

Ab Initio Theoretical Study of DyScO₃ at High Pressure

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Abstract: DyScO₃ is a member of a family of compounds (the rare-earth scandates) with exceptional properties and prospective applications in key technological areas. In this paper, we study theoretically the behavior of DyScO₃ perovskite under pressures up to about 65 GPa, including its structural and vibrational properties (with an analysis of the Raman and infrared activity), elastic response, and stability. We have worked within the ab initio framework of the density functional theory, using projector-augmented wave potentials and a generalized gradient approximation form to the exchange-correlation functional, including dispersive corrections. We compare our results with existing theoretical and experimental published data and extend the range of previous studies. We also propose a candidate high-pressure phase for this material.

Keywords: DyScO₃; high pressure; density functional theory; stability

1. Introduction

On account of their dielectric and optical properties, as well as their optimal chemical stability, the rare-earth (RE) scandates with formula REScO₃, and among them DyScO₃, have been researched for the past two decades with a view mainly to replace SiO₂ in next-generation MOSFETs [1–3]. Other usages of RE-scandates include their employment as optimal substrates for the epitaxial growth and fine-tuning of high-quality perovskite-type thin films, as well as their application in the terahertz range when embedded in perovskite heterostructures [4–10]. The enticing properties and promising scope of applications of DyScO₃ suggest extending the fundamental research on this material to include as broad a range of pressures as possible, and very recently Bura and co-workers have experimentally undertaken this goal in the range from ambient conditions up to 40 GPa [11]. In the present work we aim to provide an ab initio theoretical study, performed within the framework of the density functional theory (DFT), of the so far only observed phase of DyScO₃, in the range up to 65 GPa, with a focus on its structural, elastic, and vibrational response, to compare with the experimental results of Bura et al. [11] as well as with other previously existing works and to provide a bound to its local stability. Our results indicate the existence of local instabilities, both elastic and dynamic, above 60–65 GPa. We also provide a possible candidate high-pressure phase, a well-known post-perovskite structure. We hope that these results will stimulate future experiments which, combined with theoretical studies, may lead to a better understanding of the properties of this technologically relevant, emerging material under high pressure.

2. Theoretical Method and Details of the Calculations

All the calculations were performed within the ab initio framework of the density functional theory (DFT) [12,13], using a projector augmented wave scheme (PAW) [14,15] and the VASP computational package [16–19]. The outermost, valence electrons of each species were explicitly considered in the calculations whereas their innermost, closed-shell core electrons were treated at the PAW level. For dysprosium we have used the recommended VASP Dy_3 PAW potential, which represents Dy with partially frozen 4f



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electrons and formal valency 3, a standard way to cope with localized f electrons in DFT calculations. The cutoff in the kinetic energy of the basis's plane waves was 540 eV and the required Brillouin-zone integrations were performed with a $4 \times 4 \times 4$ grid [20]. Several levels of approximations for the exchange-correlation functional were considered in our preliminary calculations, with optimal results obtained for the PBE generalized gradient approximation form [21] including dispersive corrections within the so-called Grimme's D3 scheme [22], which was accordingly the level of approximation adopted for the production of all the results shown in this paper. The material was structurally relaxed at different volumes, sampling the region between negative pressures just below 0 GPa to about 60 GPa. Hydrostatic conditions were assumed in all cases and the criteria for stopping the structural relaxations was set to stress anisotropies less than 0.1 GPa and residual forces on atoms less than 5 meV/Å.

Further to the total energy calculations just described, which provide information on the pressure dependence of the structural properties of DyScO₃, we also performed a general study of its dynamical and elastic properties and their variation with pressure, and draw conclusions on the stability of the system based on them, as described in Sections 3.2 and 3.3. For the study of the lattice vibrations, we used the small-displacement method as implemented in the phonopy package [23], employing a $2 \times 2 \times 2$ supercells for sampling the phonons across the whole Brillouin zone. At the zone center, the symmetry of the calculated modes was analyzed to determine their possible spectroscopical activity, Raman or infra-red (IR). For the calculation of the elastic constants, we used the Le Page stress-strain method as implemented in the VASP code [24].

3. Results and Discussion

3.1. Crystallographic Description of DyScO₃ Perovskite and Evolution of Its Structural Properties with Pressure

Like many other ABO₃ oxide compounds, DyScO₃ crystallizes in the perovskite-type structure. The ideal cubic perovskite structure, space group (SG) *Pm-3m* No 221, is characterized by corner-sharing BO₆ octahedra and 12-fold coordinated A sites. For non-ideal or distorted perovskites, octahedron tilting lowers the space group symmetry of the structure [25]. For DyScO₃ that means an orthorhombic phase with SG *Pbnm* No 62 and four formula units on the conventional unit cell (see Figure 1).

Our calculated lattice parameters *a*, *b*, *c*, and unit cell volume *V* at 0 GPa are in good agreement with the experimental results, as shown in Table 1. The calculated variation with the pressure of the lattice parameters (Figure 2) also reproduces the experimental results within a wide range of pressures up to approximately 10–12 GPa, although for larger pressures the calculated values of *a* and *c* are somewhat lower than the experimental values while *b* still follows the experimental trend. These deviations may be related to the loss of hydrostaticity above 10–12 GPa [26] in the methanol/ethanol pressure-transmitting medium used in the experiments [11] but more experimental and theoretical research would be needed to shed light on this point. The calculated linear axial compressibilities are $1.469 \times 10^{-3} \text{ GPa}^{-1}$ for *a*, $0.731 \times 10^{-3} \text{ GPa}^{-1}$ for *b*, and $1.411 \times 10^{-3} \text{ GPa}^{-1}$ for *c*, and the value for the bulk modulus *B*₀ of 164.7 GPa obtained by fitting the calculated energy-volume curve to a third-order Birch-Murnaghan equation of state [27] is a bit below the experimental one in Ref. [11] (see Table 2) whereas it is in good agreement with the calculated value of 165 GPa retrieved from the Materials Project database [28,29].

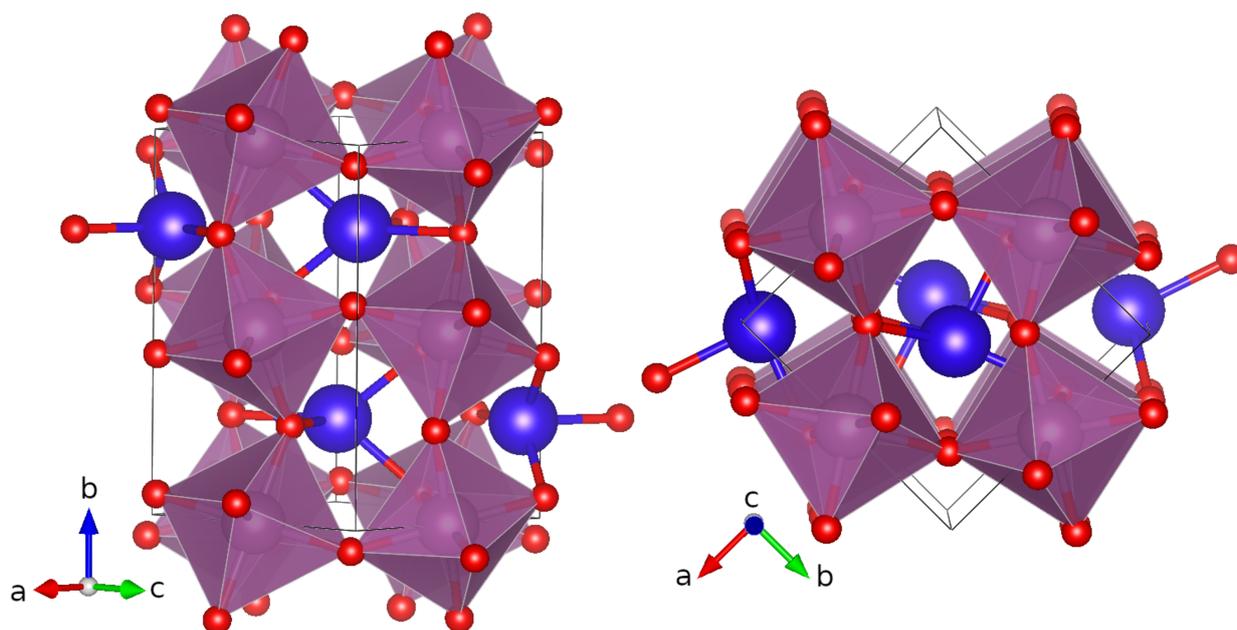


Figure 1. The structure of DyScO₃ is of the orthorhombic perovskite type. In this plot the ScO₆ octahedra of the structure are shown in magenta, with oxygen sites (in red) at their corners and Sc sites at the centers. The Dy cation sites are shown in blue.

Table 1. Calculated and experimental lattice parameters and unit cell volume at 0 GPa. Also provided are the calculated values at 60 GPa, for reference.

	This Work 0 GPa	Experiment Ref. [25]	Experiment Ref. [30]	This Work 60 GPa
a (Å)	5.421	5.443 (2)	5.442417 (54)	4.944
b (Å)	5.735	5.717 (2)	5.719357 (52)	5.484
c (Å)	7.935	7.901 (2)	7.904326 (98)	7.266
V (Å ³)	246.68	245.9 (1)	246.04	197.00

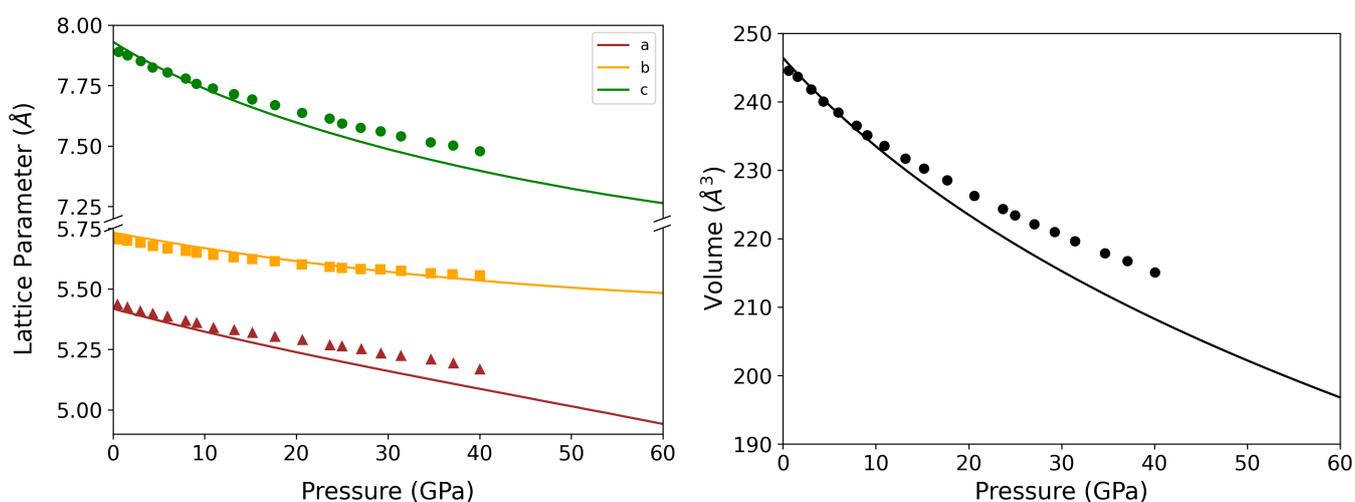


Figure 2. (left) Variation of the lattice parameters a , b , c of DyScO₃ perovskite with pressure. (right) Variation of the unit cell volume with pressure. In both panels, the curves represent our calculated values and the symbols represent the experimental values from Ref. [11].

Table 2. Comparison of the calculated and experimental unit cell volume, bulk modulus, and pressure derivative of the bulk modulus at zero pressure.

	$V_0(\text{Å}^3)$	$B_0(\text{GPa})$	B'_0
This work	246.7	164.7	4.42
Experiment Ref. [11]	245.3	189.4	7.68

The physical properties of the ABO_3 perovskites and the instabilities associated with soft modes found near the phase transitions are related to the distortions from the ideal cubic structure. Those distortions can be broken up into three components: two independent tilts of the BO_6 octahedra (with tilt angles respectively θ and ϕ), the distortion of those octahedra, and the displacement of the A cations [31–34]. In Table 3 we summarize the calculated values of the atomic positions in DyScO_3 , showing, in particular, the displacement of the Dy cation, the tilt angles, and the octahedra distortion index, compared to the values of the ideal cubic perovskite structure, and the experimental values published in Ref. [25]. (The tilt angles were calculated using Equations (A1) and (A2) in Ref. [33].)

Table 3. Summary of the atomic sites of DyScO_3 perovskite, showing the values of the internal parameters and tilt angles. The values for the ideal cubic perovskite structure are also provided, for reference.

		Ideal Cubic Perovskite	This Work 0 GPa	This Work 60 GPa	Exp. Ref. [25]
Wyckoff positions					
Dy: 4c	u_{Dy}	$\frac{1}{2}$	0.4788	0.4733	0.48262(5)
$(u_{\text{Dy}}, v_{\text{Dy}}, \frac{1}{4})$	v_{Dy}	$\frac{1}{2}$	0.4350	0.4165	0.43844(5)
Sc: 4b					
$(\frac{1}{2}, 0, 0)$					
O1: 4c	u_{O1}	$\frac{1}{2}$	0.6313	0.6312	0.6261(1)
$(u_{\text{O1}}, v_{\text{O1}}, \frac{1}{4})$	v_{O1}	0	0.0562	0.0464	0.0560(2)
O2: 8d	u_{O2}	$\frac{1}{4}$	0.1906	0.1717	0.1885(5)
$(u_{\text{O2}}, v_{\text{O2}}, w_{\text{O2}})$	v_{O2}	$\frac{1}{4}$	0.1936	0.1863	0.1937(6)
	w_{O2}	0	0.0684	0.0622	0.0658(6)
Tilt angles					
θ in [110]		0	21.26	19.10	21.00
ϕ in [001]		0	13.39	17.00	13.23

The variation of those values with pressure is shown in Figure 3. We observe that as the applied pressure increases the Dy cations move further away from the cubic perovskite positions, as the octahedra rotate further away around the [001] axes.

In Figure 4 we show the calculated pressure-evolution of near-neighbor distances in DyScO_3 . Our calculated values for the Sc-O interatomic distances, which are related to the ScO_6 octahedra, are typical for octahedrally coordinated scandium, and the octahedra show little distortion over the studied pressure range. The Dy-O distances show a larger variation and there is a gap between the first eight and the last four such distances, which makes it possible to distinguish between a first and second shell for the Dy sites and has led some authors to consider Dy as being 8-fold coordinated [25,35].

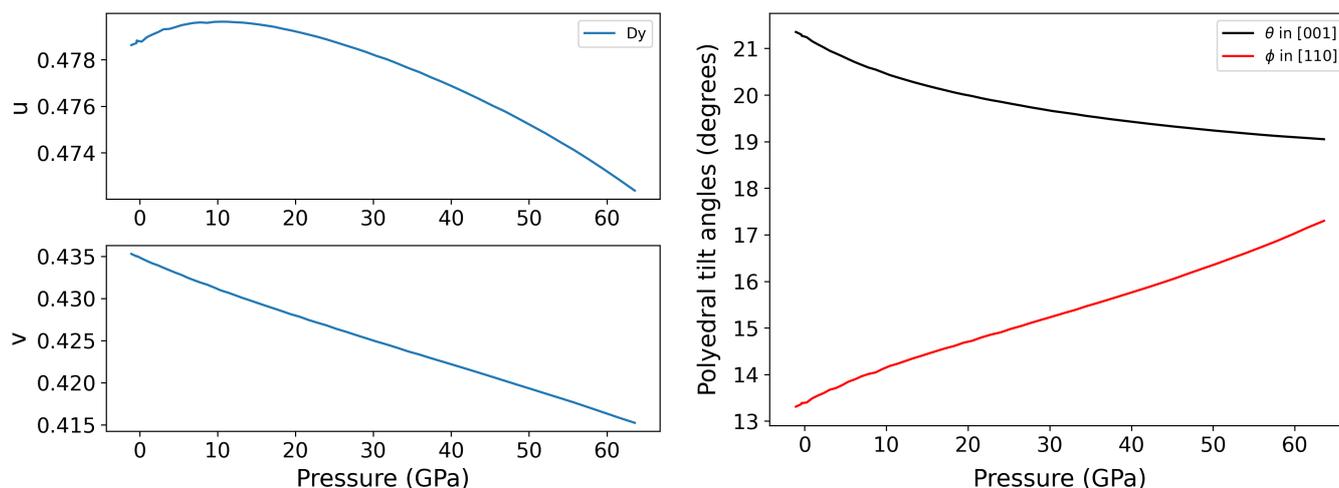


Figure 3. (left) Evolution of the internal parameters of the 4c Wickoff positions of Dy with pressure. (The values for the ideal cubic perovskite are $\frac{1}{2}$, $\frac{1}{2}$) (right) Tilt angles change with pressure.

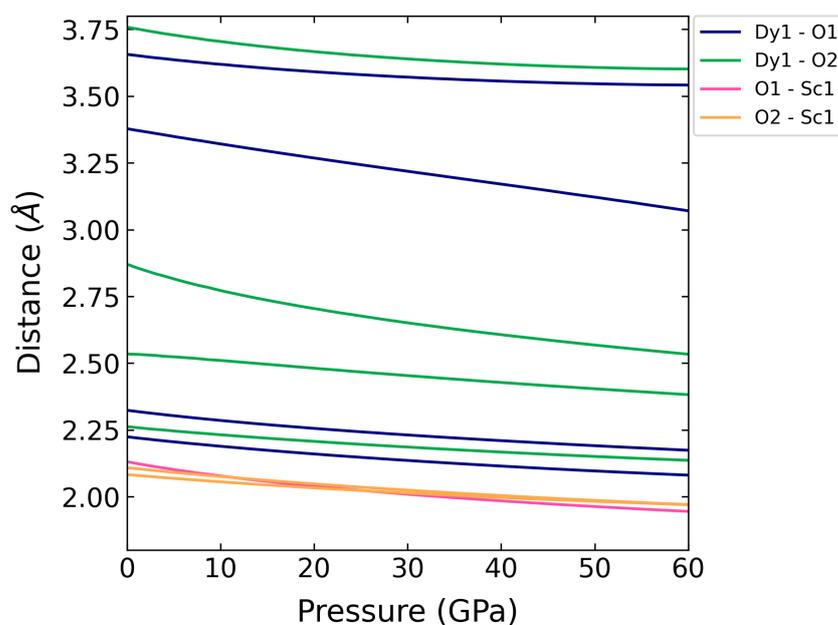


Figure 4. Evolution of interatomic distances with pressure in DyScO₃ perovskite. The O1-Sc1, O2-Sc1, and Dy1-O2 distances have all a multiplicity of two (each represented curve corresponds to two symmetrically equivalent but different bonds).

Related to the observed behavior for the cell volumes and the lattice parameters, we find that the calculated values of the polyhedra volumes stray somewhat away from the experimental ones at large pressures, but nonetheless the volume ratio of DyO₁₂ to ScO₆ stays above the value of four as suggested in Ref. [34] for a stable material. The ScO₆ octahedra are less compressible than the DyO₁₂ polyhedra, so the latter are the main contributors to the volume change of the material, see Figure 5.

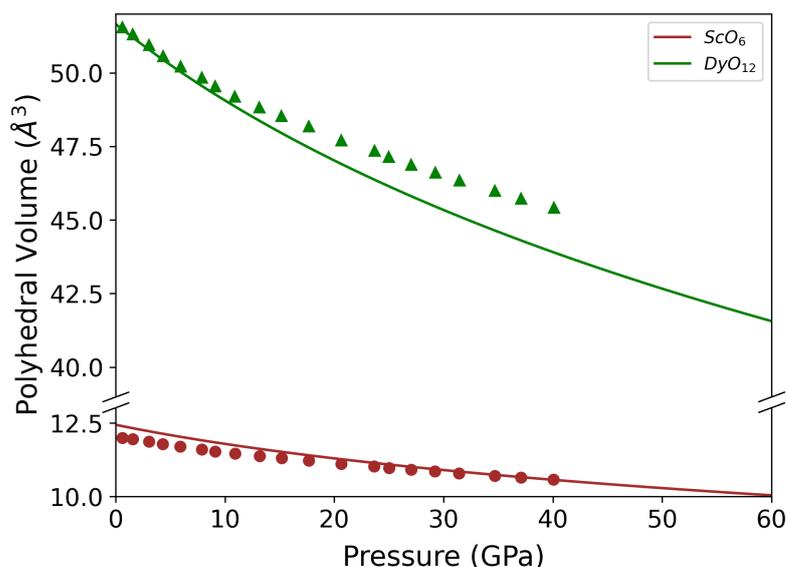


Figure 5. Variation of the polyhedral volumes of DyScO₃ perovskite with pressure. Solid curves represent the calculated values and symbols represent the experimental values from Ref. [11].

3.2. Elastic Properties

As described in Section 3.1, DyScO₃ has an orthorhombic structure with SG *Pbnm* No 62 and therefore there are nine independent elastic constants: C_{11} , C_{22} , C_{33} , C_{44} , C_{55} , C_{66} , C_{12} , C_{13} , and C_{23} [36]. The study of elastic constants is useful to understand both the mechanical properties and the structural stability of the material. Our calculated values at zero pressure of these constants for DyScO₃ are compared with other ab initio and experimental results in Table 4.

Table 4. Elastic constants (in GPa) of DyScO₃ perovskite at 0 GPa. (The *Pbnm* setting is used.)

	This Work	Experiment Ref. [37]	Theory Ref. [37]
C_{11}	281	302.4	310
C_{22}	332	345.3	339
C_{33}	213	254.6	236
C_{44}	100	103.9	96
C_{55}	76	86.0	73
C_{66}	81	84.4	77
C_{12}	121	124.4	96
C_{13}	105	130.0	132
C_{23}	119	132.5	106

The C_{11} , C_{22} , and C_{33} elastic constants, which are related to unidirectional compressions along the principal crystallographic directions, are much larger than the C_{44} , C_{55} , and C_{66} constants that are related to resistance against shear deformations. Consequently, there is a larger resistance to unidirectional compression than to shear deformation. The structure is stable at 0 GPa as it fulfills the Born stability criteria [38] given by the following expressions [39]:

$$C_{11} > 0 ; C_{11}C_{22} - C_{12}^2 > 0$$

$$C_{11}C_{22}C_{33} + 2C_{12}C_{13}C_{23} - C_{11}C_{23}^2 - C_{22}C_{13}^2 - C_{33}C_{12}^2 > 0$$

$$C_{44} > 0 ; C_{55} > 0 ; C_{66} > 0$$

When stress is applied to the crystal, the elastic properties of the material are no longer described by the elastic constants C_{ij} , and the elastic stiffness coefficients B_{ij} must be employed instead. If the external stress corresponds to a hydrostatic pressure p , the elastic stiffness coefficients for a crystal are expressed as [40,41]:

$$B_{ii} = C_{ii} - p, \quad i = 1, \dots, 6; \quad B_{ij} = C_{ij} + p, \quad i \neq j, \quad i, j = 1, 2, 3$$

where the C_{ij} are the elastic constants evaluated at the current stressed state. (As can be seen, the B_{ij} coefficients reduce to the C_{ij} when no pressure is applied.) The evolution of the elastic stiffness coefficients B_{ij} under pressure is depicted in Figure 6. The coefficients B_{22} , B_{33} , B_{12} , B_{23} , and B_{13} increase with pressure in the whole pressure range studied. On the other hand, B_{44} , B_{55} , and B_{66} decrease slowly. The B_{11} coefficient increases slightly up to 30 GPa and then decreases smoothly.

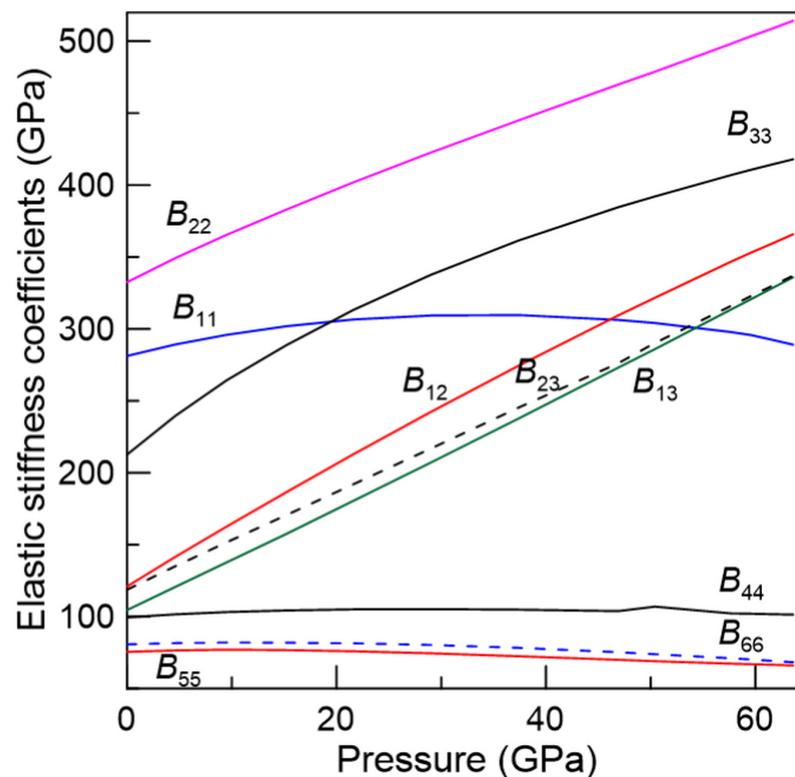


Figure 6. Calculated pressure evolution of the elastic stiffness coefficients B_{ij} of DyScO₃ perovskite.

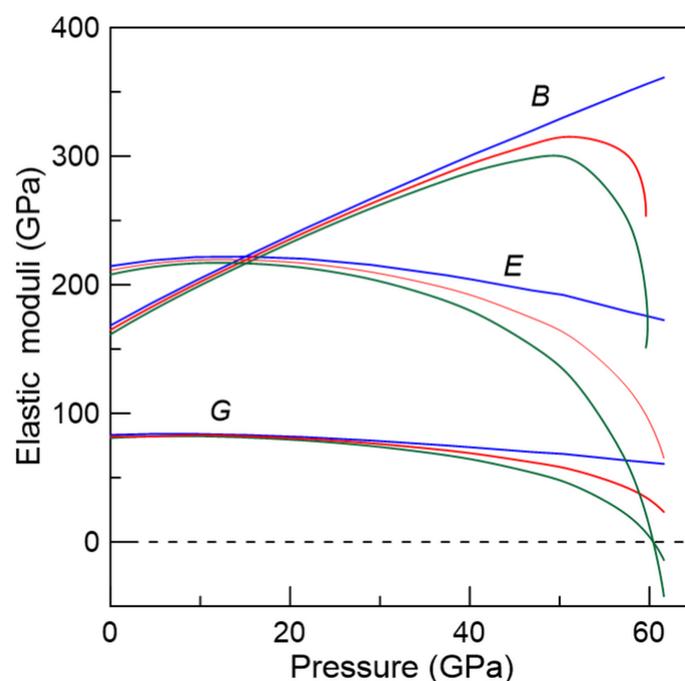
The stiffness coefficients allow us to obtain the elastic moduli that describe the major elastic properties of the material at any hydrostatic pressure. Employing standard relations for orthorhombic crystals [42], we have obtained the bulk modulus B , the shear modulus G , the young modulus E , the Poisson's ratio ν , and the B/G ratio. The values of the above-mentioned elastic moduli at zero pressure as calculated within the standard schemes of Voigt, Reuss, and Hill (with the latter one defined as the arithmetic average of the other two models [43–45]) are displayed in Table 5.

Table 5. Elastic moduli B_0 , E_0 and G_0 ; B_0/G_0 ratio; and Poisson's ratio ν_0 , calculated using the Voigt, Reuss, and Hill approximations from the calculated elastic coefficients at 0 GPa of DyScO₃ perovskite.

	Voigt	Reuss	Hill
Bulk modulus B_0 (GPa)	168.4	161.5	164.9
Shear modulus G_0 (GPa)	83.3	80.9	82.1
Young modulus E_0 (GPa)	214.5	207.9	211.2
Poisson's ratio ν_0	0.288	0.285	0.286
Bulk/Shear ratio B_0/G_0	2.02	2.00	2.01

It is worth noting that the bulk modulus, an important parameter related to the resistance of a material to compression, yields a similar value in the Hill formalism (164.9 GPa) than the one obtained from the EOS fitting (164.7 GPa), which reflects the consistency of the calculations. As the shear modulus G is associated with the resistance to plastic deformation, it is also interesting to analyze the B/G ratio. According to Pugh's criterion, if $B/G > 1.75$ the material is ductile, otherwise the material is brittle [46]. In the present case, our results indicate that DyScO₃ is a ductile material at zero pressure. The Poisson's coefficient ν is related to volume changes during uniaxial deformation. Lower coefficients indicate large volume changes during deformation; if $\nu = 0.5$ no volume changes occur during elastic deformation [40]. This parameter also provides information about the characteristics of the bonding forces. A value of $\nu = 0.25$ is a lower limit for central forces while $\nu = 0.5$ is an upper limit [40]. For DyScO₃, the calculated value of ν suggests that the forces in this material are predominantly central.

The calculated pressure evolution of the elastic moduli of DyScO₃ is presented in Figure 7. The bulk modulus B as calculated using Voigt's prescription increases in the whole pressure range under study. However, using Reuss' scheme the bulk modulus decreases dramatically around 50 GPa. This could be related to mechanical instability, as we will show in the next paragraph. As for the shear and Young moduli (G and E), E reaches a shallow maximum at about 14 GPa, then both moduli decrease with pressure and tend to get closer in values.

**Figure 7.** Pressure evolution of the elastic moduli: bulk modulus B , Young modulus E , and Shear modulus G of DyScO₃ perovskite. The blue line represents the moduli in the Voigt approximation, the green line in the Reuss approximation, and the red one in the Hill approximation.

The B/G ratio becomes larger under pressure and so the material becomes more ductile under compression. The Poisson's ratio increases rapidly with pressure, taking a value of 0.45 in the Hill expression at 60 GPa, see Figure 8.

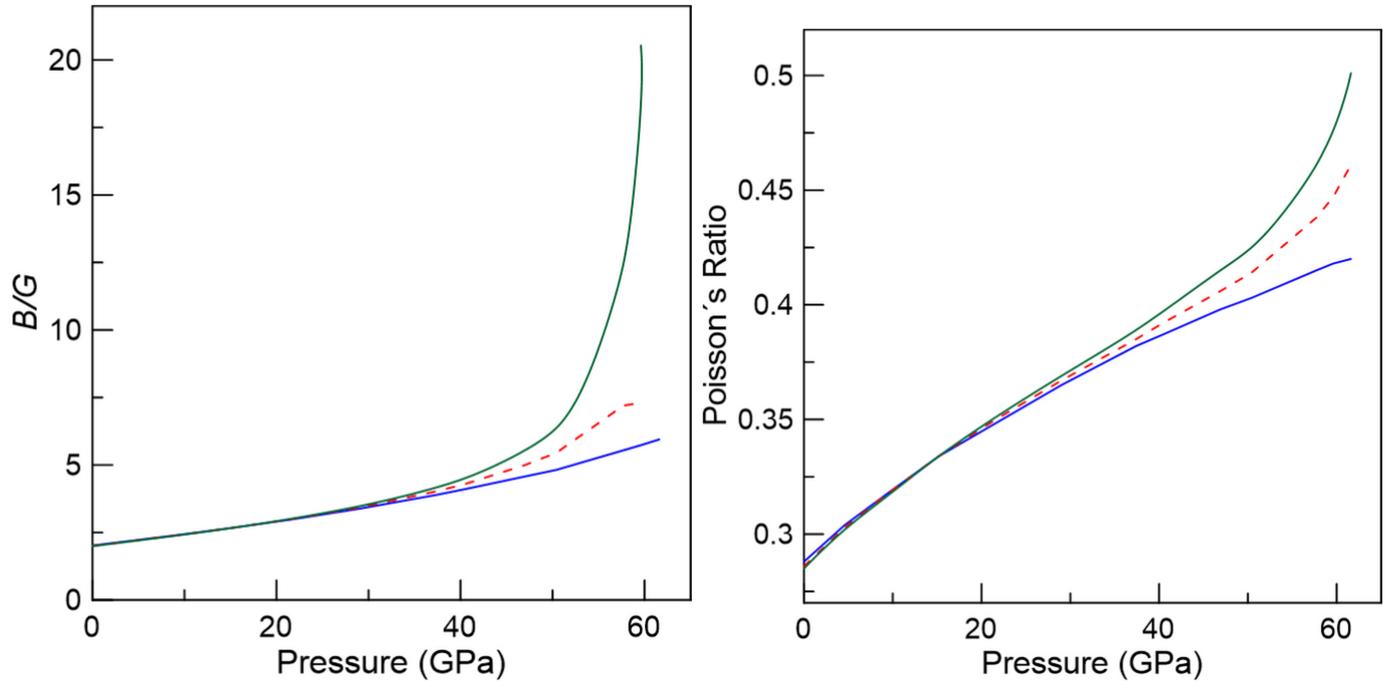


Figure 8. (left) Pressure evolution of the B/G ratio of DyScO_3 perovskite. (right) Pressure evolution of the Poisson's ratio. The blue line represents the moduli in the Voigt approximation, the green line in the Reuss approximation, and the red one in the Hill approximation.

To conclude, we analyze the mechanical stability of DyScO_3 perovskite under pressure. As stated before, when hydrostatic pressure is applied, the elastic stiffness coefficients must be employed instead of the elastic constants [38]. Therefore, the stability conditions are modified, and one needs to use the following generalized stability conditions for an orthorhombic crystal [38,47]:

$$M_1 = B_{11} > 0; M_2 = B_{11}B_{22} - B_{12}^2 > 0$$

$$M_3 = B_{11}B_{22}B_{33} + 2B_{12}B_{13}B_{23} - B_{11}B_{23}^2 - B_{22}B_{13}^2 - B_{33}B_{12}^2 > 0$$

$$M_4 = B_{44} > 0; M_5 = B_{55} > 0; M_6 = B_{66} > 0$$

The crystalline structure is mechanically stable when all the above criteria are fulfilled. The evolution of these generalized Born stability criteria M_i ($i = 1$ to 6) as functions of pressure is depicted in Figure 9 where it can be observed that the M_3 criterion is violated above 60 GPa, indicating that DyScO_3 would become mechanically unstable at such pressure. Above this pressure an amorphization or a pressure-driven phase transition can occur, a result that may entice further experimental and theoretical research along this line.

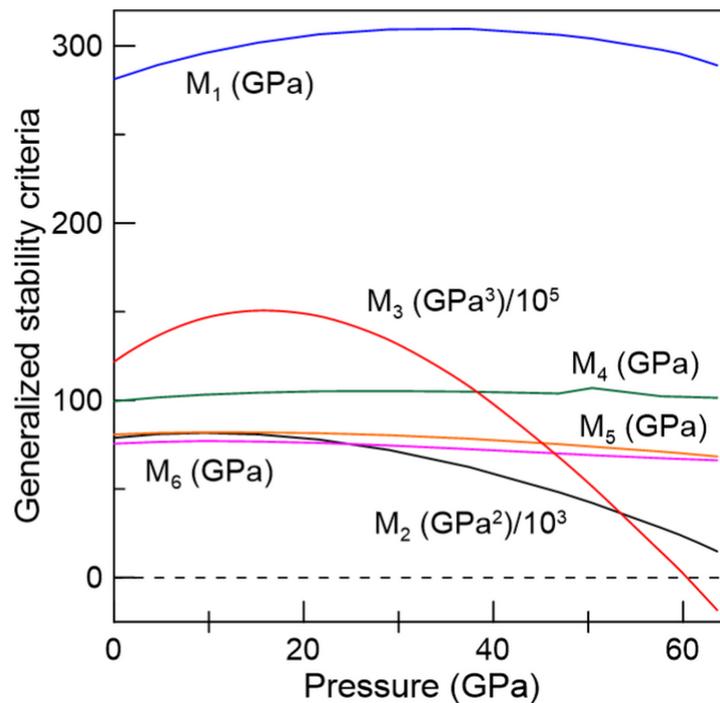


Figure 9. Generalized stability criteria as a function of pressure for DyScO₃ perovskite.

3.3. Vibrational Properties

We have also explored the vibrational properties of DyScO₃ perovskite at high pressure using the small-displacements method. The mechanical representation of DyScO₃ yields the following multiplicities for the irreducible representations of the modes at the zone center: $M = 7A_g (R) + 8A_u (S) + 5B_{1g} (R) + 10B_{1u} (IR) + 7B_{2g} (R) + 8B_{2u} (IR) + 5B_{3g} (R) + 10B_{3u} (IR)$; of which the A_g , B_{1g} , B_{2g} , and B_{3g} modes are active in Raman (R); the B_{1u} , B_{2u} , and B_{3u} modes are infrared-active (IR); and the A_u modes are silent (S). In Figure 10 we display the behavior of the phonon modes at the Γ point (Raman, infrared and silent modes) under applied hydrostatic pressure.

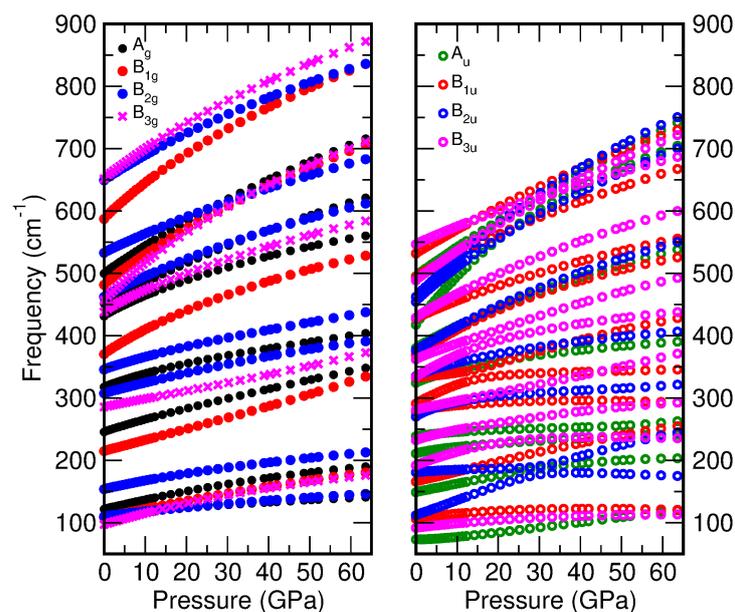


Figure 10. (left) Pressure dependence of the calculated Raman-active modes of DyScO₃ perovskite; (right) idem for the infrared and silent modes.

The present calculated values of the frequencies are in overall agreement with previous theoretical values obtained by Delugas et al. [48]. The pressure behavior in the very large region investigated in our work displays some curvature as evidenced in Figure 10 but up to about 15 GPa it can be approximated fairly well by a linear relation. The corresponding fit to the theoretical data yields the values of the pressure coefficients and calculated mode Grüneisen parameters shown in Table 6. As can be seen, they are in overall agreement with the experimental data obtained by Bura et al. [11].

The calculated phonon-dispersion curves along selected paths within the Brillouin zone with endpoints at the high symmetry points for SG 62 are plotted in Figure 11. The phonon dispersion curves at ambient pressure in these figures correspond to a dynamically stable behavior where all the frequencies are positive whereas above around 63 GPa the structure becomes dynamically unstable, displaying phonon branches with negative (imaginary) frequencies around the zone-boundary X point. This instability becomes more accentuated as the pressure is further increased, a pressure regime that we have not considered in this paper.

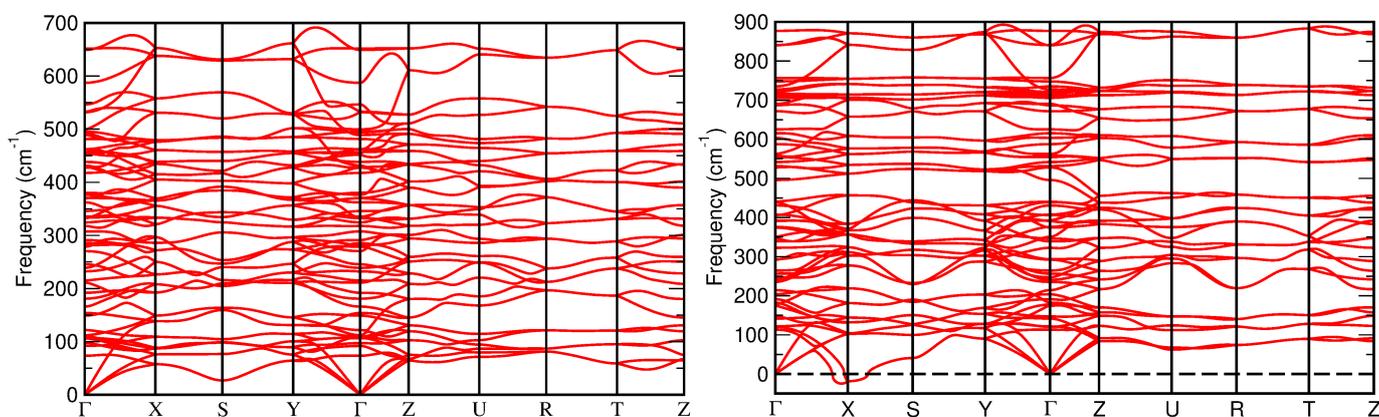


Figure 11. (left) Phonon dispersion curves of DyScO₃ perovskite at ambient pressure and (right) at 66 GPa.

To provide more detail about the nature of the lattice vibrations, the spectral distribution of the phonon density of states (phDOS) was obtained along with the partial phonon density of states, to extract the contribution of each atomic species in the local lattice configuration. Figure 12 shows phDOS plots at ambient pressure. The shape of the phDOS curve above 150 cm⁻¹ is dominated by the oxygen atoms, which are present in larger numbers. At low frequencies (up to 150 cm⁻¹) the phonon peaks correspond mainly to the heavy Dy atoms, whereas the contribution of these atoms is very small at higher frequencies. In the intermediate frequency region (from 150 cm⁻¹ to 450 cm⁻¹) there are significant contributions from both the Sc and O atoms.

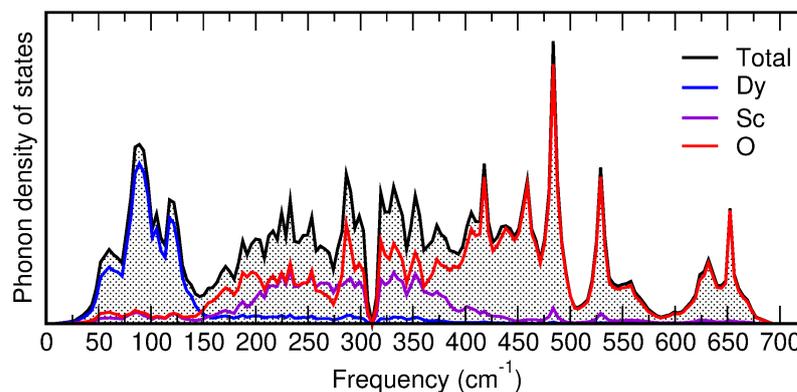


Figure 12. Total phonon density of states and contribution per species, as calculated at ambient pressure for DyScO₃ perovskite.

Table 6. Raman (R), infrared (IR), and silent (S) mode frequencies as calculated in this work for DyScO₃ perovskite, and their respective pressure coefficients obtained by fitting in the range up to 15 GPa with a linear equation $\omega(p) = \omega_0 + \alpha p$. Mode Grüneisen parameters (γ) are also provided using the calculated bulk modulus B_0 (164.7 GPa). Experimental values from Ref. [11], where available, are given in parenthesis.

Mode (R)	ω_0 (cm ⁻¹)	α (cm ⁻¹ GPa ⁻¹)	γ
A _g	107.6	0.91	1.39
	122.3	1.46	1.96
	246.6 (253.2)	1.92 (1.48)	1.28 (1.11)
	319.2 (327.7)	2.02 (1.59)	1.04 (0.92)
	433.6	3.05	1.16
	452.9 (458.3)	3.28 (2.74)	1.19 (1.13)
	501.3	4.32	1.42
B _{1g}	108.7	1.34	2.03
	215.4 (226.4)	1.75 (1.62)	1.34 (1.36)
	373.3 (380.1)	3.65 (2.90)	1.61 (1.45)
	484.2 (476)	4.61 (2.50)	1.57 (1.00)
	590.8	5.42	1.51
B _{2g}	110.7	0.75	1.12
	154.7 (157.4)	1.30 (1.15)	1.38 (1.38)
	308.3 (308.7)	1.65 (1.80)	0.88 (1.10)
	347.2	1.89	0.90
	464.1	3.11	1.10
	534.4	2.92	0.90
	651.2	3.83	0.97
B _{3g}	97.8	1.74	2.93
	286.0	1.29	0.74
	441.5 (458.3)	3.15 (2.74)	1.17 (1.13)
	459.8	5.66	2.03
	654.0	4.57	1.15
Mode (IR)	ω_0 (cm ⁻¹)	α (cm ⁻¹ GPa ⁻¹)	γ
B _{1u}	107.2	0.73	1.12
	166.2	1.58	1.56
	280.6	0.76	0.45
	294.0	2.68	1.50
	332.1	0.89	0.44
	372.2	3.71	1.64
	429.2	2.69	1.03
	495.7	3.49	1.16
	533.1	3.73	1.15
B _{2u}	112.2	2.41	3.53
	181.6	0.33	0.30
	274.7	1.81	1.08
	337.7	1.89	0.92
	377.6	3.69	1.61
	457.9	5.19	1.87
	464.0	6.03	2.14
B _{3u}	92.3	0.66	1.18
	193.4	2.02	1.72
	234.0	1.49	1.05
	281.7	1.37	0.80
	336.0	3.27	1.60
	363.5	2.39	1.08
	430.7	3.64	1.39
	490.2	4.58	1.54
	547.7	2.76	0.83

Table 6. Cont.

Mode (S)	ω_o (cm ⁻¹)	α (cm ⁻¹ GPa ⁻¹)	γ
<i>A_u</i>	72.6	0.48	1.09
	149.3	1.34	1.48
	212.5	0.91	0.70
	239.7	0.58	0.40
	326.5	1.90	0.96
	381.7	3.37	1.45
	421.7	7.31	2.85
	496.8	4.51	1.49

3.4. Prospective High-Pressure Post-Perovskite Phase

The recent experimental work of Bura et al. [11] shows that the DyScO₃ perovskite has a large stability range, up to at least 40 GPa, which are the highest pressures for which experimental results on this compound exist. For this reason, up to now we have been concerned with the local stability of the perovskite phase of DyScO₃, the only one so far observed. However, other perovskites, like MgSiO₃, undergo a pressure-induced transition to a post-perovskite phase with a *Cmcm*-type structure [49–51]. We have performed simulations on this post-perovskite structure for DyScO₃ (see Table 7) and our results show that this new phase has lower enthalpy above 14.4 GPa, see Figure 13. The study of the phonon dispersion curves at 15 GPa (Figure 14) shows that this phase is also dynamically stable. The existence of large kinetic barriers (and the enhanced local stability of the perovskite phase) might explain why this high-pressure phase was not observed in the experimental study of Ref. [11], but more experimental work, involving heating at high pressure to overcome barriers, would be needed to test this hypothesis. In the Supplementary Material we provide some extra results for this proposed high-pressure phase of DyScO₃, but its full study will not be pursued here.

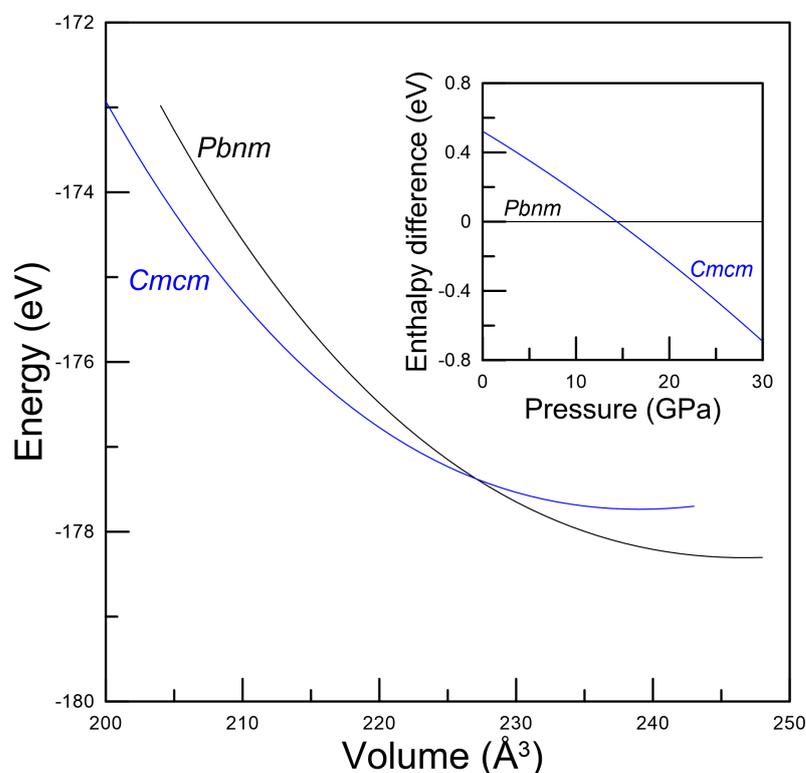


Figure 13. Energy-volume curves and (inset) enthalpy-pressure curves for the *Pbnm* (perovskite) and *Cmcm* (post-perovskite) phases of DyScO₃ (black and blue curves, respectively). The enthalpies are given with respect to that of the perovskite phase.

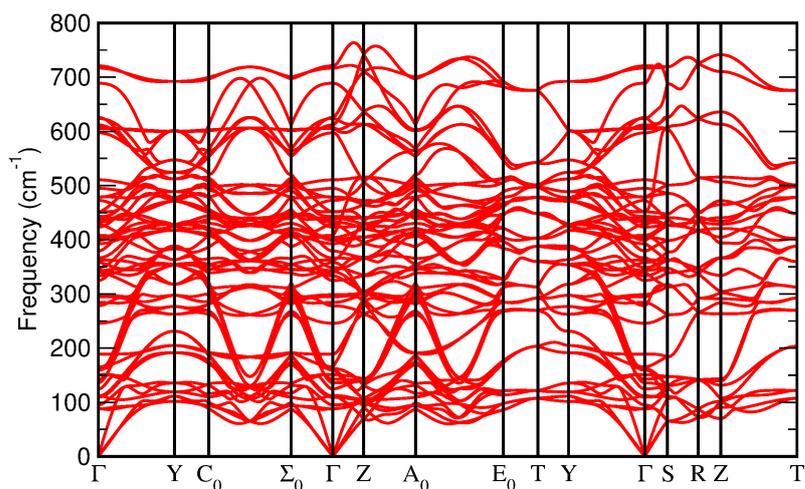


Figure 14. Calculated phonon dispersion curves of the *Cmcm* phase at 15 GPa.

Table 7. Crystallographic description of *Cmcm* at 15 GPa.

Space Group and Lattice Parameters	Wyckoff Positions
No 63 - C 2/m 2/c 2 ₁ /m (<i>Cmcm</i>)	Dy 4c (0.000, 0.756, 0.250)
$a = 3.081 \text{ \AA}$; $b = 9.901 \text{ \AA}$; $c = 7.278 \text{ \AA}$	Sc 4b (0.000, 0.500, 0.000)
$\alpha = \beta = \gamma = 90^\circ$	O1 4c (0.000, 0.414, 0.250)
$V = 222.00 \text{ \AA}^3$	O2 8f (0.000, 0.139, 0.063)

3.5. Summary and Conclusions

We have performed an ab initio study of DyScO₃ perovskite both at ambient and under hydrostatic pressure up to just above 60 GPa. We have analyzed the evolution of its structural, dynamic, and elastic properties. The calculated lattice parameters compare well with experimental values up to 10–15 GPa, above which pressures the experiments are probably affected by the loss of hydrostatic conditions of the pressure transmitting medium. The equation of state, the evolution of interatomic distances with pressure, and the compressibility were calculated. The DyO₁₂ polyhedra account for most of the volume reduction under compression. The phDOS and partial phDOS were analyzed to determine the contribution of each atomic species to the vibrational spectrum. The evolution of all the Raman, infrared and silent phonon modes was obtained along with their corresponding Grüneisen parameters. At about 63 GPa, imaginary frequencies appear around the X point of the Brillouin zone and therefore the structure becomes dynamically unstable. The elastic constants and the elastic stiffness coefficients were also determined. This allowed us to study the major elastic moduli (B , G , E , ν). The value of the B/G ratio points out that this material is ductile and that its ductility increases under pressure. Above 60 GPa the phase becomes mechanically unstable. We have also considered the possibility of a high-pressure post-perovskite phase with *Cmcm* symmetry. Our calculations suggest that it is favored over the perovskite phase above 14.4 GPa, which is a pressure within the range of observation of the perovskite phase and well below its predicted local (i.e., dynamical and elastic) instabilities. This prediction should thus entice further experimental and theoretical studies on DyScO₃.

Supplementary Materials: The following supporting information can be downloaded at: <https://www.mdpi.com/article/10.3390/cryst13020165/s1>, Figure S1. Variation of the lattice parameters a , b , c with pressure for the proposed *Cmcm* phase. Table S1. Elastic moduli B , E and G ; B/G ratio; and Poisson's ratio ν for the proposed *Cmcm* phase., calculated using the Voigt, Reuss, and Hill approximations from the calculated elastic coefficients at 15 GPa. Figure S2. Calculated pressure evolution of the elastic stiffness coefficients for the proposed *Cmcm* phase. Figure S3. M_3 generalized stability criteria as a function of pressure for the proposed *Cmcm* phase. Figure S4. Pressure dependence of the

calculated Raman-active modes of the *Cmcm* phase. Figure S5. Pressure dependence of the calculated infrared modes of the *Cmcm* phase. Figure S6. Pressure dependence of the calculated silent modes of the *Cmcm* phase. Table S2. Calculated Raman (R), infrared (IR), and silent (S) mode frequencies for the proposed *Cmcm* phase as calculated in this work, and their respective pressure coefficients obtained by fitting in the range up to 60 GPa with a quadratic equation $\omega(p) = \omega_0 + \alpha p + \beta p^2$. Figure S7. Phonon dispersion curves at 65.7 GPa for the proposed *Cmcm* phase.

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