



Article

Radiation Damage in Polyethylene Naphthalate Thin-Film Scintillators

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Abstract: This paper describes the scintillation features and the radiation damage in polyethylene naphthalate 100 μm -thick scintillators irradiated with an 11 MeV proton beam and with a 1 MeV electron beam at doses up to 15 and 85 Mrad, respectively. The scintillator emission spectrum, optical transmission, light yield loss, and scintillation pulse decay times were investigated before and after the irradiation. A deep blue emission spectrum peaked at 422 nm, and fast and slow scintillation decay time constants of the order of 1–2 ns and 25–30 ns, respectively, were measured. After irradiation, transmittance showed a loss of transparency for wavelengths between 380 and 420 nm, and a light yield reduction of $\sim 40\%$ was measured at the maximum dose of 85 Mrad.

Keywords: plastic scintillators; radiation damage; polyethylene-naphthalate; thin-film scintillator



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1. Introduction

Plastic scintillators are well-established radiation detectors in several applications, including particle and nuclear physics experiments, as well as medical physics. They feature a fast decay time, high light yield (LY) and optical transmission, ease of manufacture and low cost [1–3]. Plastic scintillators are typically made of a polymer matrix with small quantities of primary and secondary dyes [4]. When the polymer base molecules are excited by the radiation, the excitation energy is promptly transferred to the primary dye via the Foerster mechanism [5], and thus to the secondary dye (also referred to as wavelength shifter) that emits photons in a wavelength range detectable from most of the typically used photodetectors. Typical plastic scintillators’ polymer bases are aromatic hydrocarbon chemical compounds such as polystyrene (PS) or polyvinyltoluene (PVT). These are doped with fluorescent molecules such as 2,5-Diphenyloxazole (PPO) and 1,4-bis(5-phenyloxazol-2-yl) benzene (POPOP), typically in percentages of $\sim 1\%$ and $\sim 0.005\%$, respectively [6,7].

Recent works proved that undoped aromatic ring polymers commonly known as low-cost plastic materials, such as polyethylene naphthalate (PEN) or polyethylene terephthalate (PET), are good scintillation materials [8–11]. PEN is synthesized by the poly-condensation of dimethyl-2,6-naphthalenedicarboxylate and ethylene glycol. Its monomeric chemical structure consists of a double benzene ring, as shown in Figure 1.

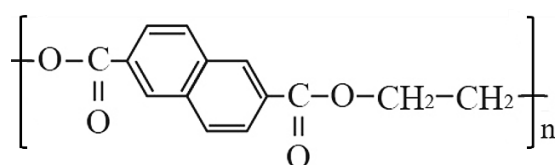


Figure 1. Chemical formula of the polyethylene naphthalate.

PEN features blue light emission with a peak at ~ 420 nm and a high LY of $\sim 10^4$ photons/MeV [9], a short decay time [12,13], and a moderately high density (1.33 g/cm^3) with respect to other common plastic scintillators, despite a relatively high refractive index [14]. Moreover, PEN has good mechanical properties that allow its production in different sizes and shapes, including thin flexible films, which could be useful in applications where curved surfaces must be instrumented. Because of these properties and ease of manufacture, PEN has recently drawn the attention of the scientific community for its use as a scintillator. As an example, some works explored its suitability for dosimetry in nuclear and medical physics [15–17] and calorimetry in collider experiments [18]. PEN is also suitable for neutron detection due to the high amount of hydrogen, which scatters neutrons through elastic scattering on protons [19]. PEN has also been proposed as a wavelength shifter for noble liquid scintillators in astroparticle underground experiments, as it fulfills the requirements for high purity and low background contamination [20–22].

Some of these applications require the detectors to operate in environments with high radiation levels while asking for long-term stability. Specifically, accelerator or collider applications demand resistance to extremely high radiation doses. As an example, the radiation background in the hadronic calorimeter of the CMS detector at the Large Hadron Collider (LHC) results in a 10 Mrad expected dose after 10 years of operation [23]. Ongoing upgrades of existing high-energy physics experiments and future particle colliders will deal with even higher radiation dose levels [24]. During operation in such harsh environments, the interaction of the radiation with the scintillation detectors may result in the change of their physical and chemical properties through several mechanisms depending on the selected material. Possible effects include polymer chain scission, breakage of bonds, cross-linking and the formation of new chemical bonds [23]. These can affect either the scintillation light mechanisms or the light transmission, resulting in a deterioration of the detector performances [25–31]. Therefore, the search for new scintillation materials capable of withstanding high radiation doses is of paramount importance.

In this work, we describe the scintillation features and the radiation damage in a $100 \text{ }\mu\text{m}$ -thick polyethylene naphthalate scintillator. We irradiated several samples of such a scintillator with an 11 MeV proton beam and with a 1 MeV electron beam at the maximum doses of 15 Mrad and 85 Mrad, respectively. We characterized the PEN emission spectrum, optical transmission, scintillation decay time and LY loss before and after the irradiation.

Our study adds to the few other works on radiation damage in PEN [32–35], and explores radiation effects on several PEN properties over a broader range of doses, with two radiation types (electrons and protons) at two different energies, paving the way for using such thin-film scintillators in applications where high levels of radiation are expected.

2. Materials and Methods

The PEN scintillator under investigation is produced by Teonex[®] as a $100 \text{ }\mu\text{m}$ -thick flexible film (mod. Q65HA). Starting from the commercialized format, we obtained, by cutting, samples of $30 \times 40 \text{ mm}^2$, as shown in Figure 2.

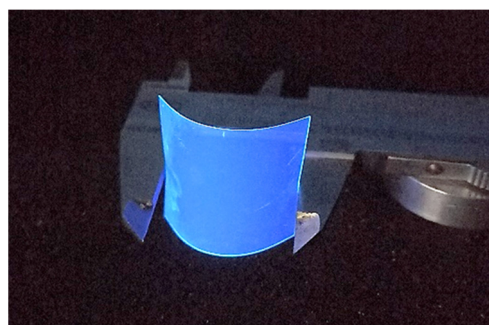


Figure 2. A $30 \times 40 \times 0.1 \text{ mm}^3$ thin-film flexible PEN scintillator illuminated with UV light.

Two irradiation campaigns were performed, one with a proton beam and the other with an electron beam. In both cases, PEN samples were irradiated in air (at room temperature), placed at the exit of the beam-line and exposed frontally to the beam. The proton irradiation was carried out at the Tandem accelerator of the INFN Laboratori Nazionali del Sud. Considering the energy loss in the beam extraction window and in the air layer between the end of the beam line and the sample, the proton energy on the sample surface was estimated to be 11 MeV. The PEN samples were exposed to a maximum proton dose of 15 Mrad (to water), with a max dose rate of 21 krad/min. The electron irradiation was performed at the ILU-6 accelerator at the Institute of Nuclear Chemistry and Technology (INCT) in Warsaw (Poland) with a 1 MeV electron beam. The samples were exposed to a maximum dose of 85 Mrad (to water) with a max dose rate of 5 Mrad/min. In both irradiations, the beams had enough energy to pass through the whole scintillator: the expected penetration ranges in polyethylene at those energies are ~ 1.1 mm and ~ 3.5 mm for protons and electrons, respectively.

The emission spectrum was measured with a spectrometer (mod. AvaSpac-2048XL from Avantes, Apeldoorn, The Netherlands) by stimulating PEN emission with a 280 nm LED source operating at 20 mW. The optical transmittance in the 300–800 nm wavelength range was measured with the same spectrometer by using a deuterium/halogen light source.

To study the scintillation pulse decay times and the LY, we exposed our PEN samples to both ^{241}Am and ^{137}Cs sources. The former decays to the second excited state of ^{237}Np with 85% absolute intensity by emitting 5.486 MeV alpha particles; the latter decays to ^{m137}Ba with 95% absolute intensity by emitting electrons with maximum energy of 0.514 MeV. Scintillation pulses were measured by coupling the PEN scintillator to a flat window Hamamatsu R5900 PhotoMultiplier Tube (PMT) with a gain of $\sim 4 \times 10^6$ biased at 900 V. The radioactive source was positioned on top of a collimator with a 2 mm hole diameter. A schematic of the experimental setup is shown in Figure 3.

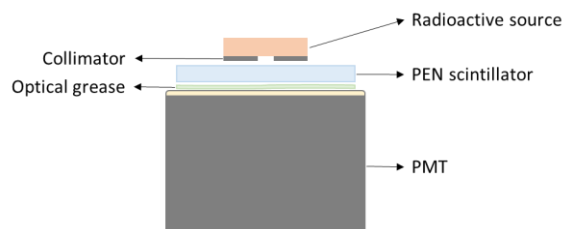


Figure 3. Schematic representation of the experimental setup for the PEN thin-film scintillators light yield measurement.

The scintillation pulse decay times were measured by acquiring the signal waveform directly from the PMT with a 1 GHz-bandwidth, 10 GS/s sampling rate LeCroy oscilloscope. As for the measurement of the signal spectra and of the LY, PMT signals were amplified and acquired via an ORTEC EASY-MCA Multichannel Analyzer (MCA).

3. Results

3.1. Preliminary Characterization

We tested the scintillation features of our 100 μm -thick film PEN samples prior to the irradiation. The samples were stimulated with a 285 nm LED and the emission spectrum was measured, as shown in Figure 4 (left). A single peak feature at 422 nm was observed, in agreement with the values reported in the literature [9]. Such a deep blue emission spectrum matches the high-quantum-efficiency spectral regions of many commercially available photomultipliers or SiPMs [36,37]. The optical transmittance, relative to air, was measured in the range from 300 to 800 nm, as shown in Figure 4 (Right). Transmittance exhibits a steep decrease for wavelengths lower than about 400 nm and is negligible below 370 nm.

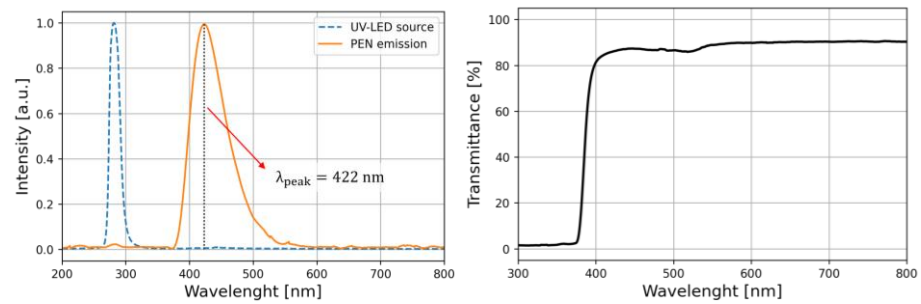


Figure 4. (Left): Emission spectrum of the PEN scintillator (orange) stimulated by a UV LED source (blue dashed curve). (Right): PEN transmittance relative to air as a function of the wavelength.

Scintillation decay time constants were measured by exciting the scintillator with alpha and beta particles from ^{241}Am and ^{137}Cs sources and coupling it to the PMT according to the description given in Section 2. The scintillation light pulses over 10,000 waveforms averages are shown in Figure 5.

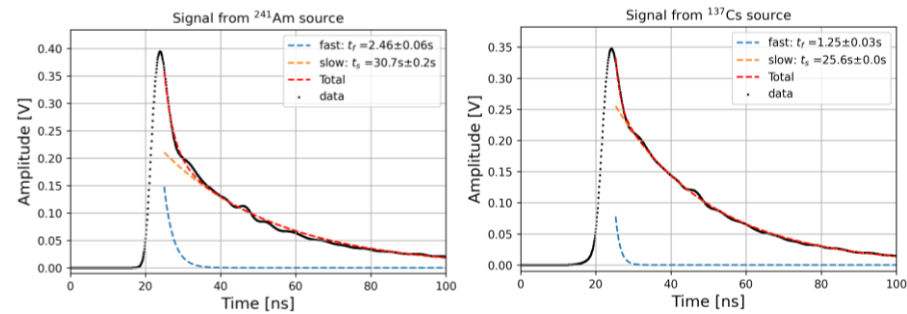


Figure 5. Signal waveform of PEN scintillators excited with ^{241}Am (left) and ^{137}Cs (right) sources. The signals shown are the averages of 10,000 waveforms.

The signal decay shape was modeled as the sum of two negative exponential functions with time constants τ_{fast} and τ_{slow} , respectively, which account for different de-excitation mechanisms [3]:

$$\text{Amplitude}(t) = \frac{A}{\tau_{fast}} e^{-\frac{t}{\tau_{fast}}} + \frac{B}{\tau_{slow}} e^{-\frac{t}{\tau_{slow}}} \quad (1)$$

By fitting the experimental data for a pristine PEN sample with Equation (1), we observed fast and slow time constants of 2.46 ns and 30.7 ns for the sample excited with ^{241}Am and of 1.25 ns and 25.6 ns for the sample excited with ^{137}Cs . Similar values have been reported in the literature for the slow decay time constants [12,13], while the fast component had not been measured before. Such small time constants strongly increase the interest in using such materials in many applications requiring good timing resolution.

3.2. Post-Irradiation Measurements

After the irradiation, PEN samples were kept in the dark with a controlled environment temperature of 25 °C. All post-irradiation measurements were performed one month after the irradiation campaigns in order to measure the persistent damage effects only. Indeed, a partial recovery mechanism of the radiation-induced damage in PEN with a characteristic time of ~10 days is reported in the literature [32].

Differences in the emission spectrum before and after the irradiation are below 0.5%. This is in agreement with the results obtained with gamma-irradiated PEN scintillators reported in Ref. [33], where only a shift of a few nanometers of the emission peak was observed for a dose of 65 Mrad.

Similarly, no degradation trend in the decay times was observed as a function of the dose, with variations within 7%. This suggests that the damage to the light emission mechanisms is negligible up to the considered doses.

The light yield spectra of the irradiated samples are shown in Figure 6. The left plots show the spectra measured by using the ^{241}Am source. A typical distribution of a monochromatic particle losing all the energy in the detector was observed; indeed, 5.486 MeV alpha particles from ^{241}Am have a range of $\sim 35\ \mu\text{m}$ in PEN. In the right plots, spectra measured using ^{137}Cs source are shown. The observed distribution is due to the energy loss of beta particles with a continuum energy spectrum, which accounts for the observed long tail.

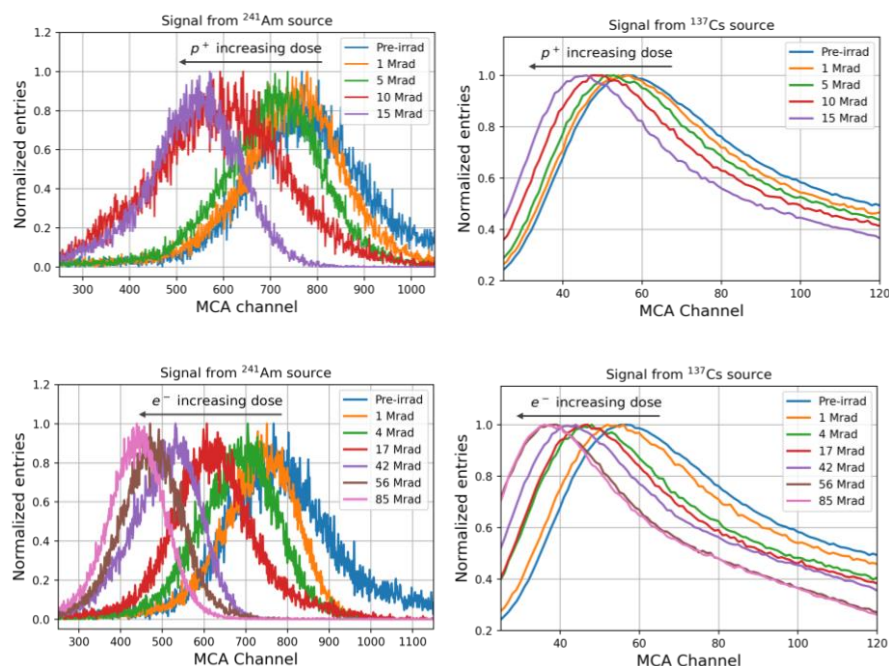


Figure 6. Light yield spectra of PEN scintillators irradiated with protons (**top**) and electrons (**bottom**), measured upon excitation with ^{241}Am (**left**) and ^{137}Cs (**right**) sources.

The effect of the irradiation is observable as a decrease in the peak centroids as the radiation dose increases.

In Figure 7, the irradiated samples LY normalized to that of a pristine scintillator as a function of the delivered dose are shown. Similar trends for both proton and electron irradiated samples can be observed as long as the same dose range is investigated. A slightly larger LY degradation is observed when measured with ^{241}Am than with ^{137}Cs . This could be ascribed to the fact that alpha particle-induced scintillation occurs near the surface of PEN samples, and thus the scintillation light is more attenuated than that due to ^{137}Cs beta particles, which excite the sample along with its entire thickness. At higher doses, the radiation damage rate slows down and LY loss reaches values close to 40%.

In Ref. [32], the authors find, for a 1 mm-thick PEN scintillator irradiated with gamma at a dose of 14 Mrad, a permanent $\sim 20\%$ LY loss, which is close to our findings at similar doses. A larger degradation is observed in Ref. [34], where after irradiation of a $9\ \mu\text{m}$ -thick PEN film with 1 MeV protons, a $\sim 55\%$ loss was measured at a dose of ~ 150 Mrad after 3 h from the irradiation.

Optical transmission measurements were repeated after irradiation as a function of the delivered dose. The results are shown in Figure 8. As the dose increases, the absorption edge shifts to longer wavelengths and an overall decrease in the transmittance is observed.

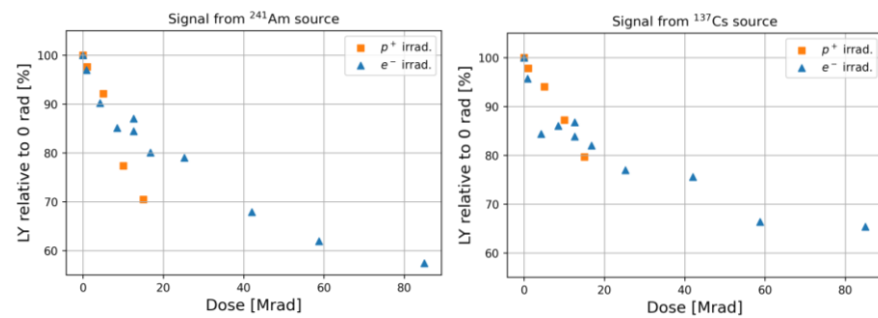


Figure 7. Irradiated PEN light yield normalized to that of a pristine scintillator as a function of the delivered dose for proton and electron irradiations. Scintillation light was measured upon excitation with (left) ^{241}Am and (right) ^{137}Cs sources.

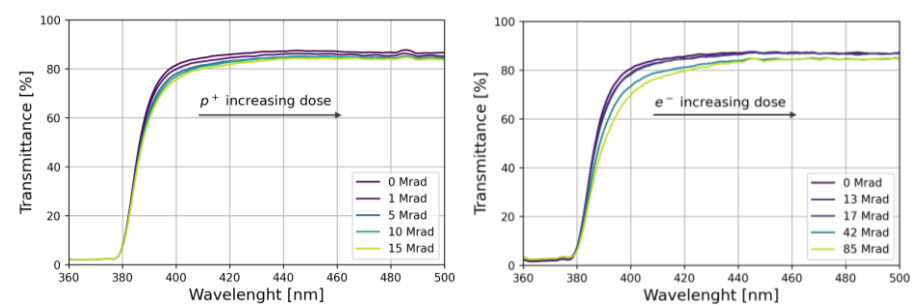


Figure 8. Light transmission spectra for PEN scintillators irradiated with protons (left) and electrons (right).

We computed the radiation-induced absorption coefficient μ as a function of the wavelength as:

$$\mu = \ln\left(\frac{T_b}{T_a}\right)/d, \quad (2)$$

where T_b and T_a are the transmission measured before and after irradiation, respectively, and d is the sample thickness ($d = 100 \mu\text{m}$). The results are shown in Figure 9, where large absorption values are observed in the 380–420 nm-wavelength region, suggesting that the effect of irradiation is the production of free radicals acting as absorption centers [38].

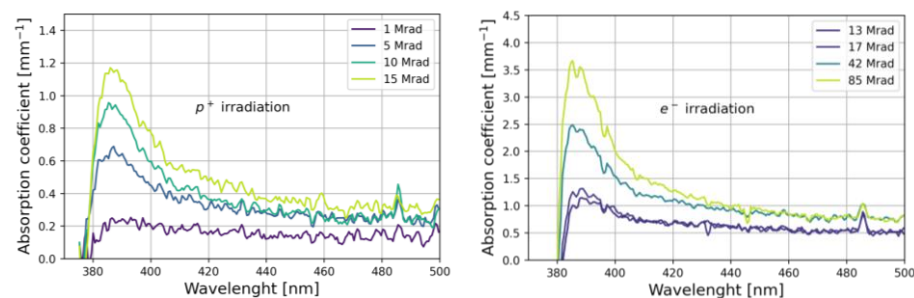


Figure 9. Radiation-induced absorption coefficient spectra for samples irradiated with protons (left) and electrons (right).

4. Conclusions

In recent years, PEN has emerged as a good scintillation material due to its behaviors such as high light yield and blue emission, which match well with the sensitivity range of many common photodetectors. Furthermore, PEN is a common and cheap material and can be produced in different sizes and shapes, including in the form of flexible films.

In this work, we report on the characterization of the scintillation features and the radiation damage in polyethylene naphthalate $100 \mu\text{m}$ -thick scintillators.

Our results revealed the excellent PEN scintillation features: a deep blue emission spectrum peaked at 422 nm and small fast and slow decay time constants of the order of 1–2 ns and 25–30 ns, respectively.

The radiation damage in PEN has been investigated by performing two irradiation tests with 11 MeV protons and 1 MeV electrons with maximum doses of 15 Mrad and 85 Mrad, respectively. Irradiation did not change the light emission spectrum and nor the scintillation decay time constants. A reduction of ~40% in the light yield was observed at the maximum delivered dose of 85 Mrad. Transmittance measurements showed a marked loss of transparency for wavelengths between 380 and 420 nm.

Based on the findings reported in this work, we foresee the possibility of using such PEN thin-film scintillators in radiation detection applications where peculiar geometries and high doses are expected.

Future works will be devoted to detailed studies of the PEN damage for different high-dose radiation types, which can be of paramount importance for applications involving the presence of more than one radiation type.

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