

Weight Function Method for Stress Intensity Factors of Semi-Elliptical Surface Cracks on Functionally Graded Plates Subjected to Non-Uniform Stresses

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Nomenclature

a	crack depth of a semi-elliptical surface crack
c	half crack length of a semi-elliptical surface crack
D_1, D_2, D_3	weight function coefficients
D_{D1}, D_{D2}, D_{D3}	weight function coefficients for deepest point
D_{S1}, D_{S2}, D_{S3}	weight function coefficients for surface point
D_{P1}, D_{P2}, D_{P3}	weight function coefficients for general point
E	Young's modulus
E_0, E_1	Young's modulus of starting face constituent, ending face constituent
E_{tip}	Young's modulus at the crack tip
E'_{tip}	modified Young's modulus at the crack tip
F	boundary correction factors
F_0, F_1	boundary correction factors for reference stress intensity factors K_{r1}^S and K_{r2}^S
h	half height of functionally graded plate
K	stress intensity factor
$K_r(a)$	reference stress intensity factor related to crack length
K_{r1}^D, K_{r2}^D	reference stress intensity factors of deepest point
K_{r1}^S, K_{r2}^S	reference stress intensity factors of surface point
K_{r1}^P, K_{r2}^P	reference stress intensity factors of general point
Q	shape factor for an ellipse
t	thickness of functionally graded plate
w	half width of functionally graded plate
Y_0, Y_1	boundary correction factors for reference stress intensity factors K_{r1}^D and K_{r2}^D
Z_0, Z_1	boundary correction factors for reference stress intensity factors K_{r1}^P and K_{r2}^P
σ_0	nominal or characteristic stress
$\sigma(x)$	local stress distribution normal to the prospective crack face
ν	Poisson's ratio
ϕ	parametric angle of an elliptical surface crack

1. Detailed Derivation Process and Explanation

1.1. Detailed Explanation of Equation (15) in the Manuscript.

The relationship between the crack opening displacement $u(x, a)$ and the weight function $m(x, a)$ was derived by Rice [5], and it is expressed as follows:

$$m(x, a) = \frac{E_{\text{tip}}}{K_r(a)} \frac{\partial u(x, a)}{\partial a} \quad (\text{S1})$$

Thus, the first derivative of the weight function with respect to x can be written in the following form.

$$\frac{\partial m(x, a)}{\partial x} = \frac{E_{\text{tip}}}{K_r(a)} \frac{\partial}{\partial a} \left[\frac{\partial u(x, a)}{\partial x} \right] \quad (\text{S2})$$

The second derivative of the weight function with respect to x can be written in the following form.

$$\frac{\partial^2 m(x, a)}{\partial x^2} = \frac{E_{\text{tip}}}{K_r(a)} \frac{\partial}{\partial a} \left[\frac{\partial^2 u(x, a)}{\partial x^2} \right] \quad (\text{S3})$$

The following derivation of the additional condition for a surface crack with depth a is from the reference [24]. Let us consider displacements u and v in the vicinity of the point $(0, 0)$. Because the x -axis is an axis of symmetry, no shear stresses will be acting there.

$$\tau(x, 0) = 0 \quad (\text{S4})$$

Along the free surface of the plate it holds:

$$\tau(0, y) = \sigma_x(0, y) = 0 \quad (\text{S5})$$

For stresses σ_x and τ developable by power series as:

$$\sigma_x = \sum_{v=0}^{\infty} \sum_{u=0}^{\infty} A_{uv} x^u y^v \quad \tau = \sum_{v=0}^{\infty} \sum_{u=0}^{\infty} B_{uv} x^u y^v \quad (\text{S6})$$

The following equation is obtained according to conditions (S4) and (S5):

$$B_{u0} = B_{0v} = A_{0v} = 0 \quad (\text{S7})$$

The relationship between deflections u , v and shear distortion γ is:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = \gamma \quad (\text{S8})$$

The following equation is obtained by taking the derivative of (S8) with respect to x and using the strain component $\varepsilon_x = \partial v / \partial x$.

$$\frac{\partial^2 u}{\partial x^2} = \frac{\partial \gamma}{\partial x} - \frac{\partial^2 v}{\partial x \partial y} = \frac{\partial \gamma}{\partial x} - \frac{\partial \varepsilon_x}{\partial y} \quad (\text{S9})$$

By using Hooke's law, we obtain:

$$\varepsilon_x = m \sigma_x + n \sigma_y; \quad \gamma = \tau / G \quad (\text{S10})$$

$$m = \begin{cases} 1/E \\ (1-\nu^2)/E \end{cases}; \quad n = \begin{cases} -\nu/E & \text{for plane stress} \\ -\nu(1+\nu)/E & \text{for plane strain} \end{cases}$$

The equilibrium condition gives:

$$\frac{\partial \sigma_y}{\partial y} + \frac{\partial \tau}{\partial x} = 0 \quad (\text{S11})$$

The following expression is obtained from Equations (S9), (S10) and (S11).

$$\frac{\partial^2 u}{\partial x^2} = \left(\frac{1}{G} + n \right) \frac{\partial \tau}{\partial x} - m \frac{\partial \sigma_x}{\partial y} \quad (\text{S12})$$

The second derivative of u at point $(0, 0)$ is as follows:

$$\frac{\partial^2 u}{\partial x^2} \Big|_{x=y=0} = \left(\frac{1}{G} + n \right) B_{10} - mA_{01} = 0 \quad (\text{S13})$$

That is, directly at the surface of the plate, the curvature of the crack contour disappears [24].

Based on Equations (S5), (S6), (S7) and (S12), it can be proved that the curvature of the crack contour at the surface ($x = 0$) vanishes, as follows.

$$\frac{\partial^2 u(x, a)}{\partial x^2} \Big|_{x=0} = 0 \quad (\text{S14})$$

Consequently, the second derivative of the weight function at $x = 0$ must also be zero. Thus, in the case of an edge or surface crack, Equation (S3) can be written as follows [6,24]:

$$\frac{\partial^2 m_D(x, a)}{\partial x^2} \Big|_{x=0} = 0 \quad (\text{S15})$$

1.2. Detailed Explanation of Equation (17) in the Manuscript.

The explanation of the sentence “Due to the weight function for the surface point of a semi-elliptical surface crack is derived from the weight function for the embedded penny-shape crack, therefore, the weight function in Equation (11) must vanish at $x=a$ [25]” in the manuscript is as follows.

The closed form weight function for an embedded circular crack (embedded penny-shape crack) is given [25]:

$$m_F(a, x, \theta) = \frac{1}{\pi \sqrt{\pi a}} \frac{\sqrt{(a^2 - x^2)}}{a^2 + x^2 - 2ax \cos \theta} \quad (\text{S16})$$

Shen et al. [25] derived the weight function for the surface point B of a semi-elliptical surface crack from Equation (S16); it is expressed as follows.

$$m_B(x, a) = \frac{2}{\sqrt{\pi x}} \left(1 - \sqrt{\frac{x}{a}} \right) = \frac{2}{\sqrt{\pi x}} \left[1 - \left(\frac{x}{a} \right)^{1/2} \right] \quad (\text{S17})$$

The weight function for the surface point in the manuscript is given by analogy with the equation (21) in reference [25].

$$m_S(x, a) = \frac{2}{\sqrt{\pi x}} \left(\frac{x}{a} \right)^{-1/2} \left[1 + D_{S1} \left(\frac{x}{a} \right) + D_{S2} \left(\frac{x}{a} \right)^2 + D_{S3} \left(\frac{x}{a} \right)^3 \right] \quad (\text{S18})$$

Since the equation (21) in the reference must satisfy the condition that the weight function is zero at the crack tip ($x = a$) [25]; therefore, the weight functions in Equations (S17) and (S18) are equal to zero at $x=a$, leading to:

$$m_S(x, a) \Big|_{x=a} = 0 \quad (\text{S19})$$

1.3. Detailed Explanation of Equations (19) and (20) in the Manuscript.

The deepest point ($\phi = \pi/2$) and surface point ($\phi = 0$) are special cases of general points [26]. Fett et al. [24] pointed out that for a surface crack with depth a , the curvature of the crack contour at the surface ($x = 0$) vanishes. That is, the condition that the curvature of the crack contour at $x = 0$ is zero should be satisfied if the general point infinitely approaches the deepest point ($\phi \rightarrow \pi/2$).

$$\left. \frac{\partial^2 u(x, a, \phi)}{\partial x^2} \right|_{x=0, \phi \rightarrow \pi/2} = 0 \quad (S20)$$

Consequently, the second derivative of the weight function of the general point at $x = 0$ should also be zero.

$$\left. \frac{\partial^2 m_{p1}(x, a, \phi)}{\partial x^2} \right|_{x=0, \phi \rightarrow \pi/2} = 0 \quad (S21)$$

Finally, we obtain:

$$\left. \frac{\partial^2 m_{p1}(x, a)}{\partial x^2} \right|_{x=0} = 0 \quad (S22)$$

As explained in the previous derivation, the weight function for the surface point must be zero at the crack tip ($x = a$) [25]. The condition that the weight function of the surface point is zero at $x = a$ should be satisfied if the general point is infinitely approaching the surface point ($\phi \rightarrow 0$).

$$m_{p2}(x, a, \phi) \Big|_{x=a, \phi \rightarrow 0} = 0 \quad (S23)$$

Consequently, the Equation (19) in the manuscript is obtained:

$$m_{p2}(x, a) \Big|_{x=a} = 0 \quad (S24)$$

Notice: For references cited in supplementary materials, please refer to the corresponding references in the manuscript.

2. Detailed Derivation Process of Equations (41), (42), (43) and (44)

$$\begin{aligned} & \left. \frac{\partial}{\partial x^2} \left\{ \sqrt{\frac{2}{\pi a \sin \phi}} \left[\left(1 - \frac{x}{a \sin \phi}\right)^{-\frac{1}{2}} + D_{p1} \left(1 - \frac{x}{a \sin \phi}\right)^{\frac{1}{2}} + D_{p2} \left(1 - \frac{x}{a \sin \phi}\right)^{\frac{3}{2}} \right] \right\} \right|_{x=0} \\ &= \left\{ \sqrt{\frac{2}{\pi a \sin \phi}} \frac{\partial}{\partial x} \left[-\frac{1}{2} \left(-\frac{1}{a \sin \phi}\right) \left(1 - \frac{x}{a \sin \phi}\right)^{-\frac{3}{2}} + \frac{1}{2} D_{p1} \left(-\frac{1}{a \sin \phi}\right) \left(1 - \frac{x}{a \sin \phi}\right)^{-\frac{1}{2}} \right. \right. \\ & \quad \left. \left. + \frac{3}{2} D_{p2} \left(-\frac{1}{a \sin \phi}\right) \left(1 - \frac{x}{a \sin \phi}\right)^{\frac{1}{2}} \right] \right\} \Big|_{x=0} \\ &= \left\{ \sqrt{\frac{2}{\pi a \sin \phi}} \left[\frac{3}{4} \left(-\frac{1}{a \sin \phi}\right)^2 \left(1 - \frac{x}{a \sin \phi}\right)^{-\frac{5}{2}} - \frac{1}{4} D_{p1} \left(-\frac{1}{a \sin \phi}\right)^2 \left(1 - \frac{x}{a \sin \phi}\right)^{-\frac{3}{2}} \right. \right. \\ & \quad \left. \left. + \frac{3}{4} D_{p2} \left(-\frac{1}{a \sin \phi}\right)^2 \left(1 - \frac{x}{a \sin \phi}\right)^{-\frac{1}{2}} \right] \right\} \Big|_{x=0} = 0 \\ \Rightarrow \quad & 3 - D_{p1} + 3D_{p2} = 0 \quad (41) \end{aligned}$$

$$\left. \left\{ \sqrt{\frac{2}{\pi a \sin \phi}} \left[\left(\frac{x}{a \sin \phi} - 1 \right)^{-\frac{1}{2}} + D_{P3} \left(\frac{x}{a \sin \phi} - 1 \right)^{\frac{1}{2}} + D_{P4} \left(\frac{x}{a \sin \phi} - 1 \right)^{\frac{3}{2}} \right] \right\} \right|_{x=a}$$

$$= \sqrt{\frac{2}{\pi a \sin \phi}} \left[\left(\frac{1}{\sin \phi} - 1 \right)^{-\frac{1}{2}} + D_{P3} \left(\frac{1}{\sin \phi} - 1 \right)^{\frac{1}{2}} + D_{P4} \left(\frac{1}{\sin \phi} - 1 \right)^{\frac{3}{2}} \right] = 0$$

$$\Rightarrow 1 + D_{P3} \left(\frac{1}{\sin \phi} - 1 \right) + D_{P4} \left(\frac{1}{\sin \phi} - 1 \right)^2 = 0 \quad (42)$$

$$K_{I_r}^P = \int_0^{a \sin \phi} \sigma_0 \left\{ \sqrt{\frac{2}{\pi a \sin \phi}} \left[\left(1 - \frac{x}{a \sin \phi} \right)^{-\frac{1}{2}} + D_{P1} \left(1 - \frac{x}{a \sin \phi} \right)^{\frac{1}{2}} + D_{P2} \left(1 - \frac{x}{a \sin \phi} \right)^{\frac{3}{2}} \right] \right\} dx$$

$$+ \int_{a \sin \phi}^a \sigma_0 \left\{ \sqrt{\frac{2}{\pi a \sin \phi}} \left[\left(\frac{x}{a \sin \phi} - 1 \right)^{-\frac{1}{2}} + D_{P3} \left(\frac{x}{a \sin \phi} - 1 \right)^{\frac{1}{2}} + D_{P4} \left(\frac{x}{a \sin \phi} - 1 \right)^{\frac{3}{2}} \right] \right\} dx$$

$$= \sigma_0 \sqrt{\frac{2}{\pi a \sin \phi}} \left[-2a \sin \phi \left(1 - \frac{x}{a \sin \phi} \right)^{\frac{1}{2}} - \frac{2}{3} a \sin \phi D_{P1} \left(1 - \frac{x}{a \sin \phi} \right)^{\frac{3}{2}} \right. \\ \left. - \frac{2}{5} a \sin \phi D_{P2} \left(1 - \frac{x}{a \sin \phi} \right)^{\frac{5}{2}} \right] \Big|_0^{a \sin \phi}$$

$$+ \sigma_0 \sqrt{\frac{2}{\pi a \sin \phi}} \left[2a \sin \phi \left(\frac{x}{a \sin \phi} - 1 \right)^{\frac{1}{2}} + \frac{2}{3} a \sin \phi D_{P3} \left(\frac{x}{a \sin \phi} - 1 \right)^{\frac{3}{2}} \right. \\ \left. + \frac{2}{5} a \sin \phi D_{P4} \left(\frac{x}{a \sin \phi} - 1 \right)^{\frac{5}{2}} \right] \Big|_{a \sin \phi}^a$$

$$= \sigma_0 \sqrt{\frac{2}{\pi a \sin \phi}} \left[2a \sin \phi + \frac{2}{3} a \sin \phi D_{P1} + \frac{2}{5} a \sin \phi D_{P2} \right]$$

$$+ \sigma_0 \sqrt{\frac{2}{\pi a \sin \phi}} \left[2a \sin \phi \left(\frac{1}{\sin \phi} - 1 \right)^{\frac{1}{2}} + \frac{2}{3} a \sin \phi D_{P3} \left(\frac{1}{\sin \phi} - 1 \right)^{\frac{3}{2}} \right. \\ \left. + \frac{2}{5} a \sin \phi D_{P4} \left(\frac{1}{\sin \phi} - 1 \right)^{\frac{5}{2}} \right] = \sigma_0 \sqrt{\frac{\pi a}{Q}} Z_0$$

$$\Rightarrow \left[1 + \frac{1}{3} D_{P1} + \frac{1}{5} D_{P2} \right] + \left[\left(\frac{1}{\sin \phi} - 1 \right)^{\frac{1}{2}} + \frac{1}{3} D_{P3} \left(\frac{1}{\sin \phi} - 1 \right)^{\frac{3}{2}} + \frac{1}{5} D_{P4} \left(\frac{1}{\sin \phi} - 1 \right)^{\frac{5}{2}} \right]$$

$$= \pi \sqrt{\frac{1}{8Q \sin \phi}} Z_0 = \sqrt{\frac{1}{Q \sin \phi}} V_0 \quad (43)$$

$$\begin{aligned}
 K_{2r}^p &= \int_0^{a \sin \phi} \sigma_0 \left(1 - \frac{x}{a}\right) \left\{ \sqrt{\frac{2}{\pi a \sin \phi}} \left[\left(1 - \frac{x}{a \sin \phi}\right)^{-\frac{1}{2}} + D_{p1} \left(1 - \frac{x}{a \sin \phi}\right)^{\frac{1}{2}} + D_{p2} \left(1 - \frac{x}{a \sin \phi}\right)^{\frac{3}{2}} \right] \right\} dx \\
 &+ \int_{a \sin \phi}^a \sigma_0 \left(1 - \frac{x}{a}\right) \left\{ \sqrt{\frac{2}{\pi a \sin \phi}} \left[\left(\frac{x}{a \sin \phi} - 1\right)^{-\frac{1}{2}} + D_{p3} \left(\frac{x}{a \sin \phi} - 1\right)^{\frac{1}{2}} + D_{p4} \left(\frac{x}{a \sin \phi} - 1\right)^{\frac{3}{2}} \right] \right\} dx \\
 &= \int_0^{a \sin \phi} \sigma_0 \left(1 - \frac{x}{a}\right) \left\{ \sqrt{\frac{2}{\pi a \sin \phi}} \left[\left(1 - \frac{x}{a \sin \phi}\right)^{-\frac{1}{2}} + D_{p1} \left(1 - \frac{x}{a \sin \phi}\right)^{\frac{1}{2}} + D_{p2} \left(1 - \frac{x}{a \sin \phi}\right)^{\frac{3}{2}} \right] \right\} dx \\
 &= \left\{ \sigma_0 \left(1 - \frac{x}{a}\right) \sqrt{\frac{2}{\pi a \sin \phi}} \left[-2a \sin \phi \left(1 - \frac{x}{a \sin \phi}\right)^{\frac{1}{2}} - \frac{2}{3} a \sin \phi D_{p1} \left(1 - \frac{x}{a \sin \phi}\right)^{\frac{3}{2}} \right. \right. \\
 &\quad \left. \left. - \frac{2}{5} a \sin \phi D_{p2} \left(1 - \frac{x}{a \sin \phi}\right)^{\frac{5}{2}} \right] \right\} \Bigg|_0^{a \sin \phi} \\
 &+ \frac{\sigma_0}{a} \int_0^{a \sin \phi} \left\{ \sqrt{\frac{2}{\pi a \sin \phi}} \left[-2a \sin \phi \left(1 - \frac{x}{a \sin \phi}\right)^{\frac{1}{2}} - \frac{2}{3} a \sin \phi D_{p1} \left(1 - \frac{x}{a \sin \phi}\right)^{\frac{3}{2}} \right. \right. \\
 &\quad \left. \left. - \frac{2}{5} a \sin \phi D_{p2} \left(1 - \frac{x}{a \sin \phi}\right)^{\frac{5}{2}} \right] \right\} dx \\
 &= \sigma_0 \sqrt{\frac{2}{\pi a \sin \phi}} \left[2a \sin \phi + \frac{2}{3} a \sin \phi D_{p1} + \frac{2}{5} a \sin \phi D_{p2} \right] \\
 &- \frac{\sigma_0}{a} \sqrt{\frac{2}{\pi a \sin \phi}} \left[\frac{4}{3} a^2 \sin^2 \phi + \frac{4}{15} a^2 \sin^2 \phi D_{p1} + \frac{4}{35} a^2 \sin^2 \phi D_{p2} \right] \\
 &= a \sigma_0 \sqrt{\frac{2}{\pi a \sin \phi}} \left[2 \sin \phi + \frac{2}{3} \sin \phi D_{p1} + \frac{2}{5} \sin \phi D_{p2} \right] \\
 &- a \sigma_0 \sqrt{\frac{2}{\pi a \sin \phi}} \left[\frac{4}{3} \sin^2 \phi + \frac{4}{15} \sin^2 \phi D_{p1} + \frac{4}{35} \sin^2 \phi D_{p2} \right]
 \end{aligned}$$

$$\begin{aligned}
& \int_{a \sin \phi}^a \sigma_0 \left(1 - \frac{x}{a}\right) \left\{ \sqrt{\frac{2}{\pi a \sin \phi}} \left[\left(\frac{x}{a \sin \phi} - 1\right)^{-\frac{1}{2}} + D_{P_3} \left(\frac{x}{a \sin \phi} - 1\right)^{\frac{1}{2}} + D_{P_4} \left(\frac{x}{a \sin \phi} - 1\right)^{\frac{3}{2}} \right] \right\} dx \\
&= \left\{ \sigma_0 \left(1 - \frac{x}{a}\right) \sqrt{\frac{2}{\pi a \sin \phi}} \left[2a \sin \phi \left(\frac{x}{a \sin \phi} - 1\right)^{\frac{1}{2}} + \frac{2}{3} a \sin \phi D_{P_3} \left(\frac{x}{a \sin \phi} - 1\right)^{\frac{3}{2}} \right. \right. \\
&\quad \left. \left. + \frac{2}{5} a \sin \phi D_{P_4} \left(\frac{x}{a \sin \phi} - 1\right)^{\frac{5}{2}} \right] \right\} \Bigg|_{a \sin \phi}^a \\
&+ \frac{\sigma_0}{a} \int_{a \sin \phi}^a \left\{ \sqrt{\frac{2}{\pi a \sin \phi}} \left[2a \sin \phi \left(\frac{x}{a \sin \phi} - 1\right)^{\frac{1}{2}} + \frac{2}{3} a \sin \phi D_{P_3} \left(\frac{x}{a \sin \phi} - 1\right)^{\frac{3}{2}} \right. \right. \\
&\quad \left. \left. + \frac{2}{5} a \sin \phi D_{P_4} \left(\frac{x}{a \sin \phi} - 1\right)^{\frac{5}{2}} \right] \right\} dx \\
&= a \sigma_0 \sqrt{\frac{2}{\pi a \sin \phi}} \left[\frac{4}{3} \sin^2 \phi \left(\frac{x}{a \sin \phi} - 1\right)^{\frac{3}{2}} + \frac{4}{15} \sin^2 \phi D_{P_3} \left(\frac{x}{a \sin \phi} - 1\right)^{\frac{5}{2}} \right. \\
&\quad \left. + \frac{4}{35} \sin^2 \phi D_{P_4} \left(\frac{x}{a \sin \phi} - 1\right)^{\frac{7}{2}} \right] \Bigg|_{a \sin \phi}^a \\
&\Rightarrow \left[1 + \frac{1}{3} D_{P_1} + \frac{1}{5} D_{P_2} \right] - \left[\frac{2}{3} \sin \phi + \frac{2}{15} \sin \phi D_{P_1} + \frac{2}{35} \sin \phi D_{P_2} \right] \\
&+ \left[\frac{2}{3} \sin \phi \left(\frac{1}{\sin \phi} - 1\right)^{\frac{3}{2}} + \frac{2}{15} \sin \phi D_{P_3} \left(\frac{1}{\sin \phi} - 1\right)^{\frac{5}{2}} + \frac{2}{35} \sin \phi D_{P_4} \left(\frac{1}{\sin \phi} - 1\right)^{\frac{7}{2}} \right] \\
&= \pi \sqrt{\frac{1}{8Q \sin \phi}} Z_1 = \sqrt{\frac{1}{Q \sin \phi}} V_1 \tag{44}
\end{aligned}$$

