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The Influence of the Thickness of Compact TiO₂ Electron Transport Layer on the Performance of Planar CH₃NH₃PbI₃ Perovskite Solar Cells

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Abstract: In recent years, lead halide perovskites have attracted considerable attention from the scientific community due to their exceptional properties and fast-growing enhancement for solar energy harvesting efficiency. One of the fundamental aspects of the architecture of perovskite-based solar cells (PSCs) is the electron transport layer (ETL), which also acts as a barrier for holes. In this work, the influence of compact TiO₂ ETL on the performance of planar heterojunction solar cells based on CH₃NH₃PbI₃ perovskite was investigated. ETLs were deposited on fluorine-doped tin oxide (FTO) substrates from a titanium diisopropoxide bis(acetylacetonate) precursor solution using the spin-coating method with changing precursor concentration and centrifugation speed. It was found that the thickness and continuity of ETLs, investigated between 0 and 124 nm, strongly affect the photovoltaic performance of PSCs, in particular short-circuit current density (J_{SC}). Optical and topographic properties of the compact TiO₂ layers were investigated as well.

Keywords: titanium dioxide; thin film; perovskite solar cells; electron transport layer

1. Introduction

One of the most reasonable responses for growing global electricity demands, respecting the need to protect the natural environment, is harnessing solar energy by means of photovoltaic devices. Besides the commonly used mature wafer-based silicon technology, the new branch of perovskite-based solar cells is very promising due to the high performance and potentially low costs of production. Since the first report in 2009 by Miyasaka's group [1], the efficiency has advanced from 3.8% to more than 25% [2]. Perovskite solar cell (PSC) technology evolved at first from titania-based dye-sensitized solar cells [3] and spread later into different architectures [4]: regular n-i-p with and without a mesoporous TiO₂ layer or inverted p–i–n (Figure 1). Although slightly higher efficiencies are obtained with mesoporous titania scaffolds [5], planar structures without them are more prospective from an industrial point of view. After it was proved that perovskites, unlike organic absorbers, generally have a long carrier diffusion length (a few hundreds of nanometers) [6], mesoporous TiO_2 was able to be omitted, allowing for the processing of whole solar cells below 200 °C. High efficiency of 21.6% for the planar n-i-p structure has been reported by Jiang et al. [7], while other authors have claimed an even higher value of 23.7% [8]. In general, the PSC structure is composed of a perovskite absorber layer placed between a selective transport layer for holes (HTL) and electrons (ETL). For the p–i–n structure, the most commonly used HTLs are PEDOT:PSS [9] and NiO_x [10], while the ETL, mostly PCBM [11], is placed on top of the HTL and perovskite. On the other hand, in the most common n-i-p architecture, titanium dioxide is predominantly used as an ETL followed by a perovskite absorber and Spiro-OMeTAD as the HTL. Although significant progress in PSCs was enabled by engineering the perovskite composition [12,13]



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and microstructures [14,15], achieving high power conversion efficiency is not possible without optimizing other cell components.

Figure 1. Schemes of the three most common architectures used in PSC technology, from left to right: planar regular (n–i–p), mesoporous (n–i–p) and planar inverted (p–i–n). The devices are composed of glass coated with transparent conductive oxide (TCO), an electron transport layer (ETL), a possibly mesoporous electron transport layer (mp-ETL), the perovskite, a hole transport layer (HTL) and the top electrode.

The ETL in PSCs transports electrons generated from the perovskite layer [16]. It also serves as a blocking layer to hinder direct contact between the holes and FTO [17]. Hence, the layer should be pinhole free to prevent the recombination of electrons and holes at the front electrodes. Compatible energy levels support the fast injection of photogenerated electrons from perovskites and low-voltage losses. It also acts as a window layer for the perovskites; thus, a bandgap above 3 eV is necessary. The layer should electrically be as thin as possible to provide fast electron transport and low resistive losses. The titaniumoxide-based ETL is also a crucial investigation area reported in a number of studies [18–22]. The titanium dioxide blocking layer can be deposited by various methods: atomic layer deposition (ALD) [23], magnetron sputtering [24], spray pyrolysis [25], electrochemical deposition [26] and spin coating, where the latter is most commonly used. Due to the relatively high surface roughness of the FTO electrode, typically ca. 20 nm (Rq), the deposition methods can be divided into methods supporting a perfectly conformal morphology, such as ALD or sputtering, and methods resulting in fewer conformal layers such as spin coating. Spin-coated compact TiO₂ layers more easily fill the troughs of the rough surface than the peaks. The optimal thickness of $c-TiO_2$ is a minimal one that supports the blocking properties at the peaks. Literature reports have announced a variety of titanium precursors and deposition conditions; however, the optimal thickness is rather rarely given explicitly. Additionally, each precursor may require different thicknesses to reach the highest cell performances, as discussed in [27], the reason being that there are intrinsic differences in material properties in comparison to the morphological aspects. For instance, Qin et al. compared three different titanium precursors: tetrabutyl titanate, titanium diisopropoxide bis(acetylacetonate), titanium isopropoxide and solvents [28], concluding the first as the best. However, the thicknesses of compared layers differed considerably. Furthermore, the confrontation of ETLs was conducted in the presence of mesoporous TiO_2 , which could slightly disrupt the thickness effects, and the surface roughness of FTO used in these studies was approximately two times lower than commonly used.

It is known that ETLs (TiO₂ or SnO₂) should be uniform, pinhole free and completely cover the surface of FTO. Therefore, the optimal thickness depends on the deposition technique and post-treatment of these layers. Many authors have reported thin ETLs in the range of 20–30 nm. For example, Xiao et al. [29] applied a 30 nm thick layer of TiO₂ for the planar n–i–p structure and obtained a power conversion efficiency of 13.9%. Saliba et al. [4] used 20–30 nm thick SnO₂ (instead of TiO₂) in the planar regular structure (n–i–p) and a 20–30 nm compact TiO₂ layer for the mesostructure for high-efficiency solar cells with a PCE \geq 20%. On the other hand, Sun et al. [30] obtained the greatest efficiency (PCE of 18.32%) for much thicker layers of 152 nm. In this case, the TiO₂ layer was additionally treated with TiCl₄. However, for the ALD layer deposition technique, the thickness of the layers could be much lower. Lu et al. [31] showed that the optimal thickness for the ALD technique is considerably smaller and equal to 10 nm (PCE of 13.6%). This suggests that the ETL may be very thin but must also be of good quality.

Considering the ETL is an important part of PSCs, in this work, we emphasize its impact on perovskite solar cell parameters. The regular planar n–i–p architecture of PCSs was studied, i.e., glass/FTO/c-TiO₂/MAPbI₃/Spiro-OMeTAD/Au. In this report, we deposited ETLs using a titanium diisopropoxide bis(acetylacetonate) (Ti(acac)₂OiPr₂) precursor solution in 1-BuOH. This set was chosen because its good performance matches the preparation simplicity and good storage stability. This stability is important in the case of experimental repeatability, as it was already reported for precursor solutions undergoing aging processes [27,32]. Thoroughly differing thicknesses of the compact TiO₂ layer were the key features enabling careful investigation of the properties of c-TiO₂ and their influence on the performance of PSCs. Complex studies enabled a better understanding of the observed effects.

2. Materials and Methods

2.1. Preparation of Compact TiO₂ Layers

For the electron transport layer (ETM), we used compact TiO₂ (c-TiO₂) layers of different thicknesses. They were deposited either on square (2 cm \times 2 cm) FTO glass substrates (Sigma-Aldrich, St. Louis, MO, USA, ~7 Ω /sq) or on a (100) polished silicon wafer. FTO glass plates were washed by dipping in hot 2% (*v*/*v*) Hellmanex[®] III (Hellma Analytics, Müllheim, Germany) solution in deionized (DI) H₂O in an ultrasonic cleaner for 5 min. After that, they were dipped in hot DI H₂O and again in isopropanol in an ultrasonic cleaner for 5 min. Finally, the plates were submerged in DI H₂O and dried. Silicon wafers were cleaned successively with acetone and isopropanol and rinsed with deionized water. Just before applying compact TiO₂ layers, FTO glass plates and Si wafers were additionally treated with O₂ plasma (Harrick Plasma, Ithaca, NY, USA) for 15 min (MID).

The c-TiO₂ precursors were prepared by mixing 75 wt % Ti(acac)₂OiPr₂ in isopropanol (IPA) with 1-BuOH. Zero, 0.243, 0.486, 0.729, 1.214, 1.821 and 2.428 g of 75 wt % Ti(acac)₂OiPr₂ was added to a volumetric flask and filled to 5 mL with 1-BuOH to make a 0, 0.1, 0.2, 0.3, 0.5, 0.75 and 1 M (mole/dm³) solution. A 70 μ L volume of precursor solution was spread on cleaned FTO/Si plates and spin coated at 2000, 3000 or 4000 revolutions per minute

(RPM) for 15 s. They were then preheated at 200 °C for 10 min in an oven and calcined at 500 °C for 30 min in a tube furnace in the presence of air.

2.2. Fabrication of Solar Cells

Deposition of lead iodide (MAPbI₃) and the hole transport layer (HTL) was performed in a nitrogen-filled glove box (MBRAUN, Garching, Germany). A ca. 1 M methylammonium MAPbI₃ precursor was prepared by dissolving 0.461 g (1 mmol) of PbI₂ and 0.159 g (1 mmol) of MAI in 700 μ L of GBL and 300 μ L of DMSO at room temperature (RT) at ca. 28–30 °C. The solution was stirred with a magnetic stirrer for ca. 3 h and then filtered through a syringe filter of a 0.45 μ m membrane pore size. Prior to deposition of the precursor, FTO/c-TiO₂ substrates were cleaned in O₂ plasma for 30 min (MID). A 70 μ L volume of the MAPbI₃ precursor was spread on the FTO/c-TiO₂ plate by two-step spin coating: 1000 RPM (200 RPM/s ramp) for 10 s and immediately 5000 RPM (1333 RPM/s ramp) for 20 s. Between 16 and 18 s of 5000 RPM spinning, 200 μ L of toluene antisolvent was added. The spin-coating program is presented in Figure S1 of the Supporting Information. After deposition of MAPbI₃, each sample was annealed at 100 °C on a hot plate for 10 min and cooled down to RT.

To deposit the hole transport layer (HTL), we used a stock solution of doped Spiro-OMeTAD in CB. A 0.368 g (0.3 mol) amount of Spiro-OMeTAD was dissolved in 5 mL of CB and doped with 148 μ L (1.01 mol) of TBP and 83 μ L of LiTFSI (0.15 mol) stock solution (520 mg/mL in acetonitrile). A 70 μ L volume of the Spiro-OMeTAD solution was spread on the MAPbI₃ perovskite layer and spin coated at 4000 RPM for 30 s. After that, the samples were taken out of the glove box, masked and coated with gold by thermal evaporation. Au electrodes had a surface of 0.25 cm² and a thickness of approximately 120 nm.

2.3. Apparatus Used in Measurements

Photovoltaic performance measurements were carried out by I–V curve tracing using a Photo Emission Tech AAA class solar simulator under standard test conditions. The microstructural and surface investigations were performed using Innova multimode atomic force microscopy and scanning electron microscopy (tabletop TM3030, Hitachi High-Tech, Tokio, Japan). The optical characterization was based on ellipsometry (SE800 PV, SENTECH Instruments, Berlin, Germany, range = 300–980 nm, incident angle = 70°) and UV–vis–NIR spectroscopy (Lambda 950S, Perkin Elmer, Waltham, MA, USA, range = 300–850 nm).

3. Results and Discussion

3.1. Fabrication of c-TiO₂ Thin Layers by Spin Coating

Spin coating is a common technique used for the deposition of thin films on solid substrates. The main advantage of this process is the ability to quickly and easily produce very fine and uniform coatings in the thickness range of micrometer to nanometer [33]. The thickness of the layer depends on many factors and can be described by the model of Mayerhofer [34]:

$$d = \left(1 - \frac{\rho_0}{\rho}\right) \cdot \left(\frac{3\eta \cdot m}{2\rho_0 \omega^2}\right)^{\frac{1}{3}} \tag{1}$$

where *d* is the thickness, ρ is the density of the solution, ρ_0 is the density of the solvent, η is the viscosity of the solution, *m* is the rate of evaporation and ω is the angular speed. In practice, the thickness of films deposited from solutions is controlled by the spinning speed (related to angular velocity) and concentration of the desired substance (related to viscosity). In this work, we coated glass/FTO substrates using a solution of titanium diisopropoxide bis(acetylacetonate) (Ti(acac)₂OiPr₂) in 1-butanol (1-BuOH) and annealed it at 200 and 500 °C. A commercially available 75% solution of Ti(acac)₂OiPr₂ in isopropanol mixes with 1-BuOH in any ratio, forming clear mixtures that do not age quickly (no rapid hydrolysis process). In this work, we investigate a series of 19 devices with a thickness of c-TiO₂ layers ranging from 0 to 124 nm, which were deposited by the spin-coating method with varying centrifugation speed and concentration of the precursor. The layer

thickness was associated with the nominal thickness of polished Si determined from ellipsometry. Figure 2 shows the dependence of the thickness of the c-TiO₂ layer on the centrifugation speed for solutions of different concentrations. It can be seen that the precursor concentration has a major influence on the thickness of the c-TiO₂ film. Generally, the c-TiO₂ layer formed at 2000 RPM can be thinned by 13–21% using 3000 RPM or by 26–31% using 4000 RPM. The full set of ellipsometry data is available in Table S1 in the Supporting Information.



Figure 2. The dependence of the thickness of the c-TiO₂ layer on the spin speed (in 10^3 revolutions per minute) in the spin-coating process for different concentrations of Ti(acac)₂OiPr₂ in the 1-BuOH precursor.

Figure 1 shows schemes of the three most common architectures of PCSs: planar regular (n-i-p), mesoporous (n-i-p) and planar inverted (p-i-n) [35]. Because this publication focuses on the compact (or blocking) TiO₂ layer, we used the simple planar regular (n-i-p) architecture of the glass/FTO/c-TiO₂/MAPbI₃/Spiro-OMeTAD/Au configuration with no mesoporous scaffold. Herein, for the most part, the light reaches the perovskite absorber through the glass substrate, FTO and electron transport layer (ETL), c-TiO₂. Light nonabsorbed during the first pass ($\lambda > 600$ nm) is mostly backscattered from the gold electrode, as it results from its complex refractive index [36].

3.2. Optical Absorption

The power conversion efficiency (PCE) of PSC strongly depends on the photocurrent, which, in turn, is directly proportional to the number of absorbed photons. It is well known that MAPbI₃ perovskite is an excellent absorber in the 300-800 nm wavelength range [37]. The onset photon absorption energy of polar MAPbI₃ perovskite is about 1.5 eV (\approx 827 nm), which is close to its electronic bandgap value of 1.55 eV (\approx 800 nm) [38]. It shows a higher quantum yield in the blue/green region and a lower yield in the red/infrared [39,40]. For this reason, the ETL material should be as transparent as possible in a wide range of spectra. The thickness of the c-TiO₂ layer clearly affects the optical transmission of perovskite, which is shown in Figure 3. The used glass/FTO substrate absorbs or reflects some of the light, ca. 20% between 450 and 850 nm, and much more in the near-ultraviolet range (300-350 nm). c-TiO₂ layers less than ca. 30 nm hardly affect the transmittance of glass/FTO, except for some absorption below ca. 380 nm, which is characteristic of anatase [41]. The effective transmittance of the samples slowly decreases as the thickness of the c-TiO₂ layer increases (Figure 3b). It is worth noting that for all samples, one can observe multiple maxima and minima above the absorption edge in the transmittance spectra (ca. >360 nm), which are especially pronounced for c-TiO₂ layers thicker than 24.4 nm. This is due to the interference effect resulting from the two interface boundaries, c-TiO₂/FTO and the FTO/glass substrate, which suggests that the surface and interface of these bilayers

are rather optically smooth [42]. The transmittance data around the absorption band edge follow the clear trend of edge red shift with TiO_2 layer thickness. Thus, taking into account the bandgap of the anatase TiO_2 layer on FTO, it should be attributed to parasitic light absorption in TiO_2 . UV–vis–NIR spectra for all the studied samples are available in the Supporting Information (Figure S2).



Figure 3. (a) Transmittance of glass/FTO/c-TiO₂ of different thicknesses obtained from the Ti(acac)₂OiPr₂ solution in 1-BuOH; (b) the dependence of the effective transmittance on the thickness in glass/FTO/c-TiO₂. Effective transmittance T_{eff} was calculated for the range of wavelengths λ from 350 to 800 nm, where T is the transmittance, and N_{ph} is the radiant flux.

3.3. Atomic Force Microscopy

Figure 4 shows that the used glass/FTO substrate exhibits a roughness of ca. 21 nm that decreases linearly with the increasing thickness of c-TiO₂, down to ca. 9.5 nm for the 88.3 nm thick layer. The topography of the FTO surface shows sharp tops and mild depressions resulting from the columnar grain structure of this material. Based on the AFM profiles, it should be emphasized that even thin c-TiO₂ layers, e.g., 6.7 nm, dull the sharp peaks effectively, while the thicker layers smooth the entire surface. Due to the right viscosity of the 1-butanol-based titanium diisopropoxide bis(acetylacetonate) sol, the layer shows more conformal growth than other sols such as ethanol-based titanium ethoxide [27]. So far, there is a significant divergence of optimal thickness according to the deposition method: ~10–15 nm for ALD and 50–60 nm for spin coating. The enhanced conformality is the key feature that allows the thinning of the blocking layer. Although the roughness of TiO₂ could also affect the crystallization of MAPbI₃, SEM images show a similar microstructure of perovskite films deposited on the substrates with different c-TiO₂ thicknesses (Figure S3).

3.4. Current–Voltage (I-V) Characteristics of Solar Cells

The electrical properties of solar cells have a direct impact on their performance. Figure 5a shows that the introduction of an even, very thin layer of c-TiO₂ into the cells results in a step increase of all basic photoelectric parameters compared to the device where perovskite was deposited directly on FTO. However, TiO₂ electron transport layers below ca. 15 nm are too thin, which lowers the cells' efficiencies. Additionally, part of such cells turned out to be faulty, and the devices were difficult to repeat. Most likely, such c-TiO₂ layers may not be completely continuous, and the pinholes present result in the fast electron–hole pair recombination at the FTO/MAPbI₃ interface. This can also provide an alternate current path for the light-generated current, lowering the Shunt resistance of the cell. The optimal ETL thickness for the studied perovskite-based solar cells was 19.5 nm, resulting in the highest power conversion efficiency (PCE) of 13.6%. Along with thickening

of the c-TiO₂ layer, the fill factor (FF) remains at a similar level (0.60–0.65), the open-circuit voltage (V_{OC}) slightly decreases from ca. 950 mV to ca. 800 mV, while the short-circuit current density (J_{SC}) drops significantly from 19.0 to 13.1 mA·cm⁻². This drop cannot be explained only by the lower effective transmittance (Figure 3b), which indicates the hindered transport properties. It is worth noting that the highest J_{SC} of 21.3 mA·cm⁻² was recorded for the cells with a 15.9 nm c-TiO₂ layer; however, the remaining parameters (FF and V_{OC}) were significantly lower than the cells with the optimal c-TiO₂ layer. As a result, we observe a gradual decrease in PCE from 13.6% for the 19.5 nm c-TiO₂ layer to 7.3% for the 124.2 nm layer. The selected electrical parameters are collected in Table 1, while the remaining are available in the Supporting Information (Table S2).



Figure 4. The dependence of the root mean square surface roughness (Rq) on the thickness of the c-TiO₂ layer deposited on the glass/FTO substrate. It was collated with the selected AFM images of the surface roughness and morphology of the glass/FTO/c-TiO₂ substrate, where the sampling area was 2 μ m × 2 μ m. The color scale for height is also shown.

Table 1. Average electrical parameters and their standard deviations for MAPbI₃ perovskite solar cells with a compact TiO_2 layer of different thicknesses (forward scan). The samples are named after the TTDB molar concentration and centrifugation speed.

Sample	d _{c-TiO2} (nm)	J_{SC} (mA·cm ⁻²)	V _{OC} (mV)	FF (—)	PCE (%)
Glass/FTO	_	4.2 ± 5.2	188 ± 5.16	0.29 ± 0.04	0.69 ± 0.80
0.1M_2kRPM	9.1	19.1 ± 0.3	891 ± 20	0.57 ± 0.02	11.12 ± 0.62
0.2M_3kRPM	15.9	21.3 ± 0.4	892 ± 28	0.62 ± 0.03	13.36 ± 1.23
0.2M_2kRPM	19.5	19.0 ± 0.2	949 ± 23	0.66 ± 0.02	13.59 ± 0.63
0.5M_4kRPM	36.9	17.4 ± 0.5	947 ± 26	0.63 ± 0.02	11.87 ± 0.61
0.75M_4kRPM	59.4	16.4 ± 0.3	906 ± 11	0.64 ± 0.01	10.81 ± 0.22
1M_2kRPM	124.2	13.1 ± 0.3	791 ± 43	0.62 ± 0.03	7.29 ± 0.88



Figure 5. (a) Influence of the c-TiO₂ layer thickness on basic photoelectric parameters: short-circuit current density (J_{SC}), open-circuit voltage (V_{OC}), power conversion efficiency (PCE) and fill factor (FF). (b) Averaged current density–voltage (J–V) curves for the selected cells with different c-TiO₂ thicknesses (scan speed = 1V/s, light soaking 10 s before measurement).

A few selected J–V curves for the studied cells are presented in Figure 5b. They directly reflect the trends observed for the electrical parameters (Figure 5a, Table 1). The shapes of the J–V curves are similar for all cells, which is reflected by similar values of FF. All devices exhibit hysteresis typical for perovskite solar cells between forward and reverse scans, the origins of which are due to ferroelectric polarization, ion migration, charge trapping and/or capacitive effects [43]. In the case of planar heterojunction PSCs under study, the hysteresis is much more pronounced for those with a thinner ETL. Although cells with a c-TiO₂ layer of 9.1 and 19.5 nm show a similarly high J_{SC}, a significantly lower V_{OC} is observed for the thinner layer. It is also clear that the V_{OC} decreases as the c-TiO₂ thickens.

There may be several reasons for the trends in the electrical properties of the studied PSCs. The thick ETL acts as an insulator for the electrons generated in the absorber, increasing the series resistance of the cell. This hinders charge collection by lowering the gradient of the carrier concentration (diffusion-based transport) and also by disturbing the electric field in the material (drift-based transport).

4. Conclusions

The presented results clearly shows that the thickness of the compact TiO₂ ETL clearly affects the performance of MAPbI₃ perovskite solar cells. We chose the simplest planar heterojunction architecture (glass/FTO/b-TiO₂/MAPbI₃/Spiro-OMeTAD/Au) to minimize the influence of factors that could interfere with the results, e.g., the presence of a mesoporous TiO₂ scaffold.

The ETL can significantly reduce the roughness of the FTO substrates. Even very thin c-TiO₂ layers (6.7 nm) dull the sharp grains of FTO, while the thicker layers smooth the entire surface. Compact TiO₂ also affects the optical transparency of the devices. The effective transmittance of glass/FTO substrates slowly decreases with the thickening of the ETL, but layers thinner than ca. 30 nm still hardly affected it.

More significant effects were observed for the current density–voltage characteristics of the cells. We found that $c-TiO_2$ of ca. 20 nm is optimal, providing the best performance of the devices. Thinner layers may result in worse efficiency of cells or difficulties with the reproducibility of the devices. An ETL that is too thin may not be completely continuous and contain pinholes, which may result in fast electron-hole pair recombination at the FTO/MAPbI₃ interface. It can also provide an alternate current path for the light-generated current, lowering the Shunt resistance of the cell. It is worth noting that even the very thin layer of c-TiO₂ markedly increases all basic photoelectric parameters compared to the device with perovskite deposited directly on FTO. On the other hand, a TiO_2 layer that is too thick also has a negative impact on the performance of PSCs, as it significantly reduces the short-circuit current density (J_{SC}) and also the open-circuit voltage (V_{OC}). Along with the thickening of the c-TiO₂ layer from ca. 20 nm to ca. 125 nm, J_{SC} decreased by 38% (from 21.3 to 13.1 mA·cm⁻²) and V_{OC} by 16% (from 950 mV to 800 mV). Yet, for most cells, the FF remained at a similar level, between 0.6 and 0.65. In effect, we noticed a clear, gradual decrease in power conversion efficiency (PCE) from 13.6% to 7.3% along with the thickening of the c-TiO₂ ETL.

Supplementary Materials: The following are available online at https://www.mdpi.com/article/10 .3390/ma14123295/s1, Table S1: Parameters from ellipsometry measurements for the studied c-TiO₂ layers on polished silicon wafers, Figure S1: Spin-coating program for the deposition of the MAPbI₃ perovskite precursor, Figure S2: UV–vis–NIR spectra of the glass/FTO/c-TiO₂ samples under study, Table S2: Electrical parameters of the studied FTO/b-TiO2/MAPbI₃/Spiro-OMeTAD/Au solar cells: short-circuit current density (J_{SC}), open-circuit voltage (V_{OC}), fill factor (FF) and power conversion efficiency (PCE), Figure S3: SEM images showing the morphology of the MAPbI₃ perovskite layer deposited on glass/FTO/c-TiO₂.

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