [111] | 21.25 [46,110] 22.46 [90] 22.74 [103] |

5.1.2. Flux-Grown α-GeO₂

The room temperature experimental values of single-crystal elastic stiffness constants C_{IJ}^{E} at constant electric field of flux-grown α -GeO₂ in MoO₃-based solvents were determined using Brillouin scattering, Table 2 [109]. In this work, three platelets were used, defined in the standard Cartesian coordinate system as *X*-plate, *Y*-plate and *Z*-plate which respectively correspond to the (100), (2–10) and (003) (*hkl*)-crystallographic planes.

These flux-grown α -GeO₂ C_{IJ} values are compared in Table 2 with recent computed [109,114,115] and experimental elastic constant data obtained from hydrothermally-grown crystals [23,24,116]. C_{IJ}^{D} are elastic stiffness constants at constant electric displacement *i.e.*, not corrected from the piezoelectric effect e_{11}^2/ϵ_{11} where e_{11} is the piezoelectric stress constant and ϵ_{11} the dielectric tensor at constant strain.

The crystal elastic constants values of the MoO₃-based flux-grown α -GeO₂ material present close similarities with most of the published C_{IJ} data on hydrothermally-grown α -GeO₂ crystals, Table 2. Therefore, the small discrepancy, registered more especially on the C_{11} and C_{14} moduli, was attributed to the strong reduction of the OH concentration in the lattice of flux-grown α -GeO₂ crystals which was believed to induce a slightly stiffer behavior [109]. The presence of OH-interactions in the crystal would increase its ionic character and consequently softer its elastic behavior.

When compared with calculated elastic stiffness constants, the flux-grown α -GeO₂ elastic moduli presented very good agreement with computed values at 0 K of Lignie *et al.* [109], Table 2. The largest discrepancy was observed for the C_{66} (nearly 13% softer for the theoretical value) while very good accordance was found for the C_{12} and C_{14} elastic constants.

The numerical value of the e_{11}^2/ϵ_{11} piezoelectric term, deduced from the C^D_{11} to C^E_{11} difference, Table 2, is 1.29(2) GPa [109]. The authors deduced for the flux-grown α -GeO₂ material a d_{11} piezoelectric strain constant of 5.7(2) \times 10⁻¹² C/N which is in between the calculated and the reported experimental value from a hydrothermally-grown α -GeO₂ crystal, Table 3.

Table 2. Computed and experimental elastic stiffness constants C_{IJ} [GPa] of α -GeO₂.

| Elastic constant | Hydrothermal-Growth | Computed values | Flux-Growth [107] |
|--|---|--|--------------------|
| $C^{\rm E}_{\ 11} \ C^{\rm D}_{\ 11}$ | 64.00 [116] 64.80 [23] 66.40 [24] 64.13 [23] | 62.90 [109] 69.90 [114] 56.15 [115] | 68.1(1) 69.3(1) |
| C ^E ₁₂ | 22.00 [116] 21.30 [24] | 25.50 [109] 8.40 [114] 12.03 [115] | 25.1(1) |
| C^{E}_{13} | 32.00 [116] | 25.70 [109] 4.10 [114] 19.39 [115] | _ |
| $ C^{\mathrm{E}}_{14} $ | 2.00 [116] 11.70 [23] 2.20 [24] | 0.60 [109] 15.6 [114] 0 [115] | ≈0 |
| C^{E}_{33} | 118.0 [116] 116.0 [24] | 116.80 [109] 91.60 [114] 99.05 [115] | 118.8(2) |
| $C^{\mathrm{E}}_{}44}$ | 37.00 [116] 37.84 [23] 26.80 [24] | 35.00 [109] 38.40 [114] 39.99 [115] | 38.6(1) |
| $C^{\mathrm{E}}_{66} \\ C^{\mathrm{D}}_{66}$ | 21.00 [116] 21.10 [23] 22.53 [24] 24.90–25.14 [23] | 18.70 [109] 30.70 [114] 22.06 [115] | 21.5(1) 22.7(1) |

Table 3. Piezoelectric strain coefficient d_{11} of α -GeO₂ and α -SiO₂ (given for comparison).

| Material | | α-GeO ₂ | | α-SiO ₂ |
|---------------------------------|---------------------|--------------------|--------------|---------------------|
| Piezoelectric constant | Hydrothermal-Growth | Predicted values | Flux-Growth | Hydrothermal-Growth |
| | 4.04 [23] | 7.43 [109] | | |
| $d_{11} (10^{-12} \text{ C/N})$ | 4.10 [24] | 4.30 [117] | 5.7(2) [109] | 2.31 [17] |
| | 8.7-9.4 [29] | 6.00 [117] | | |

The flux-grown α -GeO₂ exhibits a d_{11} piezoelectric strain constant more than twice that of α -quartz SiO₂, Table 3, confirming the improvement of the piezoelectric properties with the structural distortion in the α -quartz analogues [17,117].

5.2. High Temperature

5.2.1. Flux-Grown α-GaPO₄

High-resolution Brillouin spectroscopy studies concerning the C_U elastic stiffness constant evolution with temperature of high-temperature flux-grown α -GaPO₄ single-crystals are reported in the literature [90,103]. Since for Brillouin scattering experiments, the samples have not to be coated with metal layers (electrodes), measurements at high temperatures are not affected by foreign chemical diffusion or bad electrical signal transmission.

Plates of X (100) and Z (001) simple orientations produced from α -GaPO₄ single crystals obtained by slow cooling from 950 to 600 °C in $X_2O:3MoO_3$ fluxes (X = Li, K) were used to follow the thermal evolution of C^D_{11} and C^E_{33} elastic constant [90]. C^D_{11} was measured from room temperature up to 1000 °C to follow the α -quartz/ β -cristobalite phase transition found close to 970 °C, in good agreement with other thermal studies [30,41,42,44,46,89,101]. C^E_{33} presented a monotonous decrease with temperature up to 850 °C while C^D_{11} presented a slight variation over the 20–500 °C temperature range followed by a stronger variation when approaching the phase transition temperature at 970 °C [90].

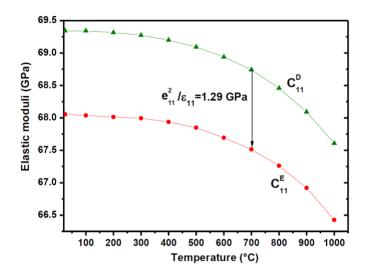
For the Brillouin scattering measurements using the backscattering geometry undertaken on a α -GaPO₄ crystal shaped as a cube with *X*-, *Y*- and *Z*-faces [103], the C^E_{11} , C^E_{33} , C^E_{44} and C^E_{12} constants showed monotonic elastic softening upon heating up to 850 °C while C^E_{66} and C^E_{14} showed continuous stiffening. Most of the C_{IJ} elastic stiffness constant derives by only a few percent upon heating, while C^E_{14} increases by about 50% and C^E_{12} decreases by about 35%. The first-order temperature coefficients $T^{(1)}_{CIJ}$ were negative for C^E_{11} , C^E_{33} , C^E_{44} and C^E_{12} elastic constants [103] in perfect agreement with previous reports on hydrothermally-grown α -GaPO₄ crystals [67,110,112,118].

5.2.2. Flux-Grown α-GeO₂

Different sets of C_{IJ} elastic stiffness constants at constant electric displacement and ambient conditions have been reported on flux- and hydrothermally-grown α -GeO₂ [23,24,109,116]. However, only one set of elastic data concerning their thermal evolution is reported in the literature [109]. It concerns a high resolution Brillouin scattering study of MoO₃-based flux-grown α -GeO₂ single crystals from room temperature up to 1000 °C. Any accident as discontinuity or phase transition was registered on the thermal evolution of the elastic constants in agreement with the conservation of the α -quartz-like structure of flux-grown GeO₂ material up to melting (1116 °C) [93,94]. The C_{11} , C_{33} , and C_{12} elastic constants show monotonic elastic softening upon heating while C_{44} and C_{66} show continuous stiffening.

The thermal evolution of the e_{11}^2/ϵ_{11} piezoelectric term, deduced from the difference between C_{11}^D and C_{11}^E , was found stable up to 1000 °C, Figure 4 [93,109] meaning that MoO₃-based flux-grown α -GeO₂ crystals exhibit piezoelectric activity even at very high temperature.

Figure 4. Thermal dependence of pure and piezoelectrically stiffened C_{11} elastic constant of flux-grown α-GeO₂. Errors are smaller than the symbol size.



6. Conclusions

In the family of α -quartz isotypes, high quality α -GaPO₄ and α -GeO₂ single crystals with a larger thermal stability and higher piezoelectric properties than α -quartz can potentially be used as piezoelectric materials for high temperatures applications. However if these single crystals are grown via hydrothermal-based methods, the presence of significant numbers of hydroxyl groups and structural defects in their network lead to a degradation of their physical properties and even to a phase transformation at relatively low temperature for α -GeO₂.

The main structural and chemical results obtained from infrared studies (OH-content), thermal behavior analysis and Brillouin scattering experiments on MoO₃-based flux-grown α -GaPO₄ and α -GeO₂ single-crystals have demonstrated OH-free high quality piezoelectric crystals and an improved thermal stability.

These results confirm on one hand the high potential of flux-grown α -GaPO₄ and α -GeO₂ single crystals as piezoelectric materials for high temperature applications and on the other hand a very powerful method of high temperature flux melted technique to grow high quality α -quartz isotype single crystals.

Moreover, as iso-structural to α -quartz, most of the device designs developed for α -quartz single crystal could be applied for α -GeO₂ with minor adaptations. The potential of α -GeO₂ single crystal for the realization of piezoelectric devices is also confirmed as its d_{11} piezoelectric constant at ambient temperature is found to be more than twice that of α -quartz. The piezoelectric property of α -GeO₂ is still conserved at high temperature as a significant piezoelectric contribution to C_{11} still exists at 1000 °C.

Author Contributions

Pascale Armand, Adrien Lignie and Marion Beaurain performed crystal growths of flux-grown piezoelectric materials, characterizations and data analysis. All author discussed the results and Pascale Armand wrote the paper.

Conflicts of Interest

The authors declare no conflict of interest.

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