



Article Finite Element Analysis for the Mechanism of Stress Wave Propagation and Crack Extension Due to Blasting of a Frozen Rock Mass

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Abstract: The propagation mechanism of explosion stress waves in frozen rock mass is the main factor affecting the blasting efficiency and safety construction of strip mines in alpine cold regions. In order to study explosion stress wave propagation and crack extension in the blasting process of frozen rock mass with ice-filled cracks, RFPA^{2D} is adopted to simulate the influence of the geometric parameters of ice-filled cracks (ice-filled crack thickness d, normal distance R from blasting hole to the ice-filled crack, and ice-filled crack angle α), loading intensity and loading rate on the explosion stress wave propagation effect and the damage range. The results show: The attenuation trend of explosion stress waves decreases gradually with an increase of thickness (e.g., In the case of R is 0.2 m, when d is 0.02 m, 0.04 m, and 0.08 m, the calculated attenuation factor of the minimum principal stress peak value is 7.128%, 18.056%, and 30.035%, respectively), and it decreases slightly with an increase of normal distance and ice-filled crack angle. The damage elements range of the ice-filled crack decreases when the ice-filled crack thickness and normal distance increases. The loading intensity and the loading rate have a significant influence on blasting hole fracture patterns. The ice-filled crack has a guiding effect on the growth of blasting cracks at the blasting hole. Nevertheless, the existence of ice-filled cracks inhibits the propagation of explosion stress waves in frozen rock mass.

Keywords: frozen rock mass blasting; ice-filled crack; explosion stress wave propagation; attenuation factor; numerical simulation

1. Introduction

The design and disaster prevention of open pit blasting mining of mineral resources in cold regions have become key issues in the field of energy safety mining. There are a large number of naturally formed intermittent joint cracks, bedding and faults in rock mass. The existence of these structural planes affects the mechanical properties, vibration, permeability, energy transfer and other properties of rock mass. Affected by low temperature, the water in the primary fissures of open pit slopes becomes ice, which forms frozen rock mass [1–4]. The propagation and attenuation of explosion stress waves in frozen rock mass is slightly different from that in conventional rock mass, which affects the blasting effect and safety of strip mining in cold regions. Therefore, it is of great significance to study the explosion stress wave propagation and crack extension of ice-filled crack rock mass under explosion loading. This study will improve the efficiency of blasting energy utilization, blasting effect and disaster prevention of rock mass engineering in cold regions [5–7].

At present, many scholars have carried out a wealth of research on stress wave propagation in jointed rock masses. In terms of theoretical calculations, it is mainly divided



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Copyright: © 2023 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). into the discontinuous displacement method, the equivalent continuous medium method, and the continuous and discontinuous coupling method [8–10] to study the propagation characteristics of stress waves. The discontinuous displacement method is mainly used to analyze the stress wave propagation in a single crack or a group of parallel cracks. The fewer cracks there are, the better the analysis effect is. The equivalent continuous medium method can quickly calculate the propagation of stress waves in rock mass under a large number of cracks and uniform distribution. The continuous and discontinuous coupling method is used to analyze macroscopic joints and mesoscopic rock fissures in rock mass. Currently, this method focuses on the one-dimensional propagation law of stress waves [11,12].

In terms of physical tests, the separation Hopkinson pressure bar (SHPB) device has become the main research method to study the propagation of explosive stress waves in jointed rock masses [13–15]. Chen et al. [16] obtained the relationship between the transmission coefficient and the contact surface by the stress wave propagation experiment in artificial rock fractures. Kumar et al. [17] investigated the rate-dependent mechanical behavior of jointed rock with a non-persistent joint with different infill conditions under varying strain rates, i.e., 10^{-4} to 130 s^{-1} using an SHPB and static uniaxial compression test set-up. Certainly, it is a good method to study the propagation characteristics of explosion stress waves through the blasting simulation test. Luo et al. [18] used the dynamic caustics test system to study the penetration process of the main crack of the slotted hole and wing cracks of different angles. They drew the conclusion that the 90° pre-crack has a certain inhibitory effect on the reflected stretching wave. Ram et al. [19] studied the interaction between explosion waves and a structure by electric explosion technology.

The numerical simulation methods, in comparison with theoretical and experimental studies, provide easier and more economical conditions for studying stress wave propagation in jointed rock masses, especially for complex cases where theoretical and experimental solutions seem impossible.

The continuum-based method mainly contains the finite element method (FEM), XFEM (extended FEM), SPH (smoothed particle hydrodynamics), etc. [20–23]. Liang et al. [24] studied the dynamic fracture properties of rocks under different static stress conditions by RFPA^{2D}, and concluded that the crack propagation path became more discontinuous and rougher in a smaller-heterogeneity parameter case. Bendezu et al. [25] obtained the advantages and limitations of three methods (XFEM, the conventional finite element method (FEM) using a remeshing technique, and the element deletion method) that simulate the evolution of a rock fragmentation process. Based on the experimentally obtained mechanical properties, experienced peak pressure values inside the rock samples and blast-induced fracture patterns, Banadaki et al. [26] calibrated the Johnson-Holmquist model parameters in ANSYS Autodyn. Zhao et al. [27] analyzed the blasting-induced fracture propagation in coal masses by LS-DYNA, considering the dynamic compressive and tensile failure.

The discontinuum-based methods include the DEM (discrete element method) and the DDA (discontinuous deformation analysis). Yari et al. [28] studied the effect of the position of the joints relative to the blast hole on the blast wave propagation by 3D DEM models. Lak et al. [29] simulated the process of extension of blast-induced fractures in rock masses by the DEM, which considered fracture propagation from both the rock mass inherent fractures and newly induced cracks. Hajibagherpour et al. [30] simulated the mechanism of rock fragmentation due to blast-induced shock waves in a single blast hole by UDEC. Ning et al. [31] extended the DDA to model rock mass fracturing by coupling the rock mass failure process and the penetration effect of the explosion gas based on a generalized artificial joint concept.

Coupled or hybrid continuum-discontinuum-based methods include the FEM-SPH method [32], the DEM-SPH method [33], the MPM (material point method), the CDEM (continuum-discontinuum element method) [34], and the combined finite-discrete element method (FDEM) [35,36], etc. Trivino et al. [37] simulated blasting-induced crack initiation and propagation in a granitic outcrop using FDEM. Zhao et al. [38] studied the blasting

effect disturbed by joint strength, joint stiffness, joint spacing, joint angle and other factors by CDEM.

The above-related research mainly focuses on the propagation characteristics of stress waves in jointed rock mass, while the explosive stress wave propagation in frozen rock masses at low temperature is slightly involved. Frozen rock mass with ice-filled cracks is very common in the mining process of mineral resources in cold regions, and its properties are different from that of conventional rock mass [39]. It is of practical value to study the propagation process of explosion stress waves in frozen rock mass with ice-filled cracks, the growth pattern of blasting cracks in frozen rock mass, and the attenuation of explosion stress waves after passing through ice, for the safe mining of strip mines in cold regions.

The aim of this study is to explore explosion stress wave propagation and crack extension in the blasting process of frozen rock mass with ice-filled cracks. The numerical model of frozen rock mass with ice-filled cracks is established by RFPA^{2D} in Section 2. The influences of the geometrical parameters of ice-filled cracks (ice-filled crack thickness D, normal distance R from blasting hole to ice-filled crack, and ice-filled crack angle α), loading intensity and loading rate on the explosion stress wave propagation effect and the damage range are mainly analyzed in Section 3. This research can provide theoretical suggestions for improving the efficiency and disaster prevention of blasting engineering in cold regions.

2. The Principle of RFPA^{2D}

2.1. Overview of Mesoscopic

Rock failure process analysis (RFPA) is used to simulate the failure process of frozen rock mass with ice-filled cracks. However, we know that the rock is a heterogeneous material filled by the disorder of micro-structures, which plays a significant role on the mechanical properties of rock [40,41]. Therefore, rock heterogeneity should be considered and implemented in the numerical model. Rock heterogeneity can be well characterized by using the statistical method. In RFPA, the numerical testing sample is composed of elements with the same shape and size. It is assumed that the distribution of elemental mechanical parameters, including the strength, Poisson ratio, elastic modulus and density, can be depicted by the Weibull distribution function [42], as follows:

$$\phi(u) = \frac{m}{u_0} \left(\frac{u}{u_0}\right)^{m-1} \exp\left[-\left(\frac{u}{u_0}\right)^m\right] \tag{1}$$

where *u* is the mechanical and dynamic properties of elements, such as elastic modulus, strength, and density; u_0 is defined as the mean value of the element parameter; *m* is a shape parameter which is defined as the homogeneity index of the material; and $\varphi(u)$ is the distribution function of mechanical properties.

2.2. Evolutionary Damage Principle of RFPA Meso-Elements

RFPA^{2D} uses the four-node iso-parametric element to describe the basic element. All of the elements are considered to be elastic and isotropic. The elastic damage constitutive method is adopted to elaborate the stress–strain relationship. The stress–strain curve of each element is considered as linearly elastic. When the damage threshold is reached, the maximum tensile stress criterion is used to judge the damage and failure of the element in tensile state, and the Mohr-Coulomb criterion is used to judge the damage and failure of the element in compressive and shear state. As the damage progresses, the elastic modulus of the element is gradually degraded. The modified elastic modulus can be expressed as follows:

$$E = (1 - D) \times E_0 \tag{2}$$

where *E* and E_0 are the elastic modulus after damage and the initial elastic modulus, respectively; and *D* is the damage variable.

Figure 1 shows the elastic damage constitutive relation of the element under uniaxial stress state. When shear failure occurs to the elements, the Mohr-Coulomb criterion is adopted:

$$F = \sigma_1 - \sigma_3 \frac{1 + \sin \varphi}{1 - \sin \varphi} \ge f_c \tag{3}$$

where σ_1 , σ_3 and f_c are the maximum principal stress, the minimum principal stress and the uniaxial compressive strength, respectively; and ϕ is the friction angle.



Figure 1. Elastic damage constitutive relation of the element under uniaxial stress state.

Damage variable *D* is introduced:

$$D = \begin{cases} 0 & \varepsilon \le \varepsilon_{c0} \\ 1 - \frac{\sigma_{cr}}{\varepsilon E_0} & \varepsilon_{c0} \le \varepsilon \end{cases}$$
(4)

where σ_{cr} is the compressive residual strength and ε_{c0} represents the maximum compressive strain.

When tensile failure occurs to the elements, the following is adopted:

$$\sigma_3 \le -f_t \tag{5}$$

Damage variable *D* is introduced:

$$D = \begin{cases} 0 & \varepsilon \leq \varepsilon_{t0} \\ 1 - \frac{\sigma_{tr}}{\varepsilon E_0} & \varepsilon_{t0} < \varepsilon \leq \varepsilon_{tu} \\ 1 & \varepsilon > \varepsilon_{tu} \end{cases}$$
(6)

where f_t and σ_{tr} are the tensile strength of rock and the residual strength of tensile damage, respectively; ε_{t0} and ε_{tu} represent the maximum tensile strain and the ultimate tensile strain, respectively. When the maximum tensile strain of the rock element is reached, it loses carrying capacity.

2.3. RFPA Solution for Dynamic Finite Element Equations

The dynamic equilibrium equation for each node in motion can be expressed as [43]:

$$M\ddot{u} + C\dot{u} + Ku = Q \tag{7}$$

where \ddot{u} , \dot{u} , and u are the constant vectors of displacement, velocity and acceleration at t, respectively. K, M and C are the stiffness matrix, mass matrix and damping matrix of the system, respectively.

By substituting the equations relating velocity, acceleration and displacement in the Newmark method:

$$\ddot{u}_{t+\Delta t} = \frac{1}{\beta \Delta t^2} (u_{t+\Delta t} - u_t) - \frac{1}{\beta \Delta t} \dot{u}_t - \left(\frac{1}{2\beta} - 1\right) \ddot{u}_t \tag{8}$$

$$\ddot{u}_{t+\Delta t} = \frac{\gamma}{\beta \Delta t} (u_{t+\Delta t} - u_t) + \left(1 - \frac{1}{\beta}\right) \dot{u}_t - \left(\frac{\gamma}{2\beta} - 1\right) \ddot{u}_{t+\Delta t} \Delta t \tag{9}$$

By substituting Equations (8) and (9) into the dynamic equilibrium Equation (7):

$$\widehat{K}u_{t+\Delta t} = \widehat{Q}u_{t+\Delta t} \tag{10}$$

$$\widehat{K} = K + \frac{1}{\beta \Delta t^2} M + \frac{\gamma}{\beta \Delta t} C$$
(11)

$$\widehat{Q}_{t+\Delta t} = Q_{t+\Delta t} + M \Big[\frac{1}{\beta \Delta t^2} (\alpha_{t+\Delta t} - \alpha_t) - \frac{1}{\beta \Delta t} \dot{\alpha}_t - \left(\frac{1}{2\beta} - 1\right) \ddot{\alpha}_t \Big]
+ C \Big[\frac{\gamma}{\beta \Delta t} (\alpha_{t+\Delta t} - \alpha_t) + \left(1 - \frac{\gamma}{\beta}\right) \dot{\alpha}_t + \left(1 - \frac{\gamma}{2\beta}\right) \Delta t \ddot{\alpha}_t \Big]$$
(12)

where β and γ are the integration coefficients of the Newmark method. When $\gamma \ge 0.5$, $\beta \ge 0.25 \times (\gamma + 0.5)^2$, the Newmark method is unconditionally stable.

3. Numerical Simulation Blasting Process Analysis of Frozen Rock Mass with Ice-filled Cracks

3.1. Model Setup

RFPA^{2D} has been widely used in the numerical simulation of stress wave propagation in conventional jointed rock masses [44,45]. Liang et al. [24] used RFPA^{2D} to carry out the dynamic fracture characteristics of fractured rock under different static stress conditions and compared it with the experimental results of Yang et al. [46]. The simulation results showed that RFPA^{2D} could well simulate crack propagation and stress wave attenuation under dynamic stress, which was in good agreement with the experimental results. In addition, some achievements have been made on the effect of ice-filled cracks on the propagation efficiency of rock explosion stress waves [47]. Therefore, explosion stress wave propagation and crack extension of frozen rock mass with ice-filled cracks during blasting is studied by RFPA^{2D}.

The blasting model of frozen rock mass is 4 m in length and 4 m in height. The mesh size is 0.01 m \times 0.01 m. Single-hole blasting is adopted, and the blasting hole radius r is 0.05 m. The blasting load is simplified into triangular waves [24,45]. The loading case is shown in Figure 2. The total duration is 3×10^{-4} s, and the single-step loading time is 2×10^{-6} s.



Figure 2. Explosion stress waves applied on the blasting hole.

The material parameters of the numerical simulation of rock and ice are shown in Table 1 [48]. Model I is used to study the influence of ice-filled crack thickness d, normal distance R from the ice-filled crack to the center of the blasting hole, loading intensity and loading rate on the propagation process of frozen rock mass explosion stress waves (Table 2). Model II is used to study the influences of the ice-filled crack angle α on the propagation process of frozen rock mass explosion stress waves (Table 2).

	Elasticity Modulus (MPa)	m	Compressive Strength (MPa)	m	Poisson Ratio	Friction Angle	Density (kg $ imes$ m ⁻³)
Rock	32,000	5	147	5	0.3	30°	2600
Ice	6000	10	8	10	0.35	26.5°	917



Table 2. Calculation Model I.

3.2. Analysis of Blasting Failure Process of Intact Frozen Rock Mass

3.2.1. Blasting Failure Process

Figure 3 shows the intact frozen rock mass explosion stress wave propagation and failure process. The cracks produced by blasting mainly occur near the blasting hole. The cracks show a uniform and divergent extension pattern. There is no ice-filled crack in the model, and explosion stress waves propagate equally on the left and right sides. At 0.54×10^{-4} s, damage elements begin to appear near the blasting hole; and at 1.48×10^{-4} s, macro blasting cracks are formed. After that, the macro cracks stop extending as the explosion stress wave has propagated out.

3.2.2. The Minimum Principal Stress at the Monitoring Point A

The monitoring point A is set (3.2 m, 2 m). In RFPA^{2D}, the pressure is positive, and the tensile stress is negative. Before 2.34×10^{-4} s, the minimum principal stress at the monitoring point A does not change, and the stress waves do not reach the monitoring point A. From 2.34×10^{-4} s to 3.04×10^{-4} s, the monitoring point is squeezed by stress waves, showing stress changes; and the peak value of the minimum principal stress is 0.576 MPa. After 3.04×10^{-4} s, the symbol of the minimum principal stress value changes,

and tensile stress is generated at the monitoring point under explosion stress waves. (Figure 4).



Table 3. Calculation Model II.

Figure 3. Intact frozen rock mass explosion stress wave propagation and failure process.



Figure 4. The minimum principal stress during blasting at intact frozen rock monitoring point.

3.3. *Ice-Filled Crack Thickness d Influence on Blasting Effect of Frozen Rock Mass* 3.3.1. Blasting Failure Process

By comparing the blasting effect of frozen rock with ice-filled cracks of 0.02 m, 0.04 m and 0.08 m thickness distributed in the same normal distance, the following observations

are obtained. With the propagation of explosion stress waves, the numerical simulation results are the same as the blasting effect of intact frozen rock mass, and the blasting crack appears first at the blasting hole (Figure 5). Taking d = 0.02 m and R = 0.2 m frozen rock mass with ice-filled cracks as an example, at 0.54×10^{-4} s, damage elements occur near the blasting hole. At 1.48×10^{-4} s, the stress wave reaches the ice-filled crack. When the explosion stress wave propagates to the ice-filled crack, reflects and transmits. The reflected wave collides with the incident wave, which reduces both the energy of the incident wave and the explosion range. As the mechanical strength of ice is less than that of rock, damage elements gradually appear in the ice-filled crack, and the range of damage elements intensifies. At 5×10^{-4} s, the stress wave passed through the ice-filled crack, and the damage range remains stable.



Figure 5. Explosion stress wave propagation and failure process of frozen rock mass with different ice-filled crack thicknesses (the normal distance R is 0.2 m). (**a1–a4**) d = 0.02 m; (**b1–b4**) d = 0.04 m; (**c1–c4**) d = 0.08 m.

By comparing Figure 5(a4,b4,c4), the damage elements range of the ice-filled crack decreases when the thickness of the ice-filled crack increases. By comparing the crack propagation of the non-ice side with the ice side, the crack propagation of the blasting is induced by the ice-filled crack.

3.3.2. The Minimum Principal Stress at the Monitoring Point A with Different Ice-Filled Crack Thicknesses

By comparing the minimum principal stress at the monitoring point A of ice-filled cracks with different thicknesses (Figure 6), the minimum principal stress presents a positive

increase at first, then decreases to a negative fluctuation. The minimum principal stress amplitude is similar to that of the intact frozen rock mass. Taking R = 0.8 m frozen rock mass with ice-filled crack as an example.



Figure 6. The minimum principal stress during blasting at frozen rock mass monitoring point with different ice-filled crack thicknesses. (a) R = 0.2 m; (b) R = 0.4 m; (c) R = 0.8 m.

When d = 0.02 m, the minimum principal stress is positive and the peak value is 0.528 MPa from 2.90×10^{-4} s to 3.72×10^{-4} s. The explosion stress waves produce extrusion effects on the monitoring point. After 3.72×10^{-4} s, the stress wave propagates through the monitoring point, tensile stress is generated on the monitoring point, and the minimum principal stress turns negative. When d = 0.04 m, the peak value of the minimum principal stress is 0.470 MPa from 2.92×10^{-4} s to 3.78×10^{-4} s. When d = 0.08 m, the peak value of the minimum principal stress is 0.395 MPa from 2.96×10^{-4} s to 3.76×10^{-4} s. The results show that the minimum principal stress decreases with an increase of ice-filled crack thickness.

At the monitoring point A, the peak value of the minimum principal stress of intact frozen rock mass is 0.576 MPa. The minimum principal stress peak value of frozen rock mass with ice-filled cracks is obviously smaller than that of intact frozen rock mass. It can be seen that the existence of ice-filled cracks inhibits the propagation of explosion stress waves in frozen rock mass.

3.4. Normal Distance R Influence on Blasting Effect of Frozen Rock Mass 3.4.1. Blasting Failure Process

Figure 7 shows the explosion stress wave propagation and failure process of frozen rock mass with different normal distance R. The crack pattern at the blasting hole is similar to that of the intact frozen rock mass, showing a uniform and divergent extension. When

d = 0.02 m and 0.04 m, the damage elements ranges have little differences. When d = 0.08 m, the damage elements range have obvious differences. The smaller the normal distance R is, the closer the ice-filled crack is to the blasting hole, the more obvious the explosion stress wave effect is. The thicker the ice-filled crack is, the more apparent that it is affected by normal distance.



Figure 7. Explosion stress wave propagation and failure process of frozen rock mass with different normal distance R. (**a**) d = 0.02 m; (**b**) d = 0.04 m; (**c**) d = 0.08 m.

3.4.2. The Minimum Principal Stress at the Monitoring Point A with Different Normal Distances

Taking d = 0.04 m frozen rock mass with ice-filled crack as an example (Figure 6). When R = 0.2 m, the peak value of the minimum principal stress is 0.472 MPa from 2.98×10^{-4} s to 3.82×10^{-4} s. The explosion stress wave produces extrusion effects on the monitoring point. After 3.72×10^{-4} s, as the stress wave propagates through the monitoring point, tensile stress is generated on the monitoring point, so it is a negative value. When R = 0.4 m, the peak value of the minimum principal stress is 0.472 MPa from 2.98×10^{-4} s to 3.82×10^{-4} s. When R = 0.8 m, the peak value of the minimum principal stress is 0.470 MPa from 2.98×10^{-4} s to 3.82×10^{-4} s to 3.80×10^{-4} s. When the normal distance between the ice-filled crack and the blasting hole increases, the reduction of the explosion stress wave propagation will decrease slightly.

3.5. *Ice-Filled Crack Angle α Influence on Blasting Effect of Frozen Rock Mass* 3.5.1. Blasting Failure Process

Figure 8 shows the explosion stress wave propagation and failure process of frozen rock mass with different ice-filled crack angles. Different angles affect the distance from the ice-filled crack to the blasting hole. When d = 0.04 m and R = 0.8 m, the larger the ice-filled crack angle is, the closer the ice-filled crack is to the blasting hole and the wider the damage range of the ice-filled crack.



Figure 8. Damage and cracking after blasting of frozen rock mass with different angles. (a) 15° ; (b) 30° ; (c) 45° ; (d) 60° ; (e) 75° .

3.5.2. The Minimum Principal Stress at the Monitoring Point A with Different Loading Angles

Taking the frozen rock mass blasting process with the ice-filled crack angle 15° as an example. Before 2.92×10^{-4} s, the minimum principal stress does not change. From 2.92×10^{-4} s to 3.80×10^{-4} s, the peak value of the minimum principal stress is 0.462 MPa. After 3.80×10^{-4} s, the minimum principal stress turns negative. When the ice-filled crack angle is 30° , 45° , 60° , 75° , the peak value of the minimum principal stress is 0.464 MPa, 0.482 MPa, 0.475 MPa and 0.500 MPa, respectively (Figure 9). An increase of the ice-filled crack angle causes a slight decrease of the minimum principal stress peak value.

3.6. Loading Intensity Influence on Blasting Effect of Frozen Rock Mass 3.6.1. Blasting Failure Process

The mechanical properties of rock materials are affected by the loading peak value, which might influence the explosion stress wave propagation efficiency and blasting effect. Taking d = 0.04 m and R = 0.4 m frozen rock mass as an example to simulate the blasting process with different loading peak values. The failure patterns in loading cases I, II and III at selected times were compared in Figure 10.

With the propagation of the explosion stress wave, the blasting crack appears first at the blasting hole. At 0.54×10^{-4} s, the blasting hole begins to show damage elements. At 1.48×10^{-4} s, damage elements gradually appear in the ice-filled crack first. The blasting crack pattern at the blasting hole is obviously different due to different loading intensities.

At 3.00×10^{-4} s, the explosion stress wave continues to impact the ice-filled crack, and the range of damage elements intensifies. At 5×10^{-4} s, the stress waves have passed through the ice-filled crack and the damage range remains stable. By comparison with Figure 10(a4,b4,c4), as the loading intensity increases, the blasting crack pattern at the blasting hole becomes more apparent, and the damage elements range expands. Also, the effect of explosion stress wave propagation is different.









3.6.2. The Minimum Principal Stress at the Monitoring Point A with Different Loading Intensities

Taking the results of frozen rock mass with ice-filled crack blasting process when the loading intensity of the explosion stress wave is 15 MPa as an example. Before 2.96×10^{-4} s, the minimum principal stress has no change. From 2.96×10^{-4} s to 3.80×10^{-4} s, the peak value of the minimum principal stress is 0.362 MPa, and the explosion stress wave produces extrusion effects on the monitoring point. After 3.80×10^{-4} s, the minimum principal stress is negative. As the explosion stress wave propagates through the monitoring point, tensile stress is generated on the monitoring point. When the explosion stress wave loading intensity is 10 MPa, the minimum principal stress peak value is 0.233 MPa. When the explosion stress wave loading intensity is 20 MPa, the minimum principal stress peak value is 0.472 MPa (Figure 11). When the loading peak value of the explosion stress wave is larger, the effect of the explosion stress wave on the rock medium is more obvious.



Figure 11. The minimum principal stress during blasting at frozen rock mass monitoring point with different loading intensities.

3.7. Loading Rate Influence on Blasting Effect of Frozen Rock Mass

3.7.1. Blasting Failure Process

Different blasting loading rates also have influence on the blasting effect. Taking d = 0.08 m and R = 0.4 m frozen rock mass with ice-filled crack as an example, the failure modes of loading cases I, IV and V at selected times are compared, as shown in Figure 12.

At 0.54×10^{-4} s, the blasting hole begins to produce damage elements. At 1.48×10^{-4} s, damage elements gradually appear in the ice-filled crack first. The blasting crack pattern at the blasting hole is different due to the different loading rates. At 3.00×10^{-4} s, the explosion stress wave continues to impact the ice-filled crack, and the range of damage elements intensifies. The blasting hole crack continues to expand. At 5×10^{-4} s, the stress waves have passed through the ice-filled crack and the damage range remains stable. By comparison with Figure 12(a4,b4,c4), as the loading rate decreases, the crack length at the blasting hole increases, and the ice-filled crack shows more damage elements.

3.7.2. The Minimum Principal Stress at the Monitoring Point A with Different Loading Rates

By comparing the minimum principal stress on the monitoring point with different loading rates in Figure 13, the following results are observed. Before 2.96×10^{-4} s, the minimum principal stress stays unchanged. From 2.96×10^{-4} s to 3.80×10^{-4} s, the peak value of the minimum principal stress is 0.362 MPa in loading case I. In loading cases IV and V, the variation of the minimum principal stress at the monitoring point A fluctuates dramatically. As the loading waveform is different, the minimum principal stress appears hysteresis, and the transmitted waves generated by the explosion stress waves through the ice-filled cracks are also different.



Figure 12. Damage and cracking after blasting of frozen rock mass with different loading rates. (**a1-a4**) Loading case I; (**b1-b4**) Loading case IV; (**c1-c4**) Loading case V.



Figure 13. The minimum principal stress during blasting at ice-filled crack frozen rock mass monitoring point with different peak loading values.

3.8. Evaluation of the Explosion Stress Wave Attenuation Factor

In order to intuitively evaluate the effects of different ice-filled crack conditions on the explosion stress wave, the attenuation factor B of the explosion stress wave through frozen rock mass is calculated as follows,

$$\mathbf{B} = \left| \frac{\beta_c - \beta_i}{\beta_i} \right| \times 100\% \tag{13}$$

where β_c is the peak value of the minimum principal stress at monitoring point of frozen rock mass under different conditions; and β_i is the minimum principal stress value at monitoring points of intact frozen rock mass, which is 0.576 MPa.

The attenuation trend of the explosion stress wave decreases gradually with an increase of thickness, and decreases slightly with an increase of direction, distance and ice-filled crack angle (Table 4).

 Table 4. Explosion stress wave attenuation at monitoring points of frozen rock mass under different conditions.

No.	Width (m)	R (m)	Angle (°)	Loading Case	Attenuation (%)
1	0.02	0.2	0	Ι	7.128
2	0.02	0.4	0	Ι	7.292
3	0.02	0.8	0	Ι	8.333
4	0.04	0.2	0	Ι	18.056
5	0.04	0.4	0	Ι	18.056
6	0.04	0.8	0	Ι	18.403
7	0.08	0.2	0	Ι	30.035
8	0.08	0.4	0	Ι	29.689
9	0.08	0.8	0	Ι	31.424
10	0.04	0.8	15	Ι	19.792
11	0.04	0.8	30	Ι	19.444
12	0.04	0.8	45	Ι	16.319
13	0.04	0.8	60	Ι	17.535
14	0.04	0.8	75	Ι	13.194

4. Conclusions

In this work, the stress wave propagation and blasting crack extension mechanism of frozen rock mass with ice-filled cracks are analyzed by RFPA^{2D}. The findings facilitate blasting design and disaster prevention in cold region strip mining. Specifically, the explosion stress wave propagation effect and the damage range are obtained, which consider the geometrical parameters of ice-filled cracks (ice-filled crack thickness D, normal distance R from blasting hole to ice-filled crack, and ice-filled crack angle α), loading intensity and loading rate. The following conclusions can be drawn:

- (1) The divergent crack extension pattern is always maintained at the blasting hole. Affected by explosion wave stress, the damage elements range of the ice-filled crack decreases when the ice-filled crack thickness and normal distance increase.
- (2) The attenuation trend of explosion stress waves decreases with an increase of icefilled crack thickness, and decreases slightly with an increase of normal distance and ice-filled crack angle.
- (3) The loading intensity and the loading rate have a significant influence on blasting hole fracture patterns. The damage elements range of the ice-filled crack is enlarged when the loading intensity and the loading rate increase. The propagation effect of explosion stress waves is also different.
- (4) The ice-filled crack has a guiding effect on the growth of the blasting crack at the blasting hole. Nevertheless, the existence of ice-filled cracks inhibits the propagation of explosion stress waves in frozen rock mass.

It should be noted that due to the lack of access to information of a real project, in this research, only numerical modelling has been done. It is better to combine these results with experiment and field validation to achieve a more comprehensive method to study frozen rock masses blasting process.

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