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# Computer Simulations of a Twist Bend Nematic ( $N_{TB}$ ): A Coarse-Grained Simulation of the Phase Behaviour of the Liquid Crystal Dimer CB7CB

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**Abstract:** In recent years, a number of achiral liquid crystal dimer molecules have been shown to exhibit nematic–nematic phase transitions. The lower temperature phase has been identified as the  $N_{TB}$  phase, which demonstrates emergent chirality in the spontaneous formation of a heliconical structure. Recent fully atomistic simulations of the molecule CB7CB (1,7-bis-4-(4'-cyanobiphenyl) heptane), a dimer with an odd number of carbon spacers between the mesogenic parts of the molecule, have captured the  $N_{TB}$ –N–I phase sequence, providing a picture of the order at a molecular level. In this paper, we use atomistic simulations of CB7CB to develop a coarse-grained model using systematic coarse graining in the  $N_{TB}$  phase. We use both force matching (in the form of the MS-CG method) and iterative Boltzmann inversion (IBI) methodologies. Both techniques capture the heliconical order within the  $N_{TB}$  phase. Moreover, the model developed via force matching is shown to provide an excellent representation of the atomistic simulation reference model and, remarkably, demonstrates good transferability across temperatures, allowing the  $N_{TB}$ –N and N–I phase transitions to be simulated. We also compare results with those of a Martini 3-based coarse-grained model.

**Keywords:** liquid crystals; twist-bend nematic; coarse-graining; dimer; molecular dynamics



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## 1. Introduction

For almost 100 years, the two known nematic liquid crystal phases exhibited by low molecular weight mesogens were the standard uniaxial (apolar) nematic, N, phase and its chiral variant: the chiral nematic (or cholesteric)  $N^*$  phase [1]. However, theory has suggested the possibility that other nematics can exist, including biaxial nematics [2], polar nematics [3] and nematics in which the direction of the director vector varies spatially [4]. Many of these have now been discovered in low molecular weight nematics [5–9], in addition to polymers [10], lyotropic systems [11] and colloidal liquid crystals [12–14].

A particularly interesting member of the “new nematic family” is the twist–bend nematic phase [15], which was independently predicted by Meyer and Dozov [4,16]. This phase exhibits a nanoscale helicoidal structure with a repeat length of just a few nanometres (i.e., a few molecules) [4]. It was first seen in a liquid crystal dimer, 1,7-bis-4-(4'-cyanobiphenyl) heptane, CB7CB [8,15,17,18], where a flexible tail (of odd parity) links two rigid mesogenic moieties, and has been seen in other dimeric mesogens [19–28], bent-core species [29,30] and supramolecular hydrogen-bonded systems [9,31–33]. Possible future applications have been reported in areas such as optical devices [34–37], gels [38], photoswitchable adhesives [39] and photoalignment technology [40].

Molecular simulation provides a very effective method of studying liquid crystal phases [1,41–44]. Insights into liquid crystalline structure have arisen from simulation studies at a range of different levels, from all-atom molecular dynamics studies [45–50] to a variety of simpler coarse-grained studies in which some of the details of the molecular structure are sacrificed to concentrate on the role played by factors such as anisotropic shape or the anisotropy of attractive interactions [51–57]. At a molecular level, simulations

can provide information about the local structure of a phase, including the role of dipole correlation [58,59] and microphase separation [41] and the importance of molecular shape on molecular packing [42]. Moreover, simulations can provide insights into structural properties of liquid crystals, such as order parameters [60,61], elastic constants [62] and flexoelectric coefficients [63,64], together with dynamical quantities such as diffusion [65] and rotational viscosity [66].

Recently, several particle-based simulation studies have targeted the twist–bend nematic phase. Simulations show that the phase can arise simply from the efficient packing of (crescent-shaped) hard particles [67]. However, within studies that have attempted to reproduce the  $N_{TB}$  via atomistic studies [18,68], it becomes clear that the phase itself displays a remarkably subtle ordering of molecules. Essentially, molecular order can be influenced subtly by changes in dipoles, by the preferred angle of bend within an LC dimer (or LC oligomer), by the presence of certain functional groups or by changing molecular flexibility, as borne out by experimental studies of the relationship between transition temperature and  $N_{TB}$  stability [69–71].

Valuable insight from colloidal simulations reveals that for hard-bent spherocylinders, biaxial, twist–bend and splay–bend nematic phases must all compete with stable smectic phases. However, polydispersity in particle length or curvature can destabilize smectic ordering, allowing these nematic phases to be seen. It is interesting to note that in thermotropic low molecular weight organic molecules, molecular flexibility provides a mechanism to introduce polydispersity of shape with the same role of raising the relative free energy of competing smectic phases, allowing unusual nematics to be seen [68].

In a recent study [68], we carried out extensive all-atom molecular dynamics simulations of the bent dimer CB7CB and characterised its properties as a function of temperature. Despite enormous advances in the availability of computer time and parallel simulation algorithms to exploit this, all-atom models remain extremely computationally expensive and at the boundary of what is currently possible from a simulation point of view. However, in the current paper, we attempt to develop systematic coarse-grained models for CB7CB that capture the molecular ordering and phase behaviour seen in atomistic simulation studies. The aims of such models are to allow the subtle details of molecular order to be captured (i.e., molecular flexibility, dipole-dipole correlation, packing) but in a model that is orders of magnitude faster to simulate than an all-atom simulation.

Models such as the ones developed here open up new possibilities for studies of large systems sizes and long simulation times, far beyond what is currently possible at the all-atom level.

## 2. Computational

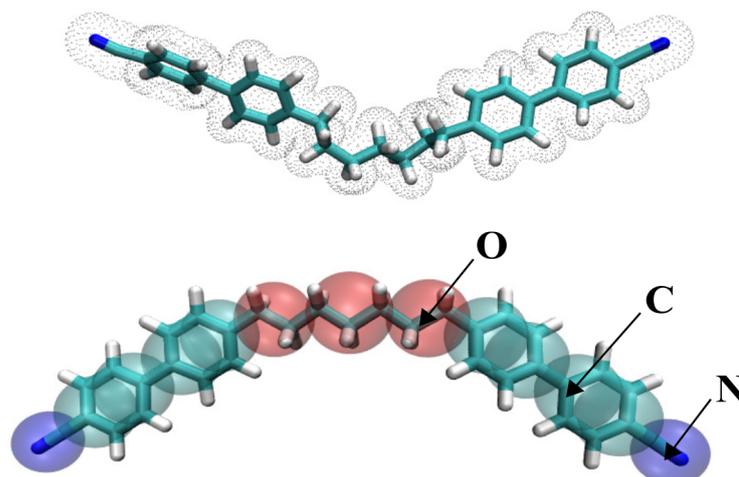
### 2.1. Choice of Methodology and CG Mapping

The coarse-grained molecular simulation models developed in this work are based on the structure of an underlying all-atom model for CB7CB, provided by our earlier simulation work [68]. In brief, our earlier calculations used a fully atomistic model composed of 512 molecules based on the GAFF force field. Simulations were carried out in the isobaric–isothermal ensemble at a range of temperatures and 1 bar of pressure in a simulation box with an approximate aspect ratio of 1:1:2, which employed anisotropic box scaling. This model demonstrated good agreement with experiment in relation to the structure of the  $N_{TB}$  phase of CB7CB, though with transition temperatures that are slightly higher than the experiment (we know from the earlier work of Boyd and co-workers [47,48,72,73] that the GAFF force field provides a good foundation for simulating liquid crystal phases but that this force field leads to molecules that are a little stiff and hence to molecules that are slightly too rigid and therefore phases that exhibit transition temperatures that are slightly too high in comparison with experiment). The atomistic reference phase transitions are ( $T_{N_{TB}-N} = 435$  K and  $T_{N-I} = 460$  K) in comparison to experiment ( $T_{N_{TB}-N} = 376$  K and  $T_{N-I} = 389$  K). However, the good agreement between experimental and atomistic

simulation in terms of pitch and average conical tilt angles provides an excellent starting point for helping to develop a coarse-grained model of CB7CB.

For the current work, we used snapshot configurations from the  $N_{TB}$  phase at a simulated temperature of 370 K (1 bar pressure) over a 1  $\mu$ s simulation run (the longest fully atomistic study of a liquid crystal reported to date) to provide reference data for atomistic radial distribution functions and forces. For our reference system, the heliconical nanostructure of the  $N_{TB}$  phase was formed spontaneously on cooling from a nematic phase, and the change in local order of CB7CB was seen reversibly in crossing and recrossing the transition by raising/lowering the temperature.

Figure 1 shows the mapping scheme used for this study. The choice of a 3:1 mapping scheme leads to a 13-site coarse-grained (CG) model and 6 pairwise interactions to fit. Within the 3:1 scheme, the combination of separate N–N, N–C and C–C interactions should be sufficient to capture the local packing structure of mesogenic units. We note that a slightly more fine-grained 2:1 mapping scheme would result in a need for additional bead types. While this may lead to greater accuracy, it would introduce greater complexity in the fitting procedures at a far greater computational cost. In particular, for the iterative Boltzmann integration (IBI) method (discussed below), a scenario with more than 6 types of beads (i.e., >15 interaction pair-types) becomes extremely difficult to fit.



**Figure 1.** (Top): Snapshot conformation of a single CB7CB molecule showing the 3D atomistic structure. (Bottom): Coarse-grained mapping scheme showing the three types of coarse-grained beads used in this work (termed N, C, O). Each bead incorporates heavy atoms and attached hydrogens.

## 2.2. Coarse-Grained Methods Used: IBI, FM, Martini 3

Coarse-grained models for the IBI scheme were developed using the VOTCA-CSG package, version 1.4.1 [74–76] to determine nonbonded parameters. Force matching (FM) was carried out using the MS-CG method within a hybrid force matching (hFM) scheme [75,77]. In this method, force matching is carried out to determine the nonbonded potentials allowing for forces arising from explicit bonded potentials. This method was carried out using the BOCS (bottom-up open-source coarse-graining software) package [78]. For both methods, a linear pressure correction was used. This was done so that a pressure of  $\sim 1$  bar could be obtained for simulations at the state point at which coarse-graining was carried out. Potentials were tailored to go smoothly to zero at a cut-off of 1.2 nm. For this work, IBI and FM used the atomistic reference coordinates from the  $N_{TB}$  phase at 370 K, employing atomistic trajectories containing 1000 frames. In the calculations, we have followed the general approaches outlined by Potter et al. [79] for simple mixtures, which have been used recently for chromonic liquid crystals [80,81].

CB7CB was coarse-grained using Martini 3 nonbonded potentials. The Martini 3 force field consists of a substantial library of Lennard–Jones bead types. This study was completed before the formal publication of Martini 3 [82], and we used the parameters released

in the open beta version of the Martini 3 force field [83]. We use the bead types N = TC6, C = SC4, O = TC1. For Martini 3 simulations, the standard Martini cut-off of 1.1 nm was used for all nonbonded interactions [84].

As an initial test of the Martini parameters, we carried out simulations of cyanobiphenyl and heptane as pure compounds and as a 50:50 mixture and compared these with results from all-atom force field studies using GAFF. Here, we used a standard constant- $NpT$  simulation approach [47,48,85] to calculate  $\Delta_{\text{vap}}H$  from bulk and single molecule simulations and obtained  $\Delta_{\text{sol}}G$  from decoupling the intermolecular interactions of a single molecule from the solvent using the Bennett acceptance ratio (BAR) method [86] following the simulation approach of Yu and co-workers [49,80]. The results (shown in Table 1) indicate that Martini 3 does quite a good job in terms of correctly capturing the thermodynamics of these molecular fragments that control the interactions of CB7CB molecules. However, we note that for bulk liquids, Martini 3 models tend to pack better than real molecules, and this leads to an over-estimate of liquid densities [87]. We note in passing that preliminary work in our laboratory suggests that it may be possible to systematically optimise Martini by fitting the heat of vaporization at 298.15 K together with the density at a range of state points, using the ForceBalance [88] software to carry out systematic optimisation. However, this approach is very computationally intensive as it requires many separate CG simulations for each optimisation step. Alternatively, progress is currently being made using the SAFT- $\gamma$  equation of state approaches to carry out top-down coarse graining [89–92], and this method has been applied recently to liquid crystals [93].

**Table 1.** Enthalpies of vaporization ( $\Delta_{\text{vap}}H/\text{kJ mol}^{-1}$ ) and free energies of solvation  $\Delta_{\text{sol}}G/\text{kJ mol}^{-1}$ ) for all-atom GAFF (AA) and Martini 3 CG models of cyanobiphenyl/heptane mixtures.

Property	Experiment	AA Model	Martini 3
$\Delta_{\text{vap}}H$ (cyanobiphenyl)	55.5 [94]	$52.3 \pm 0.4$	$47.5 \pm 0.4$
$\Delta_{\text{vap}}H$ (heptane)	36.4 [95]	$38.5 \pm 0.5$	$38.7 \pm 0.4$
$\Delta_{\text{sol}}G$ (cyanobiphenyl, $x_{\text{CB}} = 1$ )	–	$-38.4 \pm 0.10$	$-31.91 \pm 0.06$
$\Delta_{\text{sol}}G$ (heptane, $x_{\text{heptane}} = 1$ )	–	$-19.81 \pm 0.14$	$-27.4 \pm 0.2$
$\Delta_{\text{sol}}G$ (cyanobiphenyl, $x_{\text{CB}} = 0.5$ )	–	$-32.9 \pm 0.3$	$-30.49 \pm 0.07$
$\Delta_{\text{sol}}G$ (heptane, $x_{\text{CB}} = 0.5$ )	–	$-22.5 \pm 0.2$	$-26.91 \pm 0.08$

All coarse-grained simulations used the GROMACS molecular dynamics package, with simulations carried out in the canonical ensemble. For the initial model development, 512 molecules were used. This was later scaled up to 4096 molecules for larger simulations of the FM model. Simulations used a time step of 2 fs for initial setup and equilibration and 5 fs for production runs and used a stochastic dynamics integrator with a time constant of 2 ps.

The CG interactions potential can be written as

$$E^{\text{CG}} = \sum_{\text{bonds}} K_r(r - r_{\text{eq}})^2 + \sum_{\text{angles}} K_\theta(\theta - \theta_{\text{eq}})^2 + \sum_{\text{dihedrals}} k_\phi(1 + \cos(n_d\phi - \phi_d)) + \sum_{i>j}^N E_{ij}, \quad (1)$$

where  $r_{\text{eq}}$ ,  $\theta_{\text{eq}}$  are the natural bond lengths and angles, respectively,  $K_r$ ,  $K_\theta$  and  $K_\phi$  are the bond, angle and torsional force constants, respectively, and  $\phi$ ,  $n_d$  and  $\phi_d$  represent a dihedral angle, its periodicity and the phase angle associated with it, respectively.  $E_{ij}$  are nonbonded energies for particles  $i$  and  $j$ , which in the case of the IBI and FM methods are numerical potentials and in the case of Martini 3 are Lennard–Jones 12:6 potentials. In this work, we make the approximation that the influence of partial charges across the atomistic molecule can readily be coarse-grained into  $E_{ij}$ .

### 2.3. Cost of Computational Models

A few comments on the relative computational costs of the three models are appropriate. In terms of setup costs, the Martini 3 model has a significant advantage in being an “off-the-shelf” approach that usually does not require an underlying atomistic reference. However, in the current work, it was still relatively important to do some atomistic work (see below) to ensure that the combination of inter and intramolecular potentials gives a good representation of molecular shape. The IBI and FM approaches both require the significant expense of an underlying reference atomistic model for at least one state point. However, the IBI approach requires significantly more computational work than FM in the coarse-graining procedure (in order to fit the multiple partial radial distribution functions for all pair interactions required in IBI). In terms of production runs, the three models have very similar computational costs as these costs are primarily related to the time step used and the number of coarse-grained sites. Here, the stiffness of the intramolecular potentials places a limit on the length of the time step.

### 2.4. Bonded Potentials

For systems such as this, it is possible to obtain numerical bonded potentials directly from the atomistic calculations via the Boltzmann inversion of bond, bond angle and dihedral angle distributions for the coarse-grained sites, or to use numerical potentials from IBI. However, in the case of coarse-graining in the  $N_{TB}$  phase, these potentials would be likely to inherit a chiral bias from the symmetry of the phase. Hence, we use a simpler approach. In the spirit of the Martini model, we employ simple harmonic potentials for bonded parameters, determining optimum parameters by reference to the bond, angle and dihedral distribution functions, together with further testing on single molecule simulations. For the latter, we wanted to capture a good representation of the overall molecular shape in the absence of surrounding molecules, as represented by the correct bend angle between mesogenic moieties within the molecule and the angle distribution function (which is dependent on the conformational distribution). Simulations at 370 K, using the final parameters in Tables 2–4, reproduce the preferred atomistic bend angle of  $127^\circ$  and show no preference for right-handed or left-handed conformations in single molecule calculations. The latter was particularly important as we want chirality in the model to be able to arise without any biasing potentials and for the model to be able to show both achiral (nematic) and chiral ( $N_{TB}$ ) phases. However, we note that the coarse-grained model is slightly stiffer than its atomistic counterpart, exhibiting a full width at half maximum (FWHM) for the bend angle distribution of  $15^\circ$  (see ESI Figure S2) in comparison to  $40^\circ$  for the atomistic model. Moreover, in the atomistic model [68], we detected a few highly bent “hairpin” conformations arising from multiple gauche conformations in the chain. These rare conformations are not captured in the coarse-grained model.

**Table 2.** Bond parameters for CG models of CB7CB.

Atoms	Length/nm	$K_r/\text{kJ mol}^{-1} \text{nm}^{-2}$
N–O	0.25	15,000
C–C	0.22	15,000
C–O	0.24	10,000
O–O	0.32	10,000

**Table 3.** Angle parameters for CG models of CB7CB.

Atoms	Angle/deg	$K_\theta/\text{kJ mol}^{-1} \text{deg}^{-2}$
N–C–C	180	5000
C–C–C	180	5000
C–C–O	170	750
C–O–O	170	750
O–O–O	170	750

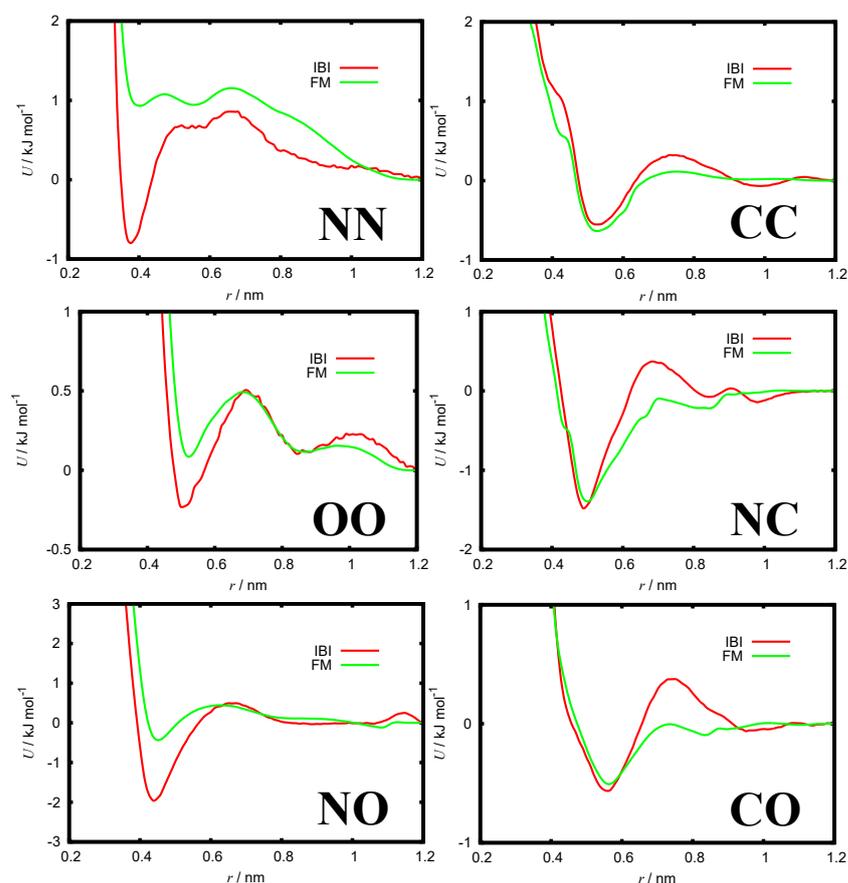
**Table 4.** Dihedral angle parameters for the CG models of CB7CB.

Atoms	Dihedral/deg	$K_\phi/\text{kJ mol}^{-1}$
C–C–O–O	180	25
C–O–O–O	180	25

### 3. Results and Discussion

#### 3.1. CG Potentials

Figure 2 shows the numerical coarse-grained potentials generated through the IBI and FM methodologies at 370 K. The six nonbonded potentials are very similar to those typically found through bottom-up coarse-graining techniques and in some cases include more than one well, reflecting the different correlations seen between coarse-grained sites in the  $N_{TB}$  phase. Both methods capture the structure of the phase (see below) but interestingly do so by a completely different combination of pair potentials. For example, the N–N potential for FM is totally repulsive, but the N–N for IBI gives a deep attractive well. However, here it is a combination of pair potentials that lead to the structure of the fluid, and we expect different combinations of potentials to be able to generate very similar pair statistics in this complex case [96].

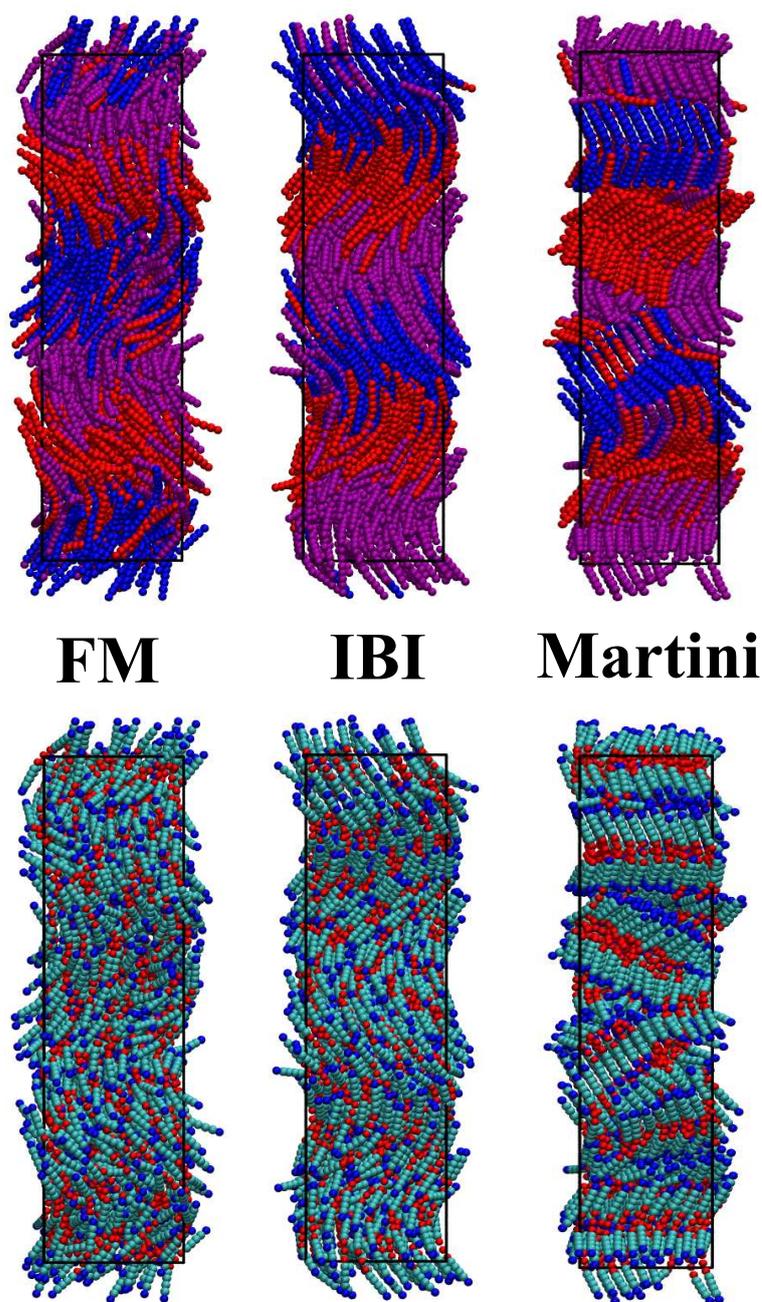
**Figure 2.** Coarse-grained nonbonded potentials obtained from IBI and FM for CB7CB at 370 K.

#### 3.2. Phase Behaviour of the Coarse-Grained Models at 370 K

The simulation snapshots for the three models are shown in Figure 3. In the top part of the figure, we colour molecules blue, purple or red. Here, colour-coding picks out variations in the azimuth angles ( $\phi$ ) calculated from the molecular long axes (i.e., the unit vector  $\hat{a}$ ) in the range:  $-180^\circ \leq \phi < -60^\circ$ ,  $-60^\circ \leq \phi < 60^\circ$  and  $60^\circ \leq \phi < 180^\circ$ . From visual inspection, we find both the IBI and FM models exhibit a  $N_{TB}$  phase, where the mesogens form a heliconical structure. As with the reference atomistic structure, the helix

found here is left-handed. However, we note that, experimentally, there are equal domains of opposite handedness. We might expect to be able to see either-handedness forming within a relatively small simulation box; however, the periodic boundary conditions and the relatively small size of the system (in comparison to experimental domains) guarantee the presence of only one chiral domain.

The Martini 3 simulations however form a crystalline smectic structure where there is a far stronger positional correlation between molecules than seen in either the IBI or FM models. This leads to the loss of smooth helicoidal ordering at 370 K, though interestingly, we still see a molecular tilt and a precession of the molecular tilt angle along the long axis of the simulation box.



**Figure 3.** Simulation snapshots of coarse-grained models of CB7CB at 370 K for the FM, IBI and Martini 3 CG models. (**Top**): Colouring by molecular orientation. (**Bottom**): Colouring according to bead type.

The molecular order at 370 K can be characterised by the order parameters  $\langle S_2 \rangle_a$  and  $\langle S_2 \rangle_b$  for the long axis vector  $\hat{a}$  and a second short-axis vector  $\hat{b}$ .  $\hat{b}$  is defined as the unit vector which bisects the two vectors between the terminal N sites and the centre of mass (COM) of the molecule. Values of  $S_2$  for a chosen axis vector can be determined from the diagonalisation of the order tensor

$$Q_{\alpha\beta} = \frac{1}{N} \sum_{i=1}^N \frac{1}{2} (3u_{i\alpha}u_{i\beta} - \delta_{\alpha\beta}), \quad (2)$$

where  $N$  is the number of molecules,  $\delta_{\alpha\beta}$  is the Kronecker delta and  $u_i$  is the axis vector  $\hat{a}$  or  $\hat{b}$  for molecule  $i$ . The largest eigenvalue provides the value of  $S_2$  and its associated eigenvector is the director,  $\hat{n}$ . While  $\langle S_2 \rangle_a$  is traditionally used to measure the degree of order in a normal nematic phase,  $\langle S_2 \rangle_b$  can measure the polar correlations or smectic-like ordering normal to the long axes.

Values of  $\langle S_2 \rangle_a$  and  $\langle S_2 \rangle_b$  are given in Table 5. The degree of order is slightly larger in the coarse-grained models than that seen atomistically. Most likely, this is due to a slightly lower degree of flexibility in the CG models. The FM model performs best, and  $\hat{b}$  for this model is very close to the value obtained in the all-atom model. The values of the pitch obtained in the simulation are also given in Table 5. The coarse-grained results are very close to all-atom results, which in turn are in good agreement with the experimentally measured pitch for CB7CB of  $\sim 8$  nm [8,18]. One should note that for the FM and IBI models, the pitch is pre-programmed inasmuch as the underlying structure of the phase is captured by the coarse-graining procedure at 370 K.

The average values of the conical tilt angle,  $\theta$ , are given in the final column of Table 5.  $\theta$  is calculated as the angle between the unit vector  $\hat{a}$  and the helical axis. The values we obtain are in good agreement with the atomistic values. Both atomistic and coarse-grained simulation results are slightly higher than the experimental value of  $\sim 25^\circ$  [18,97,98].

Although not capturing the  $N_{TB}$  phase, the Martini 3 model is interesting in the sense that it shows a molecular tilt, the direction of which precesses through the system despite the formation of smectic layers. From the Martini 3 snapshot structure in Figure 3, we note the similarity of this model to previous models of hard-bent spherocylinders. Chiappini et al. [99] have shown that, for the latter, the  $N_{TB}$  phase must compete with an array of smectic phases unless the smectics are destabilized by polydispersity in the particle length or by curvature. While the Martini 3 model contains attractive interactions in addition to repulsions, the molecular packing is still dominated by shape. When attractive interactions between molecules overcome some of the polydispersity of shape (arising from conformational flexibility), the system readily forms a smectic in preference to the  $N_{TB}$  phase. We note also that the local ordering in this phase has similarities to systems that show the segregation of mesogens into ferroelectric domains [45,100–103]. Moreover, the phase organisation observed here is similar to the twist-bend smectic phase reported in dimers with a terminal alkyl chain [104–106].

A key difference between the Martini 3 model and the two bottom-up models arises from the coarse-grained potentials themselves. Within Martini 3, the strongest like-like interactions are highly favourable, but cross-interactions are generally less favourable. This tends to promote microphase separation between similar regions of a molecule. However, stronger cross-interactions are able to destabilise smectic ordering and this, in turn, leads to greater stability of nematic phases. As discussed above, the destabilisation of a smectic phase can also be achieved via polydispersity of particle/molecule size, shape or curvature. For CB7CB molecules, conformational disorder provides this polydispersity, but for the Martini 3 model, some of this is lost in the formation of crystalline smectic structures.

Finally, it is interesting to discuss the role of the N–N interactions (corresponding to the cyano-group self-interaction) within the bottom-up models. In the FM model, this interaction is completely repulsive, and the N–C interaction is stronger than the C–C interaction. In adjacent molecules, this leads to either a preference for anti-parallel correlations of cyano groups or parallel shifted correlations (i.e., these potentials encode the local effects of dipole-

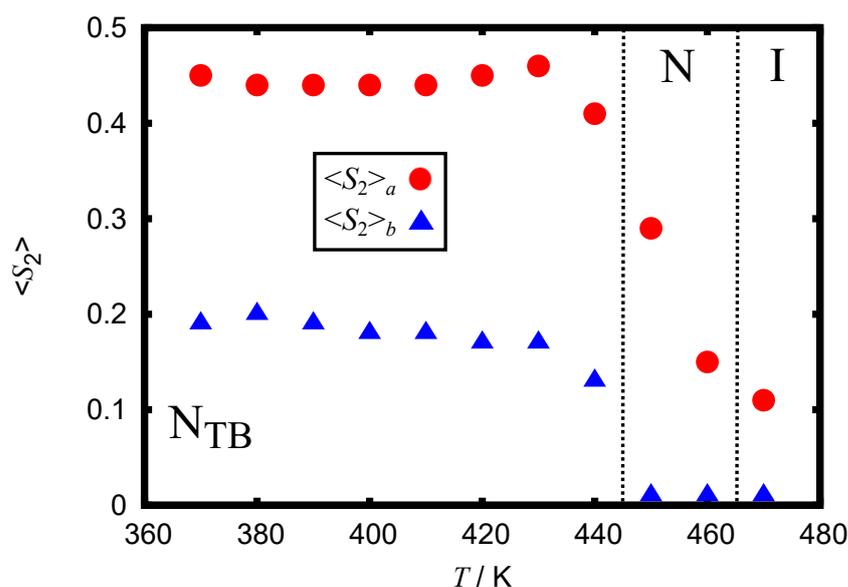
dipole interactions seen in the atomistic model). These interactions promote nematic ordering at the expense of smectics. While the IBI achieves a similar effect with different potentials, the stronger IBI N–N interaction leads to an overall stronger correlation between adjacent molecules and an orientational order that is slightly higher than seen in the FM model.

**Table 5.** The order parameters ( $\langle S_2 \rangle_a$  and  $\langle S_2 \rangle_b$ ), helical pitch ( $p$ ) and conical tilt angle ( $\theta$ ) for the AA and 3 CG models of CB7CB at 370 K.

Model	$\langle S_2 \rangle_a$	$\langle S_2 \rangle_b$	$p/\text{nm}$	$\theta/^\circ$
AA	0.36	0.16	8.35	29
FM	0.44	0.18	8.35	32
IBI	0.51	0.28	8.42	31
Martini 3	0.49	0.30	8.5	32

### 3.3. Further Simulation of the FM Model

From the results in Section 3.2, the FM model is judged to be the most representative CG model of the three studied for CB7CB. This model was selected for a heating sequence to produce a full phase diagram. Here, we test the transferability of the model to different temperatures to see whether it can capture the N or I phases and the phase transition temperatures. Starting from a configuration at 370 K, the system was heated at intervals of 10 K for 200 ns at each temperature within canonical ensemble simulation runs. The phase diagram produced is shown in Figure 4, in which the stability of the  $N_{TB}$ , N and I phases are indicated. The  $N_{TB}$ –N transition was defined as the temperature at which the helical order is lost and  $\langle S_2 \rangle_b \approx 0$ . Remarkably (given the limited temperature transferability often seen for such models), the FM CG model successfully captures the phase transitions in the N and I phases. The transition temperatures seen are  $T_{N_{TB}-N} \sim 445$  K, which is about 10 K higher than the all-atom model and  $T_{N-I} \sim 465$  K. Such good transferability to the N and I phases was not expected, particularly considering that the potentials were parametrised in the  $N_{TB}$  phase. However, this probably indicates that the local interactions between molecules within the two nematics phases are not dramatically different, and the CG potentials are sufficiently detailed to represent that structure well (see above). We note additionally that the similar probability distribution functions for the molecular bend angle measured from atomistic simulations in the two phases [68] might also be a contributing factor.



**Figure 4.** Temperature dependence of the orientational order parameter,  $\langle S_2 \rangle$ , for unit vectors  $\hat{a}$  and  $\hat{b}$  for the FM CG model.

Within the  $N_{TB}$  phase itself, the value of  $\langle S_2 \rangle_a$  does not increase with temperature, contrary to the results from atomistic simulation [68]. For the FM model, this is almost certainly down to coarse-graining at just one temperature and therefore does not capture the temperature dependence of the conformations seen in the all-atom chains. In the latter, conformations change subtly with temperature, and this leads to subtle changes of molecular tilt and a decrease in  $\langle S_2 \rangle_a$  as the  $N_{TB}$  is cooled.

On cooling from the isotropic phase with the FM model (using 200 ns simulations and 20 K temperature intervals and a longer 400 ns in cooling from 450 K to 430 K), we are able to spontaneously recover the N and  $N_{TB}$  phases. The same transition temperatures are obtained on cooling as those seen on heating (see ESI Figure S1 for the phase diagram and simulation snapshots).

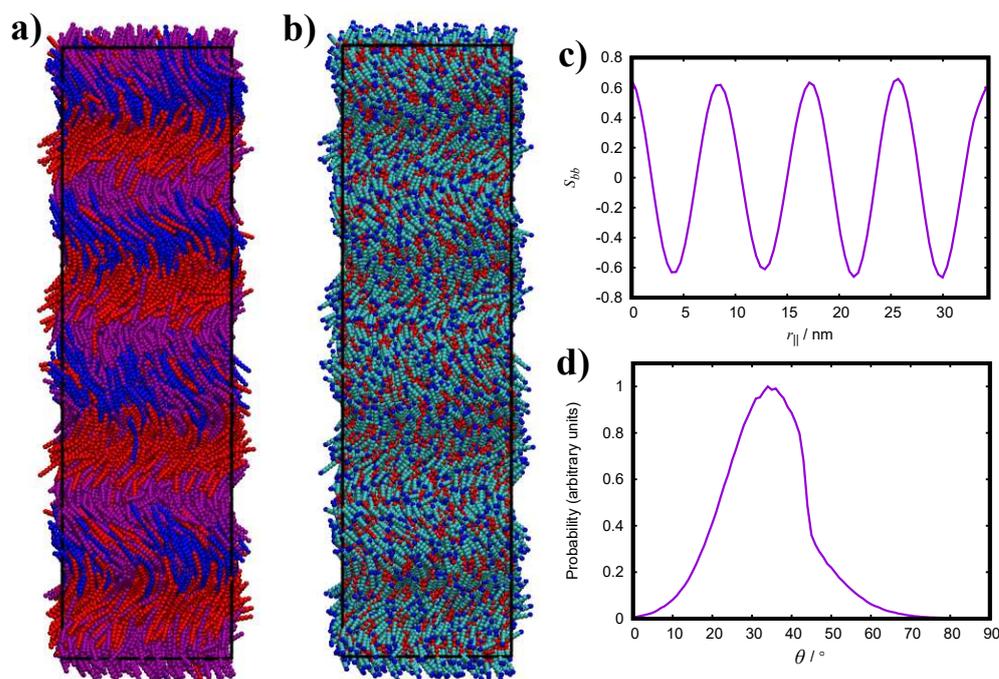
One motivation for coarse-graining is to facilitate the simulation of systems beyond the length and time scales accessible to all-atom molecular dynamics. This is particularly important for complex phases, such as the  $N_{TB}$  phase, where larger systems move the simulation away from a regime where small system sizes could induce system size effects or (worse) introduce artificial order into a system. Accordingly, we set up larger simulations of 4096 molecules for the FM model at 470 K by replicating our 512 molecules system in each direction and then carrying out long runs to ensure molecular order within the system was fully equilibrated for this large system size. Figure 5 shows the results from a 200 ns simulation of this system. The calculated values of  $\langle S_2 \rangle_a = 0.44$  and  $\langle S_2 \rangle_b = 0.18$  are in full agreement with the smaller system, and the helicoidal ordering, within a phase of flexible molecules showing translational disorder, is fully evident from the snapshot pictures in Figure 5a,b. For this system, we have calculated the orientational correlation function,  $S_{bb}$  (Figure 5c), as a function of intermolecular distance along the helical axis,  $r_{\parallel}$  [107].  $S_{bb}(r_{\parallel})$  describes the polar correlations through the system,

$$S_{bb}(r_{\parallel}) = \sum_i^N \sum_{j \neq i}^N c_{ij} (\hat{\mathbf{b}}_i \cdot \hat{\mathbf{b}}_j), \quad (3)$$

where

$$c_{ij} = \delta(|z_{ij}| - r_{\parallel}) / \sum_i^N \sum_{j \neq i}^N \delta(|z_{ij}| - r_{\parallel}), \quad (4)$$

where  $z_{ij}$  is the distance with respect to the  $z$  components for the COMs of molecules  $i$  and  $j$ , and  $N$  is the number of molecules. The results show a beautiful precession of  $\hat{\mathbf{b}}$  vectors about the helical axis through the system, as expected in the  $N_{TB}$  phase [67,108]. For this large system, we are able to capture four pitch lengths within the simulation box. Figure 5d shows the distribution of the conical tilt angle  $\theta$  averaged over all molecules. It shows quite clearly that there is a range of tilt angles for molecules, i.e., considerable disorder, as one might expect for a fluid phase. However, the average tilt of  $29^\circ$ , in combination with the helicoidal twist clearly demonstrates how the structure changes dramatically in comparison to a normal nematic phase.



**Figure 5.** Simulation snapshots from a 4096–molecule system of the FM CG model, with (a) orientational and (b) bead-type colourings. (c) Orientational correlation ( $S_{bb}$ ) as a function of the helical axis,  $r_{\parallel}$ , and the (d) distribution of the conical tilt angle,  $\theta$ .

#### 4. Conclusions

We have developed coarse-grained models for the mesogen CB7CB using bottom-up (FM and IBI) and top-down (Martini 3) models.

Using 13 CG sites and three CG bead types, both FM and IBI deliver good CG models of the  $N_{TB}$  phase. In both cases, the results are in good agreement with the reference all-atom model for CB7CB. Of these two techniques, FM gives a slightly better representation of the key structural quantities such as orientational order parameters. Further simulations of the FM model show that it demonstrates remarkably good transferability across temperatures, exhibiting reversible phase transitions to a normal nematic, N, and an isotropic phase on heating. Cooling runs spontaneously recapture the original  $N_{TB}$  phase. Large simulations of 4096 FM molecules in the  $N_{TB}$  phase show very similar results to a smaller system of 512 molecules but exhibit helicoidal ordering over 4 pitch lengths. In this phase, the molecules show a wide spread of conical tilt angles, demonstrating that considerable disorder occurs within the microstructure of a molecular  $N_{TB}$  phase.

The FM model shows quite good agreement with experimentally measured quantities for CB7CB. The helical pitch seen (8.35 nm) is very similar to the experimental value ( $\sim 8$  nm), and we see an average conical tilt of molecules within the simulated  $N_{TB}$  phase ( $\sim 32^\circ$ ), similar to that seen experimentally ( $\sim 25^\circ$ ). We note that the simulated transition temperatures predictions are a little high ( $T_{N_{TB}-N} \sim 440$  K and  $T_{N-I} \sim 465$  K) in comparison to experiment ( $T_{N_{TB}-N} = 376$  K and  $T_{N-I} = 389$  K). Here, the discrepancy is due to the underlying atomistic model ( $T_{N_{TB}-N} = 435$  K and  $T_{N-I} = 460$  K), which predicts slightly too stiff molecules, resulting in slightly higher density for liquid crystal phases. However, capturing the spontaneous formation of a helicoidal twist on cooling from the nematic represents a significant success for this coarse-grained model. The simulations show the significance of the average bent molecular shape in allowing the molecules to adopt a twisted superstructure, without significant changes in conformational distribution in comparison to the nematic phase. As we discussed above, the molecular flexibility is playing a significant role in destabilizing smectic phase formation and allowing the twist bend nematic structure to arise.

With the Martini 3 model, we were unable to capture the  $N_{TB}$  phase. Here, the system loses the disorder seen in the bottom-up models and forms crystalline smectic ordering while retaining the molecular tilt and the precession of the tilt angle. However, models of this type may well be valuable in studying a range of new smectic phases, such as the heliconical tilted smectic C ( $SmC_{TB}$ ) phase [104].

In future work, it would be interesting to explore further the intermolecular forces that give rise to  $N_{TB}$  behaviour. In many systems, subtle changes in phase behaviour are seen with small changes in chemical structure. A careful comparison study of the influence of intermolecular forces arising from different functional groups may well be informative in the molecular engineering of future compounds to either include (or exclude) the  $N_{TB}$  phase in a sequence of liquid crystal phases.

**Supplementary Materials:** The following supporting information can be downloaded at: <https://www.mdpi.com/article/10.3390/cryst13030502/s1>, Figure S1: Simulation snapshots for 512 molecules of the FM CG model, obtained via cooling from an isotropic liquid. Results show: (a) the initial I phase at 470 K, (b) the N phase at 450 K, (c) the  $N_{TB}$  phase at 370 K, (d) the average orientational order parameter,  $\langle S_2 \rangle$ , for unit vectors  $\hat{a}$  and  $\hat{b}$  obtained from a cooling run starting in the I phase between the temperatures 470–370 K., Figure S2: Bend angle distribution function calculated from the angle between rigid mesogenic units. Simulations taken from the bulk phase simulation of the FM CG model at 370 K.

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**Data Availability Statement:** ESI is included covering order parameter data from cooling runs of the FM model, snapshots of the FM  $N_{TB}$ , N and I phases and the bend angle distribution function.

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